

# Fingerprinting Higgs Suspects at the LHC

*Planck 2012  
Warsaw, 31 May 2012*

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CERN-TH

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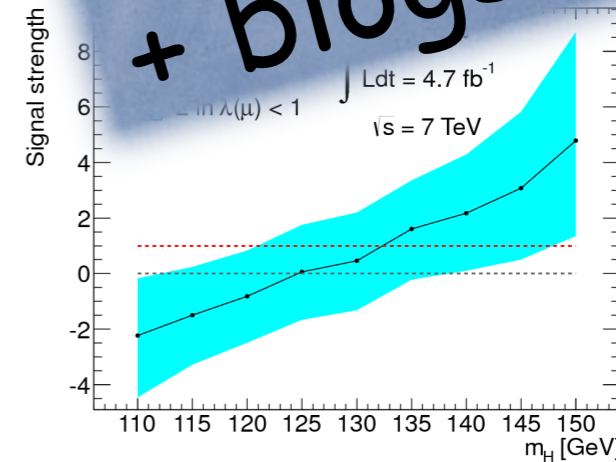
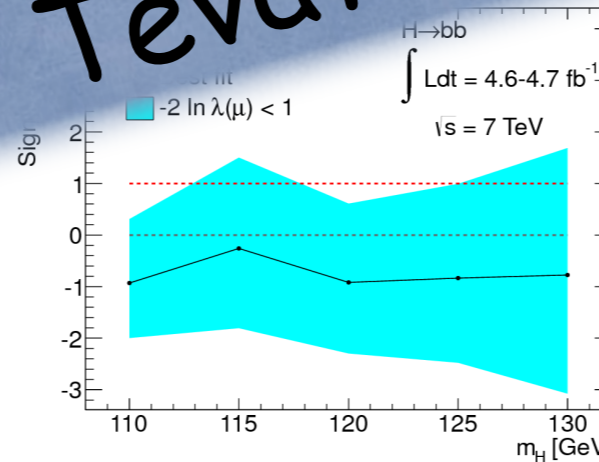
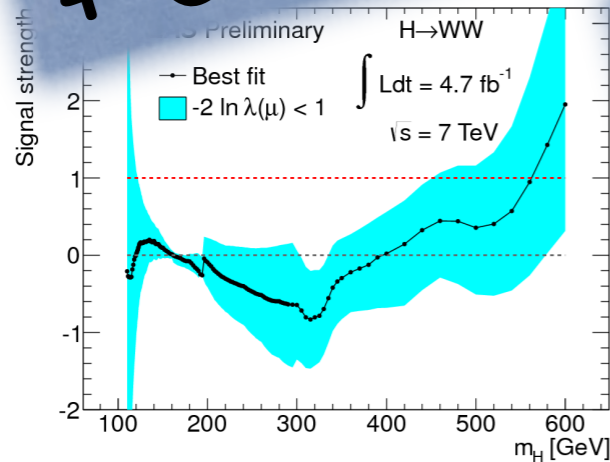
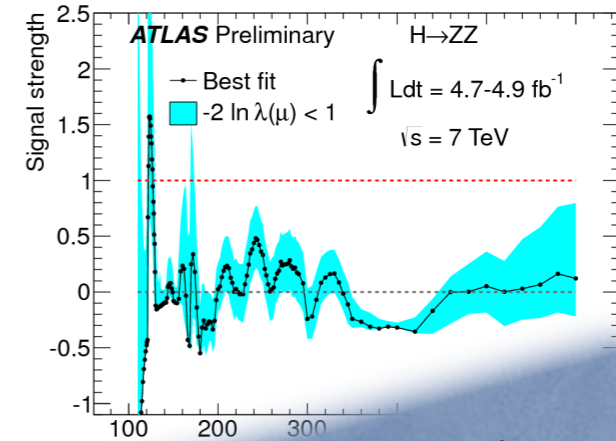
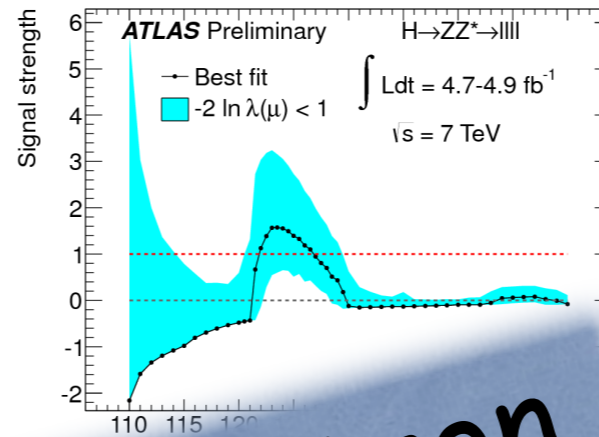
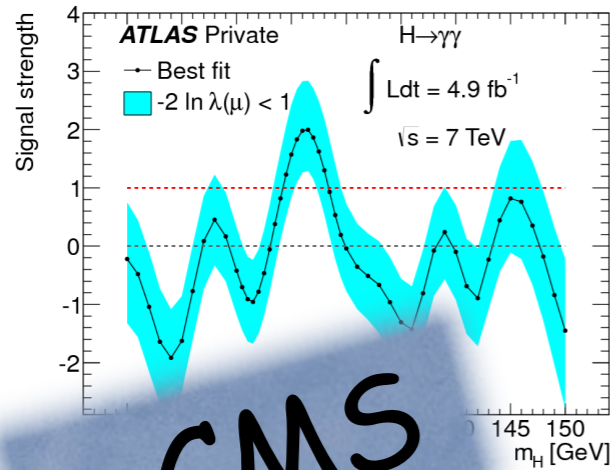
# The Question(s) of the Year

- Is N. Sarkozy going to be re-elected?
- Is the € going to survive?
- Is the Higgs boson going to be discovered?
  
- Where should you invest your money?
  - *Greece?*
  - *Spain/Italy?*
  - *Poland?*
  - *Germany?*
- Which physics to expect Beyond the Standard Model?
  - *Technicolor/Higgsless?*
  - *(C)MSSM?*
  - *(Natural) SUSY?*
  - *Composite Higgs?*
  - *SM?*

# Facts that we have to live with

signal strength

$$\mu_i = \frac{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i)}{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i) |_{\text{SM}}}$$



+ CMS

+ Tevatron

+ blogs...

Still not enough information:

Correlations?  
 Exact likelihoods?



# Chiral Lagrangian for a light Higgs

$$\begin{aligned}
 \mathcal{L} = & \frac{1}{2}(\partial_\mu h)^2 - \frac{1}{2}m_h^2 h^2 - \frac{d_3}{6} \left( \frac{3m_h^2}{v} \right) h^3 - \frac{d_4}{24} \left( \frac{3m_h^2}{v^2} \right) h^4 \dots \\
 & - \left( m_W^2 W_\mu W_\mu + \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right) \left( 1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right) \\
 & - \sum_{\psi=u,d,l} m_{\psi^{(i)}} \bar{\psi}^{(i)} \psi^{(i)} \left( 1 + c_\psi \frac{h}{v} + c_{2\psi} \frac{h^2}{v^2} + \dots \right) \\
 & + \frac{g^2}{16\pi^2} \left( c_{WW} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{ZZ} Z_{\mu\nu}^2 + c_{Z\gamma} Z_{\mu\nu} \gamma_{\mu\nu} \right) \frac{h}{v} + \dots \\
 & + \frac{g^2}{16\pi^2} \left[ \gamma_{\mu\nu}^2 \left( c_{\gamma\gamma} \frac{h}{v} + \dots \right) + G_{\mu\nu}^2 \left( c_{gg} \frac{h}{v} + c_{2gg} \frac{h^2}{v^2} \dots \right) \right] \\
 & + \frac{g^2}{16\pi^2} \left[ \frac{c_{hhgg}}{\Lambda^2} G_{\mu\nu}^2 \frac{(\partial_\rho h)^2}{v^2} + \frac{c'_{hhgg}}{\Lambda^2} G_{\mu\rho} G_{\rho\nu} \frac{\partial_\mu h \partial_\nu h}{v^2} + \dots \right] \\
 & + \dots
 \end{aligned}$$

See Contino's talk

A few (reasonable) assumptions:

□ spin-0 & CP-even



$\gamma\gamma$



WW & ZZ

□ custodial symmetry



EWPD

□ no Higgs FCNC



Flavor

[ RC Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089 ; Azatov, R.C., Galloway, JHEP 1204 (2012) 127 ]



# Chiral Lagrangian for a light Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \frac{1}{2}m_h^2 h^2 - \frac{d_3}{6} \left( \frac{3m_h^2}{v} \right) h^3 - \frac{d_4}{24} \left( \frac{3m_h^2}{v^2} \right) h^4 \dots$$

$$- \left( m_W^2 W_\mu W_\mu + \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right) \left( 1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right)$$

$$- \sum_{\psi=u,d,l} m_{\psi^{(i)}} \bar{\psi}^{(i)} \psi^{(i)} \left( 1 + c_{\psi} \frac{h}{v} + \dots \right)$$

still too much freedom  
 $\Downarrow$   
 dynamical assumptions needed  
 $\Downarrow$   
 to explore deformations of the SM  
 $\left[ c_{gg} \frac{h}{v} + c_{2gg} \frac{h^2}{v^2} \dots \right]$

$$+ \frac{g^2}{16\pi^2} \left[ \frac{c_{hhgg}}{\Lambda^2} G_{\mu\nu}^2 \frac{(\partial_\rho h)^2}{v^2} + \frac{c'_{hhgg}}{\Lambda^2} G_{\mu\rho} G_{\rho\nu} \frac{\partial_\mu h \partial_\nu h}{v^2} + \dots \right]$$

+ ...

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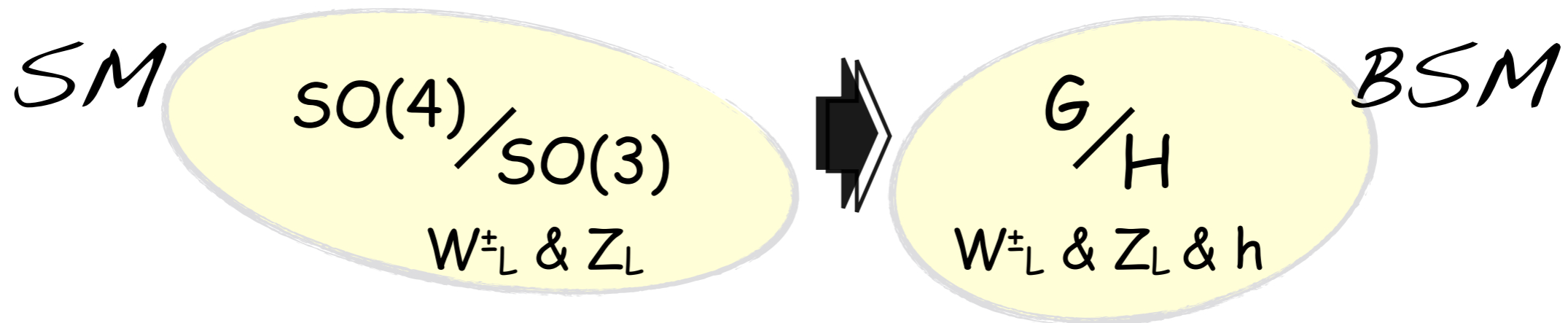


# The New Physics Mass Gap

One solution to the hierarchy pb:

Higgs transforms non-linearly under some global symmetry

Higgs=Pseudo-Goldstone boson (PGB)



Examples:  $SO(5)/SO(4)$ : 4 PGBs =  $W^\pm_L, Z_L, h$

Minimal Composite Higgs Model

Agashe, Contino, Pomarol '04

$SO(6)/SO(5)$ : 5 PGBs =  $H, a$

Next MCHM

Gripaios, Pomarol, Riva, Serra '09

$SU(4)/Sp(4, \mathbb{C})$ : 5 PGBs =  $H, s$

$SO(6)/SO(4) \times SO(2)$ : 8 PGBs =  $H_1 + H_2$

Minimal Composite  
Two Higgs Doublets

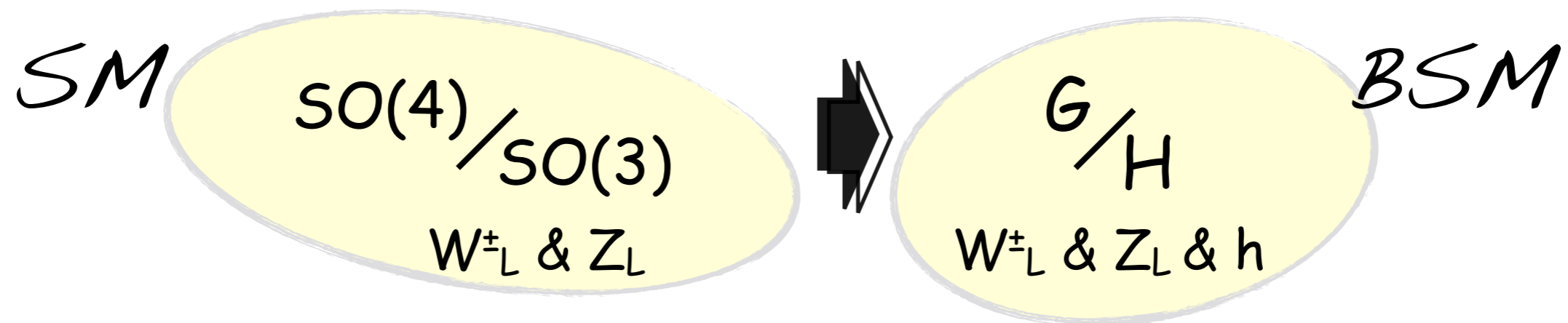
Mrazek, Pomarol, Rattazzi, Serra, Wulzer '11



# The New Physics Mass Gap

One solution to the hierarchy pb:  
Higgs transforms non-linearly under some global symmetry

Higgs=Pseudo-Goldstone boson (PGB)



*How can we tell the difference with the SM Higgs?*

*What are the experimental constraints?*

# SILH Effective Lagrangian

Giudice, Grojean, Pomarol, Rattazzi '07

At the moment, we don't need to know what the Higgs is made of  
 $\Rightarrow$  chiral Lagrangian for the composite Higgs  $\Leftarrow$

## ■ Genuine strong operators (sensitive to the scale $f$ )

$$\frac{c_H}{2f^2} \left( \partial^\mu |H|^2 \right)^2$$

~~$$\frac{c_T}{2f^2} \left( H^\dagger \overleftrightarrow{D}^\mu H \right)^2$$~~  
 custodial breaking

$$\frac{c_y y_f}{f^2} |H|^2 \bar{f}_L H f_R + \text{h.c.}$$

$$\frac{c_6 \lambda}{f^2} |H|^6$$

## ■ Form factor operators (sensitive to the scale $m_\rho$ )

$$\frac{i c_W}{2m_\rho^2} \left( H^\dagger \sigma^i \overleftrightarrow{D}^\mu H \right) (D^\nu W_{\mu\nu})^i$$

$$\frac{i c_B}{2m_\rho^2} \left( H^\dagger \overleftrightarrow{D}^\mu H \right) (\partial^\nu B_{\mu\nu})$$

$$\frac{i c_{HW}}{m_\rho^2} \frac{g_\rho^2}{16\pi^2} (D^\mu H)^\dagger \sigma^i (D^\nu H) W_{\mu\nu}^i$$

$$\frac{i c_{HB}}{m_\rho^2} \frac{g_\rho^2}{16\pi^2} (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu}$$

minimal coupling:  $h \rightarrow \gamma Z$

loop-suppressed strong dynamics

$$\frac{c_\gamma}{m_\rho^2} \frac{g_\rho^2}{16\pi^2} \frac{g^2}{g_\rho^2} H^\dagger H B_{\mu\nu} B^{\mu\nu}$$

$$\frac{c_g}{m_\rho^2} \frac{g_\rho^2}{16\pi^2} \frac{y_t^2}{g_\rho^2} H^\dagger H G_{\mu\nu}^a G^{a\mu\nu}$$

Goldstone sym.



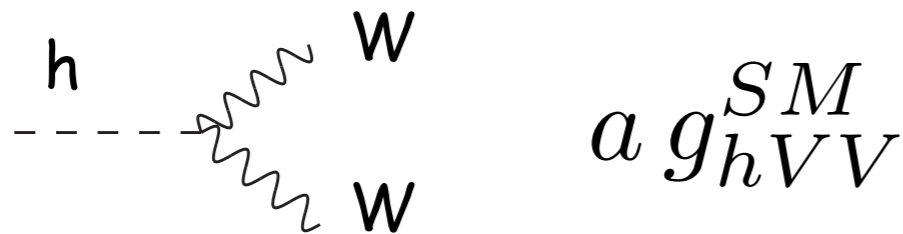
# 2 parameter Higgs physics @ LHC 2011-2012

$$\frac{c_H}{2f^2} (\partial^\mu |H|^2)^2$$

$$\xi = v^2 / f^2$$

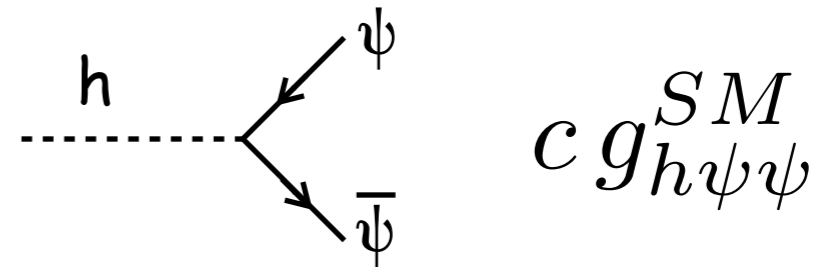
$$\frac{c_y y_\psi}{f^2} |H|^2 \bar{\psi}_L H \psi_R$$

Controls the  $hWW, hZZ$  couplings



$$a = 1 - c_H \xi / 2$$

Controls the  $h\psi\psi$  couplings



$$c = 1 - (c_H + 2c_y) \xi / 2$$

Explicit (and calculable) models built in  $AdS_5$  spacetimes

Agashe, Contino, Pomarol '04

Contino, Da Rold, Pomarol '06

MCHM5

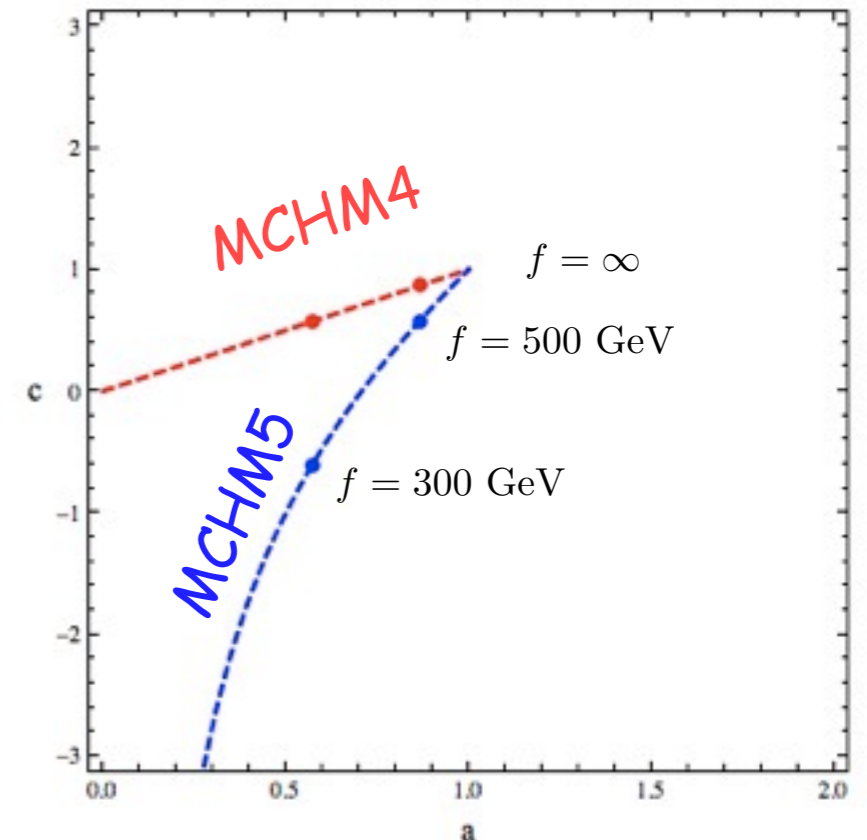
$$a = \sqrt{1 - \xi} \quad c = \frac{1 - 2\xi}{\sqrt{1 - \xi}}$$

MCHM4

$$a = \sqrt{1 - \xi} \quad c = \sqrt{1 - \xi}$$

disfavored by EW data ( $Zbb$ )

SM is recovered as a limit when the compositeness scale is well above weak scale



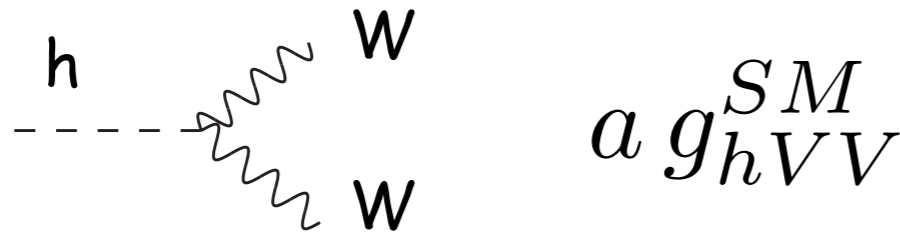
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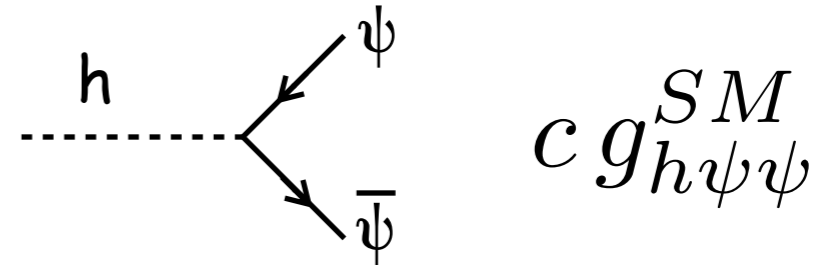
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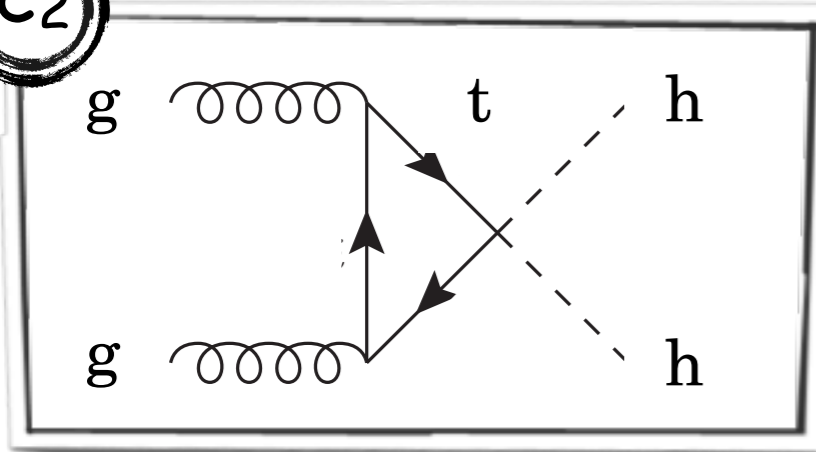
Controls the hψψ couplings



$$c = 1 - (c_H + 2c_y) \xi / 2$$

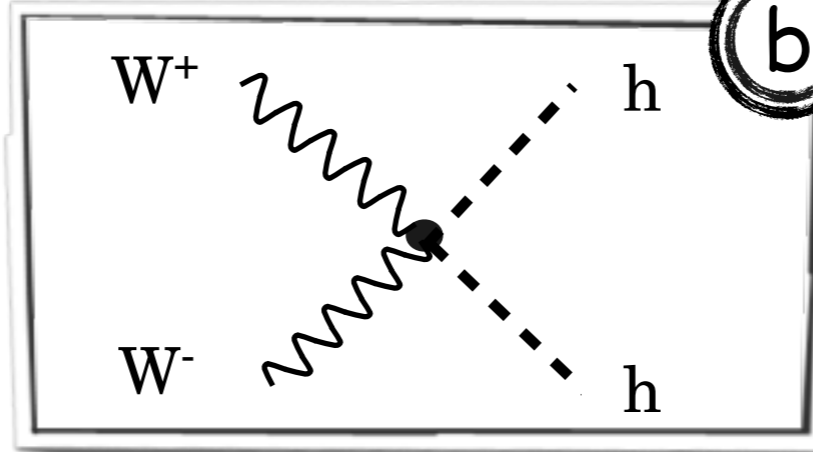
other couplings are very interesting as they are directly testing non-linearities/strong interactions of the Higgs

C2



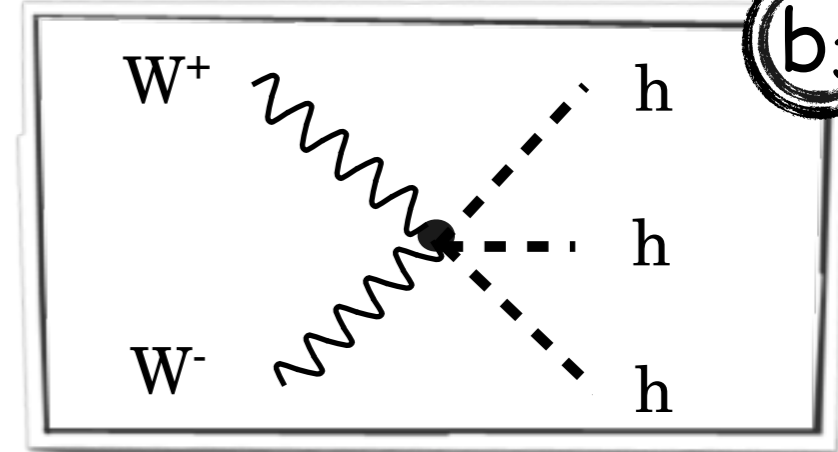
Gröber, Mühlleitner '10  
Contino et al '12  
Gillioz et al 'to appear

b



Contino, Grojean,  
Moretti, Piccinini, Rattazzi '10

b3



Contino, Grojean, Pappadopulo,  
Rattazzi, Thamm 'to appear

but they are not on agenda of the current LHC run



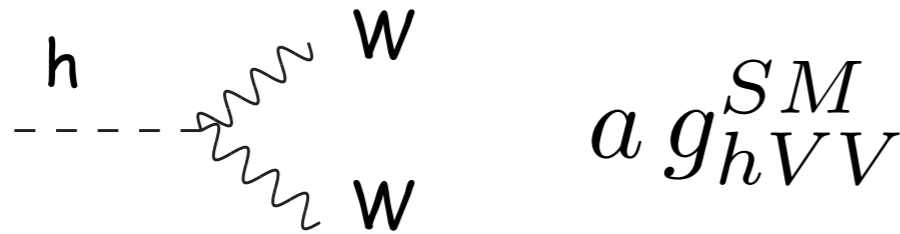
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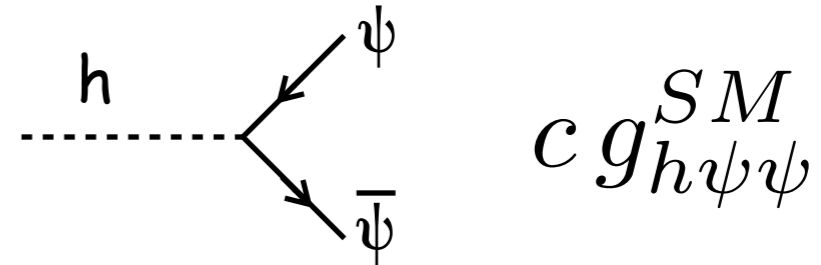
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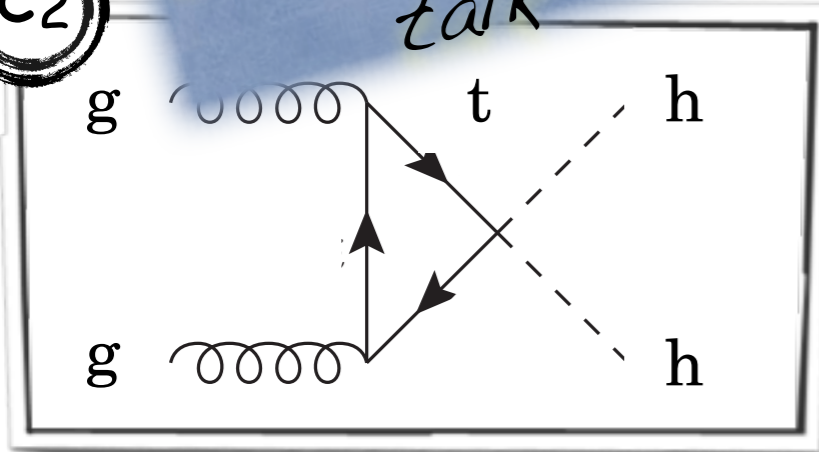
Controls the hψψ couplings



$$c = 1 - (c_H + 2c_y) \xi / 2$$

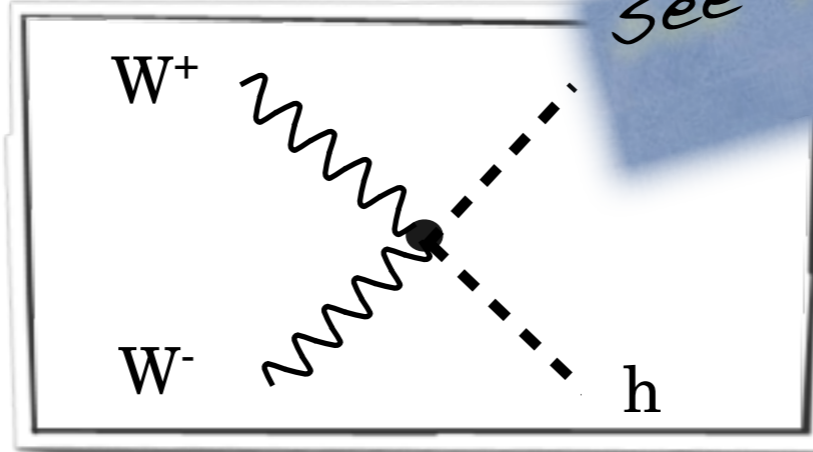
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**C2**



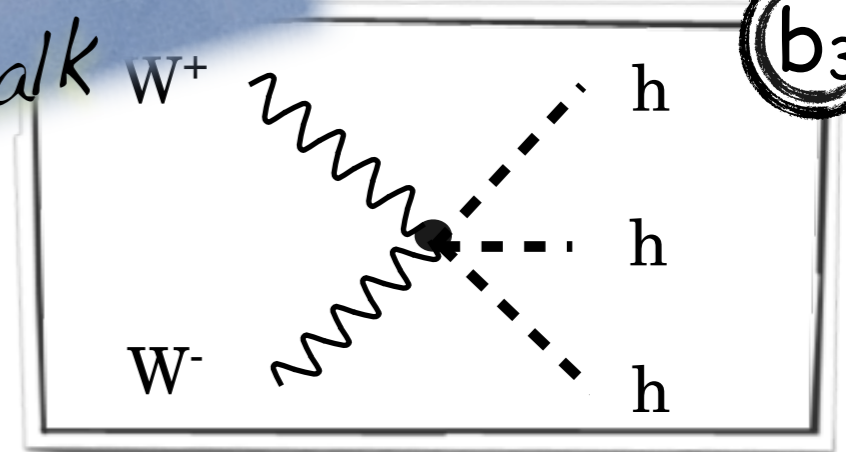
Gröber, Mühlleitner '10  
Contino et al '12  
Gillioz et al 'to appear

see talk



Contino, Grojean,  
Moretti, Piccinini, Rattazzi '10

see talk



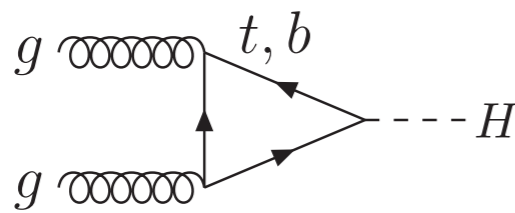
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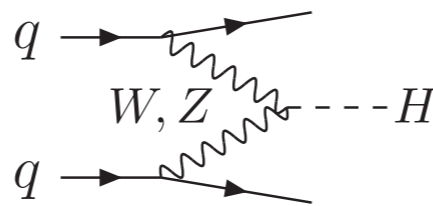
# 2 parameter Higgs physics @ LHC<sub>2011-2012</sub>

- Higgs couplings modified w.r.t. SM but same kinematics  
(particular to single Higgs process - with more than one Higgs, sensitive to derivative couplings)
- Background processes unaffected

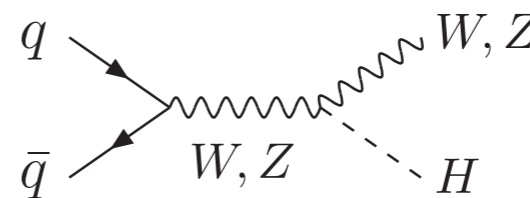
$\Downarrow$   $\Downarrow$   $\Downarrow$   
 simple rescaling of SM searches



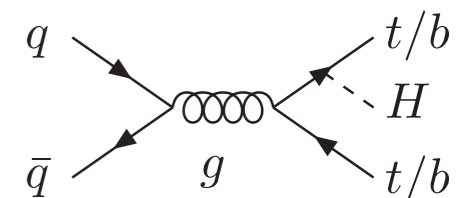
$c^2$



$a^2$



$a^2$



$c^2$

$$\frac{\sigma_{NLO}^{SM}}{\sigma_{NLO}}$$



The QCD NLO rescale trivially in the flavor universal limit.  
Not the EW NLO



$$\Gamma(H \rightarrow f\bar{f}) = c^2 \Gamma^{SM}(H \rightarrow f\bar{f}),$$

$$\Gamma(H \rightarrow VV) = a^2 \Gamma^{SM}(H \rightarrow VV),$$

$$\Gamma(H \rightarrow gg) = c^2 \Gamma^{SM}(H \rightarrow gg),$$

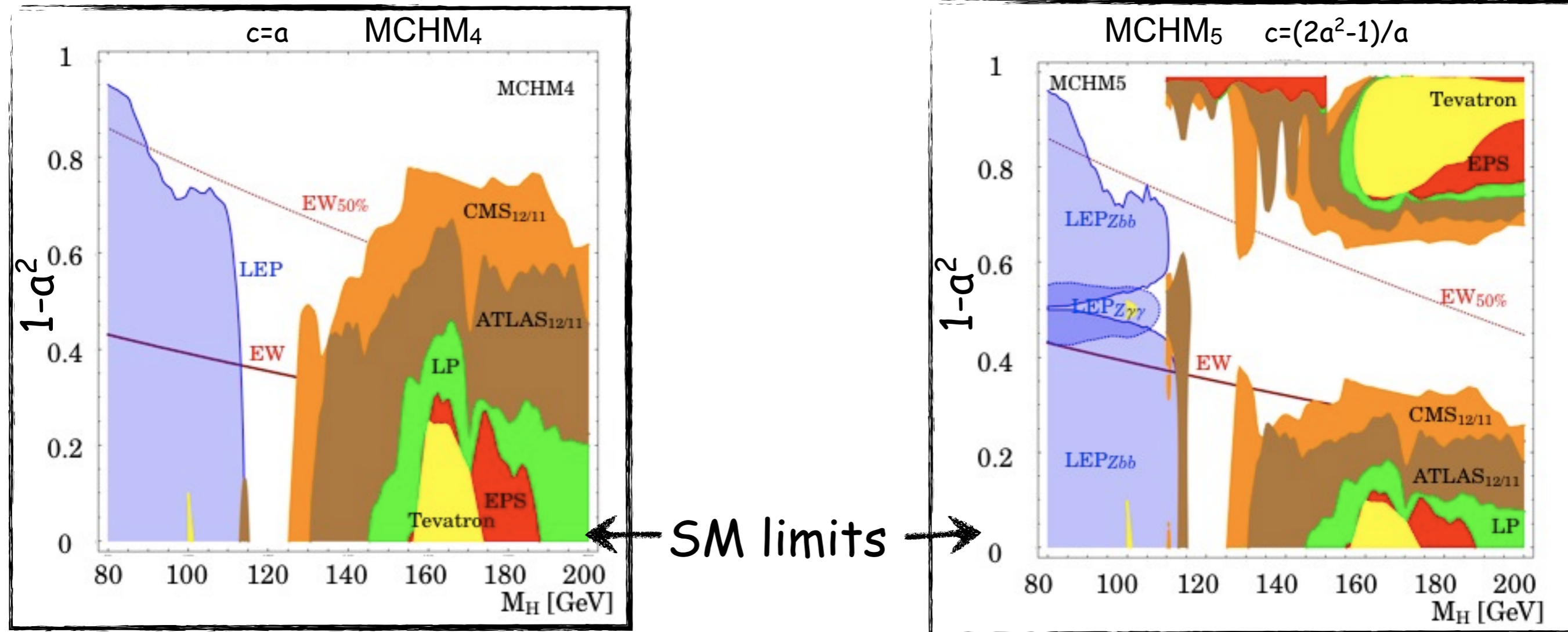
$$\Gamma(H \rightarrow \gamma\gamma) = \frac{(cI_\gamma + aJ_\gamma)^2}{(I_\gamma + J_\gamma)^2} \Gamma^{SM}(H \rightarrow \gamma\gamma),$$

$$\simeq (1.26a - 0.26c)^2 \text{ for } m_h = 125 \text{ GeV}$$

# Deformation of the SM Higgs: current constraints

the SM exclusion bounds are easily rescaled in the  $(m_H, a)$  plane

Espinosa, Grojean, Muehlleitner '11



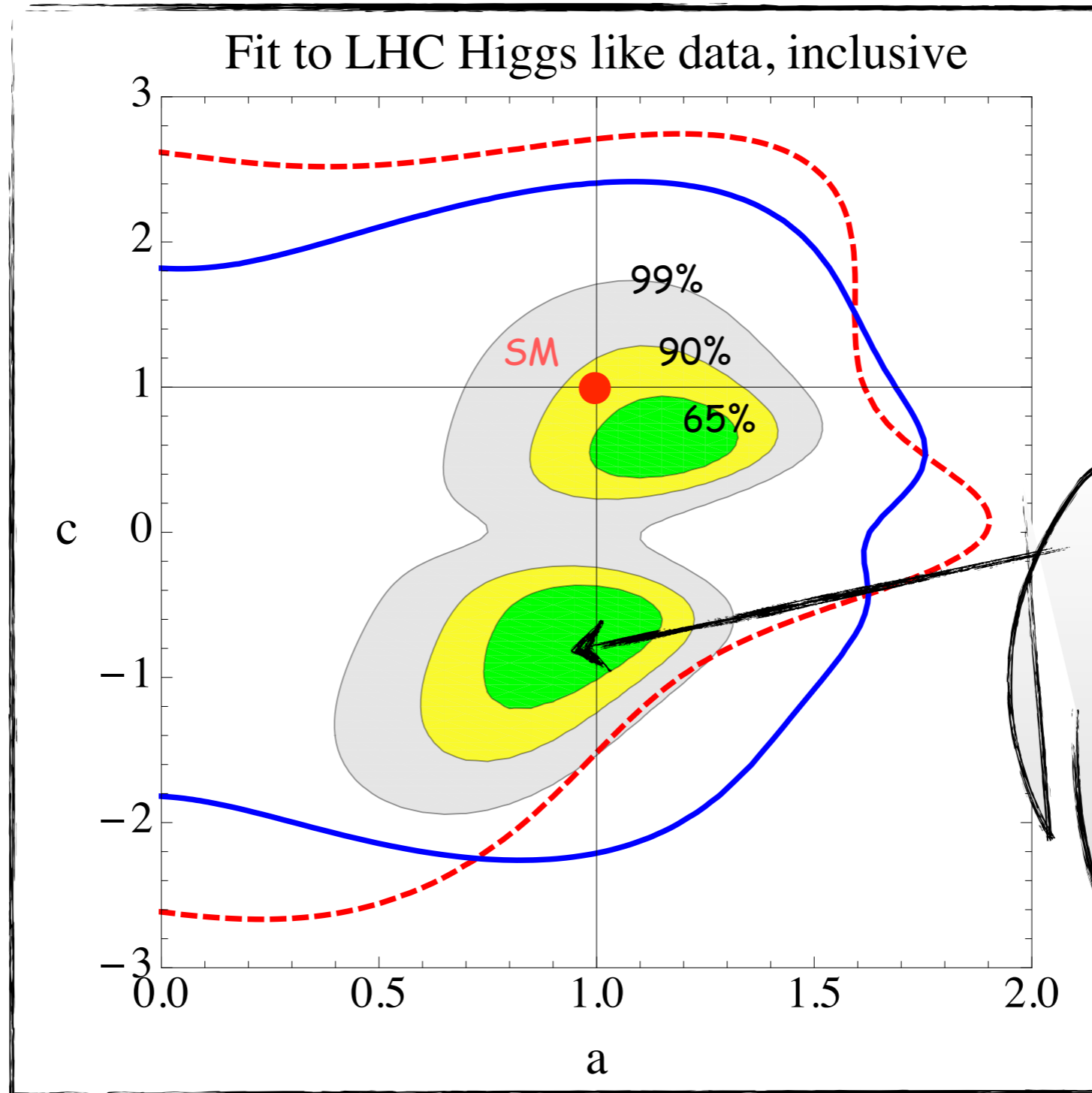
the LHC can do much more than simply excluding the SM Higgs

for similar analysis, see also [Azatov, Contino, Galloway '12](#)



# Model independent $\chi^2$ fit to LHC excess @ 125

Espinosa, Grojean, Muhlleitner, Trott '12



note: a fermiophobic Higgs is disfavored by data (mostly VBF channels)

"disfermiophilia"

the current data prefers "negative" coupling to fermions  
 $\approx$   
 positive interference between top and W in  $\gamma\gamma$  channel

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 Atlas 95%CL exclusion

—  
 CMS 95%CL exclusion

SM 82%CL  
 away from  
 best fit point

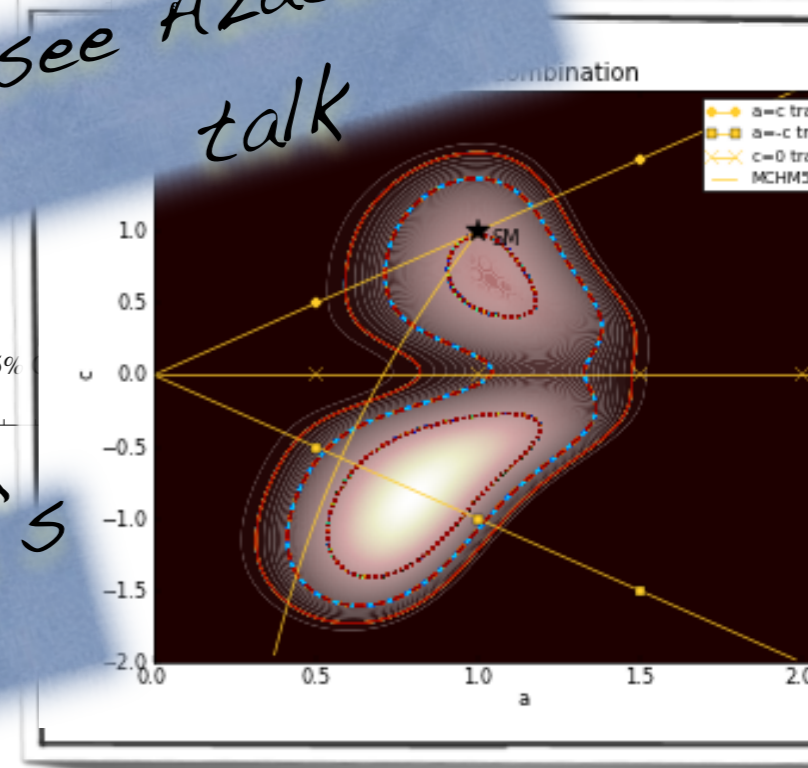
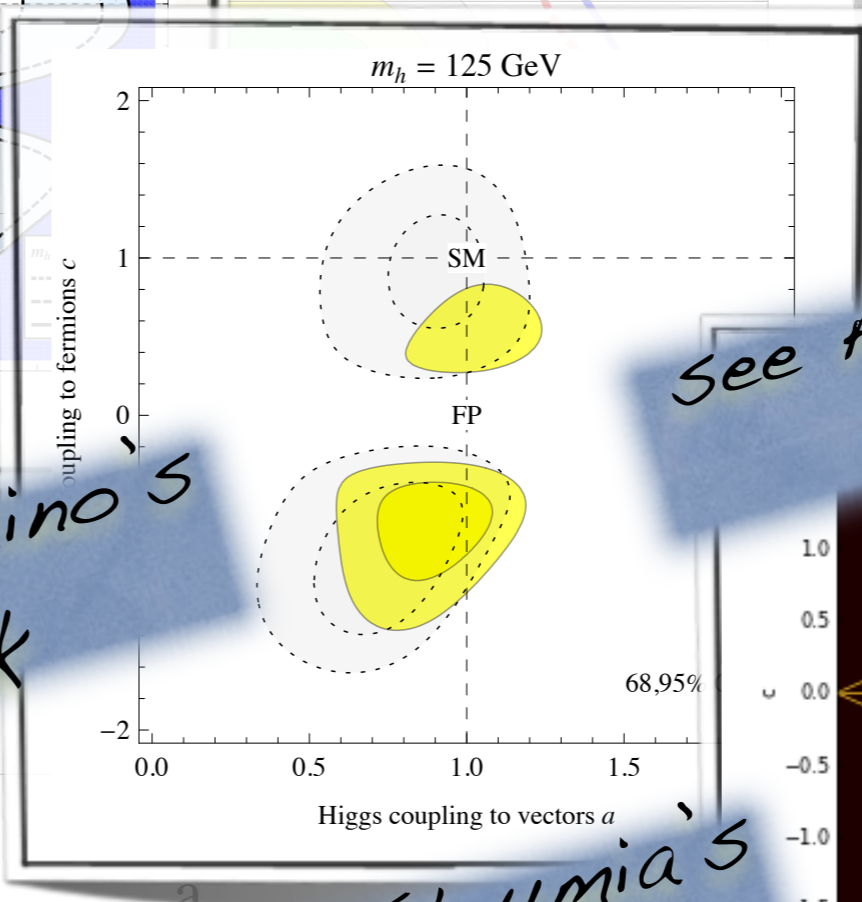
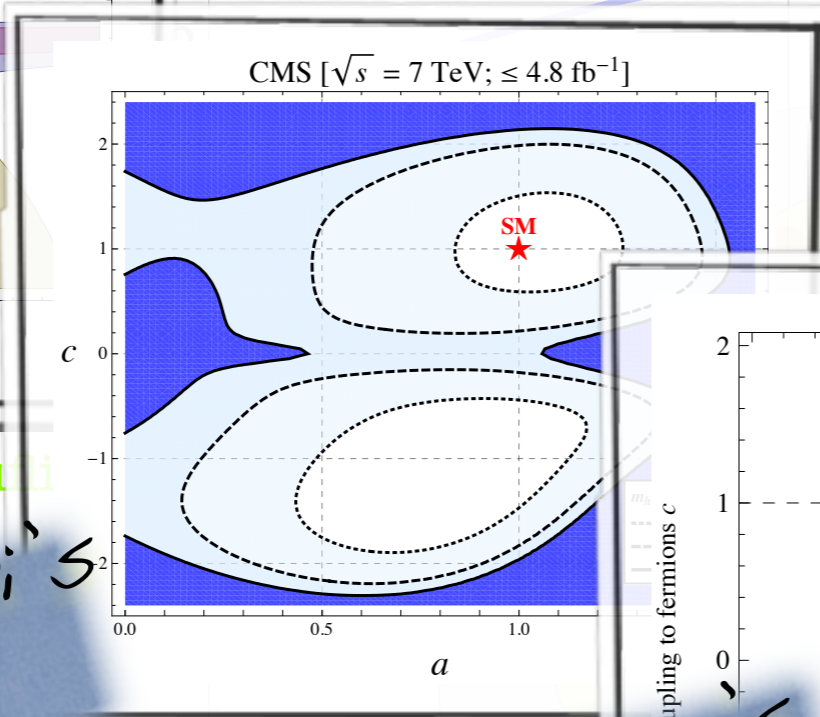
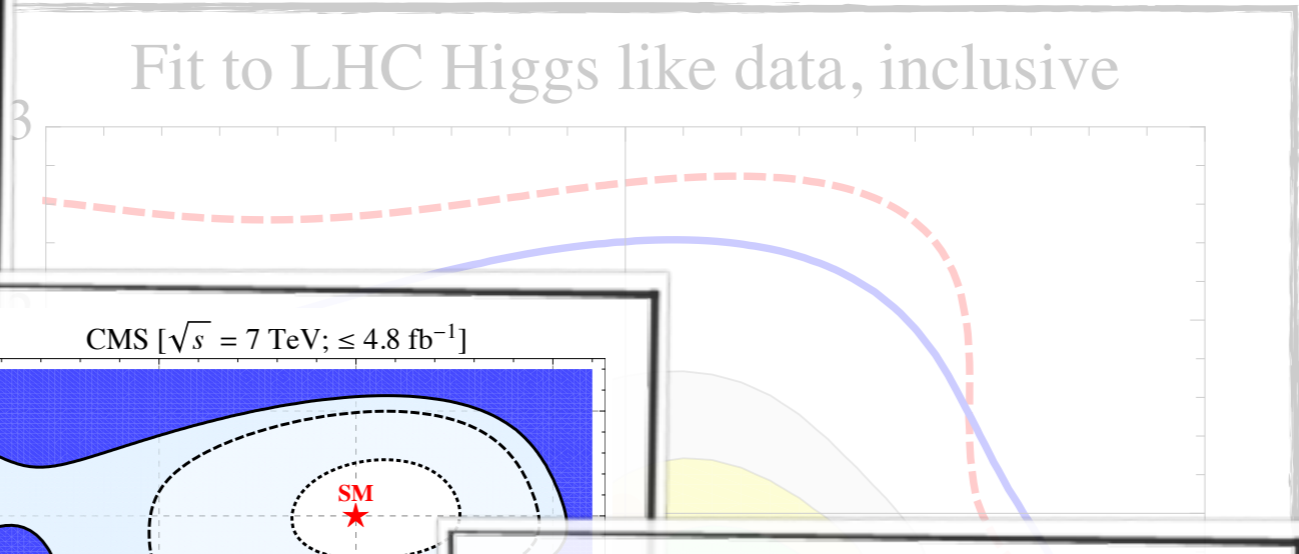
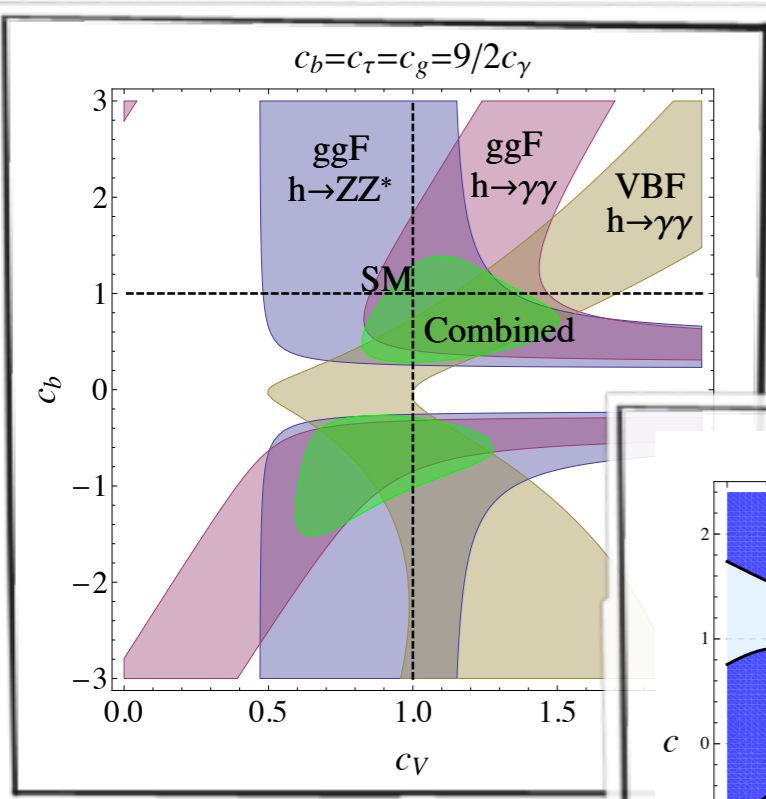
Two minima:

$(a,c)=(1.13,0.58)$   
 $\chi^2=2.86$

$(a,c)=(0.96,-0.64)$   
 $\chi^2=1.96$

# Model independent $\chi^2$ fit to LHC excess @ 125

Espinosa, Grojean, Muhlleitner, Trott '12



Carni, Falkowski, Kuziakov, Volansky '12

see Falkowski's talk

Azatov, Contino '12

see Contino's talk

Giardin '12

see Strumia's talk

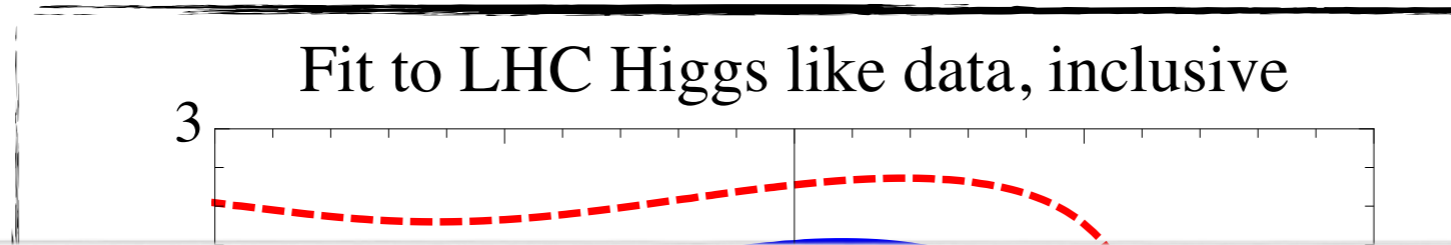
see Azatov's talk

Ellis, You '12

Zurich, 22<sup>th</sup> May 2012

# Model independent $\chi^2$ fit to LHC excess @ 125

Espinosa, Grojean, Muhlleitner, Trott '12



note: a fermiophobic Higgs is disfavored by

many issues to "validate" this kind of fits:

1. exact likelihood (departure from Gaussians...)
2. correlations among channels
3. combination of rescaled channels
4. cut efficiencies
5. ...
6. missing parameters:  $c_t \neq c_b$ ,  $a_W \neq a_Z$ ...

(a,c)=(0.96,-0.64)

$\chi^2=1.96$

for similar analyses, see also

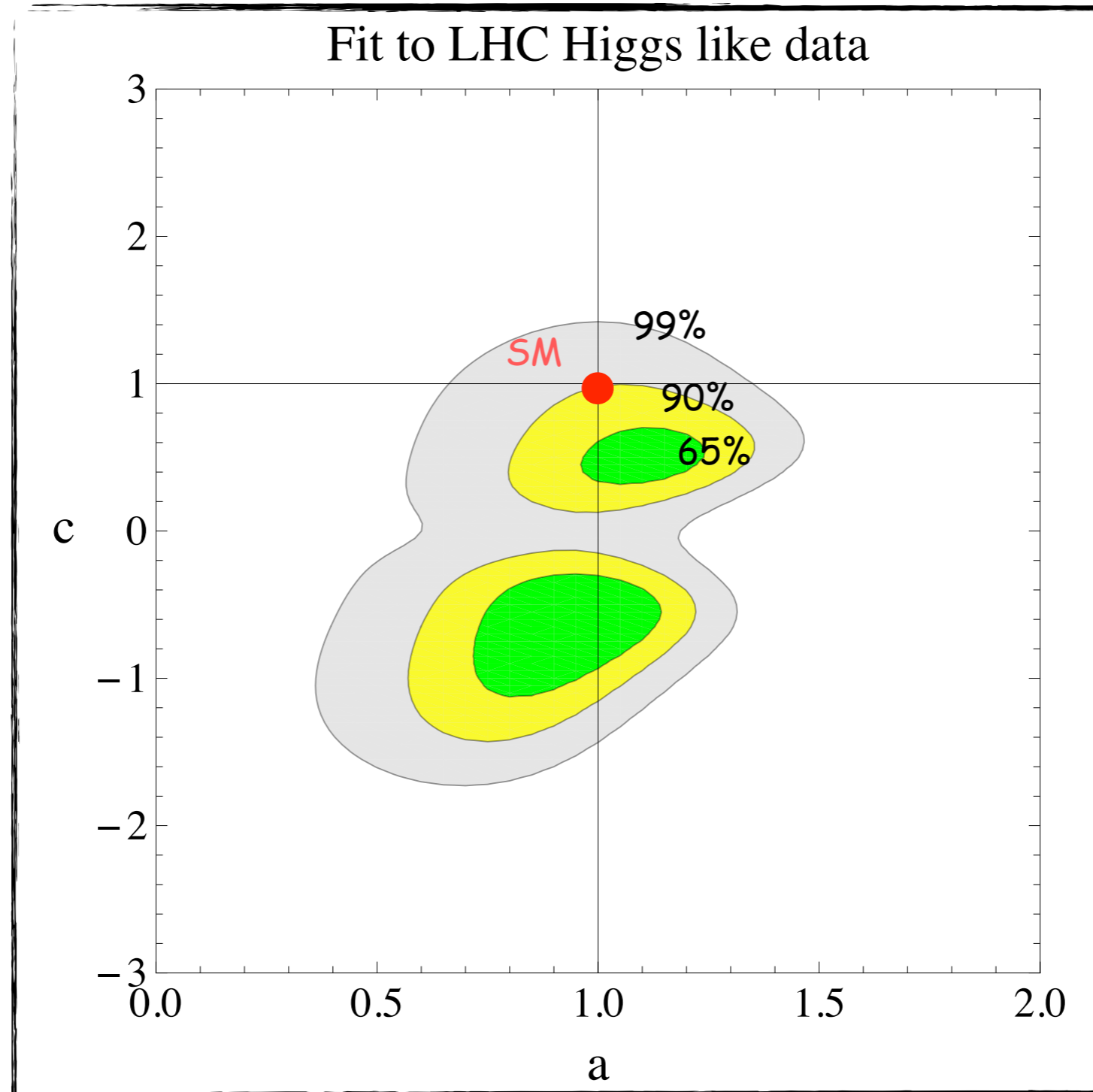
Azatov, Contino, Galloway '12

Carni, Falkowski, Kuflik, Volansky '12



# Model independent $\chi^2$ fit to (Moriond) LHC data

Espinosa, Grojean, Muhlleitner, Trott '12



note: a fermiophobic Higgs is disfavored by data (mostly VBF channels) at 97%CL

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Atlas 95%CL exclusion

—

CMS 95%CL exclusion

SM 88%CL  
away from  
best fit point  
( $\sim 2\sigma$ )

Two minima:

$(a,c)=(1.18,0.55)$   
 $\chi^2=7.5$

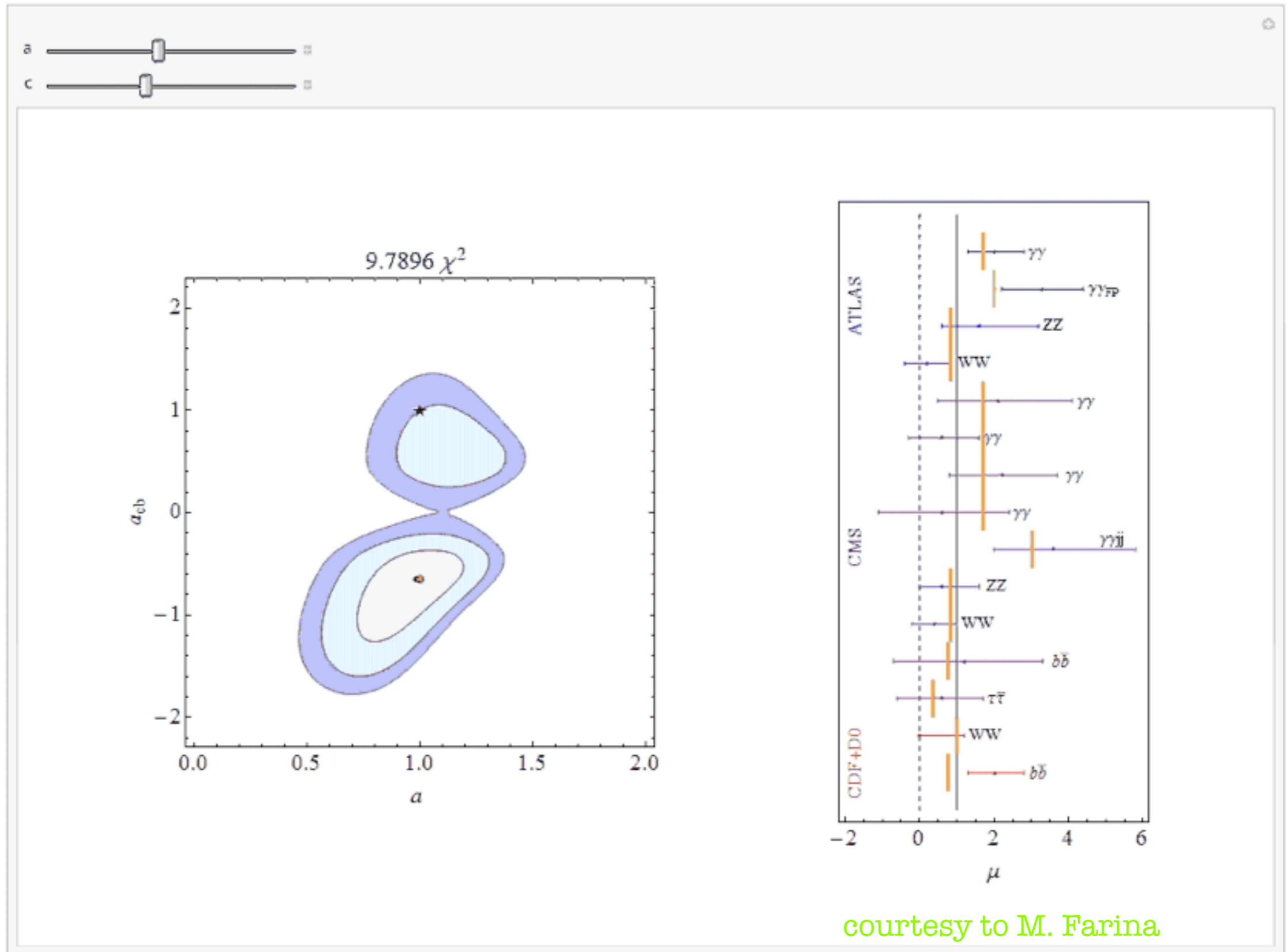
$(a,c)=(0.99,-0.64)$   
 $\chi^2=6.3$

for similar analyses, see also

Azatov, Contino, Galloway '12

Carni, Falkowski, Kuflik, Volansky '12

# Model independent $\chi^2$ fit to LHC data





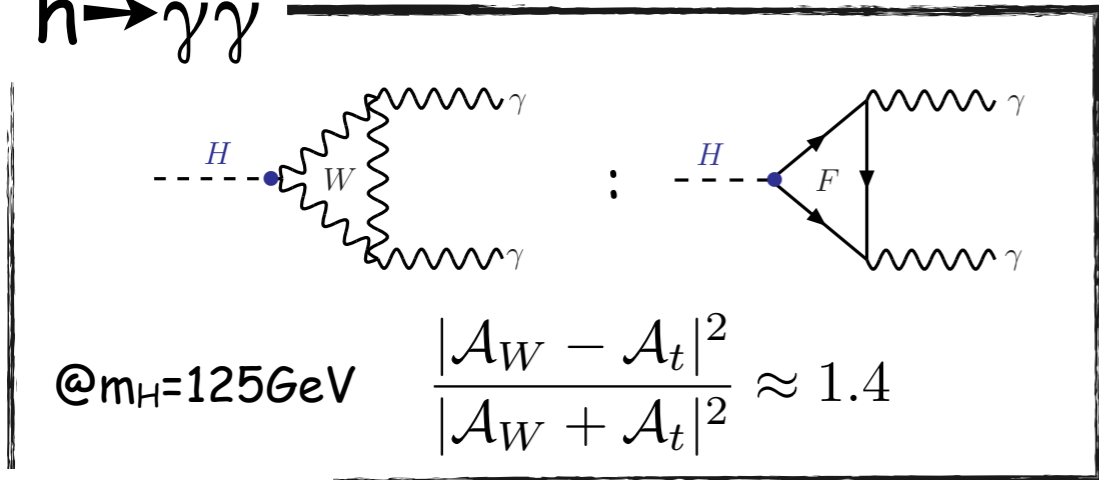
# Fermiophilia or Disfermiophilia?

Farina, Grojean, Maltoni,  
Salvioni, Thamm 'in progress

difficult!

difference is physically relevant only in the presence of  
strong interference with single  $h\psi\psi$  coupling

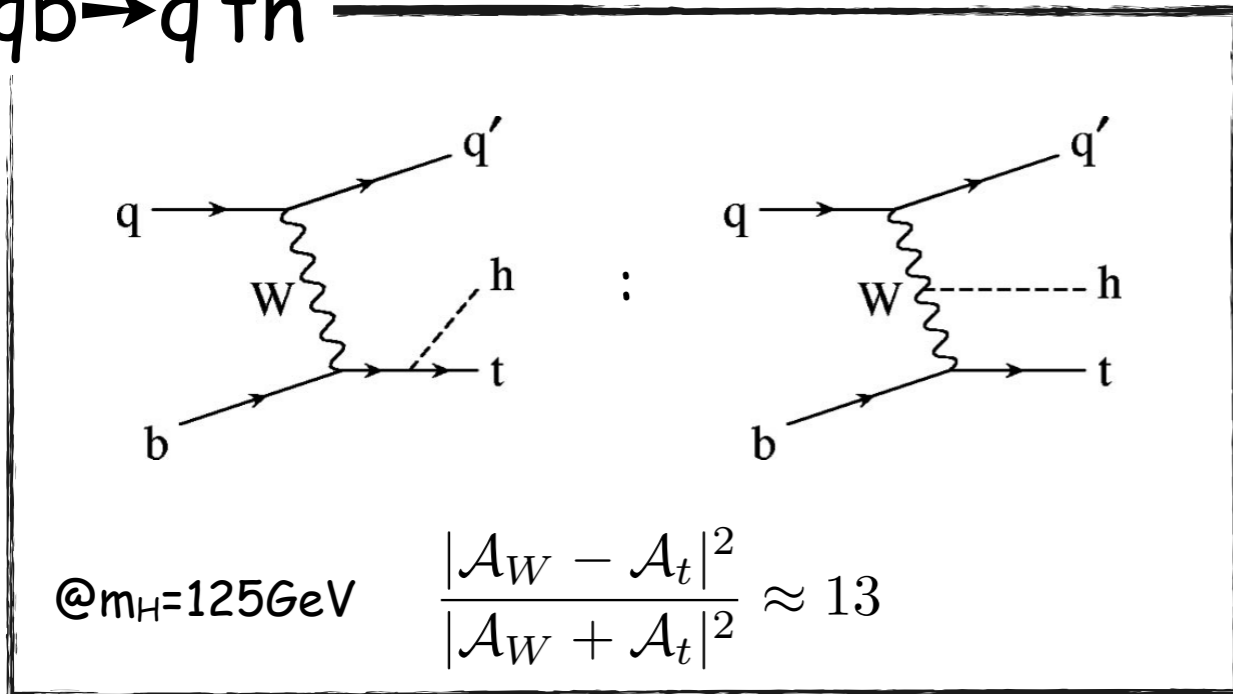
$h \rightarrow \gamma\gamma$



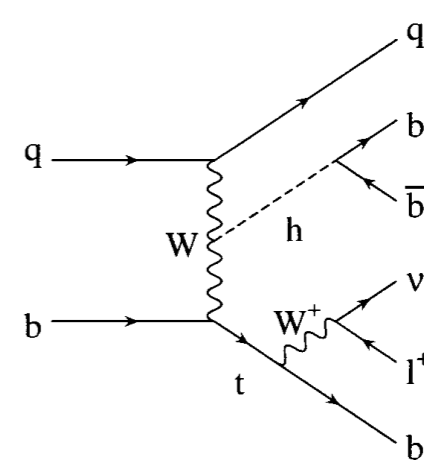
rare decay and one needs  
some largish luminosity to  
be sensitive to the sign of  $c$

see Azatov's  
talk

$qb \rightarrow q'th$



look at final state:  $3b + 1 \text{ fwd jet} + l^\pm + p^T$ .



Maltoni, Stelzer, Willenbrock '01

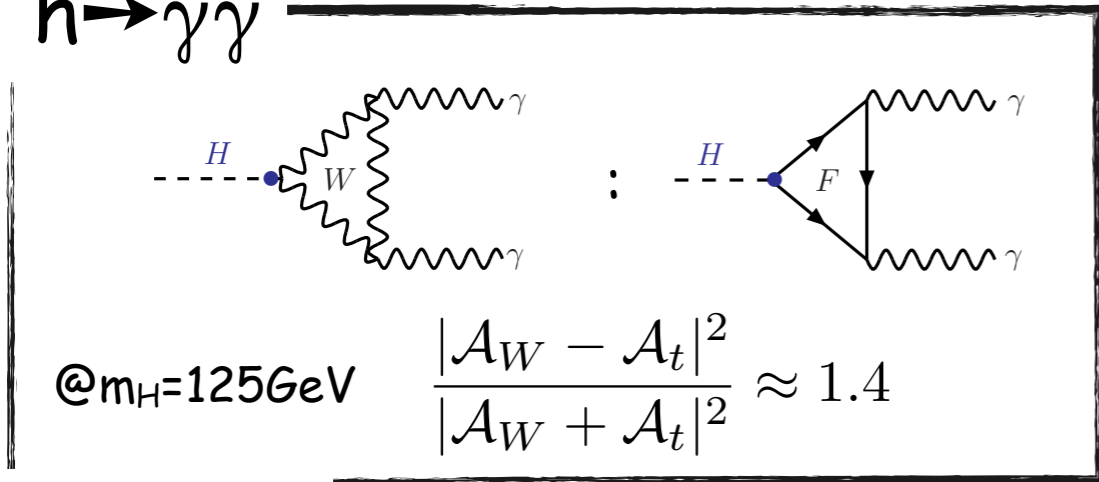
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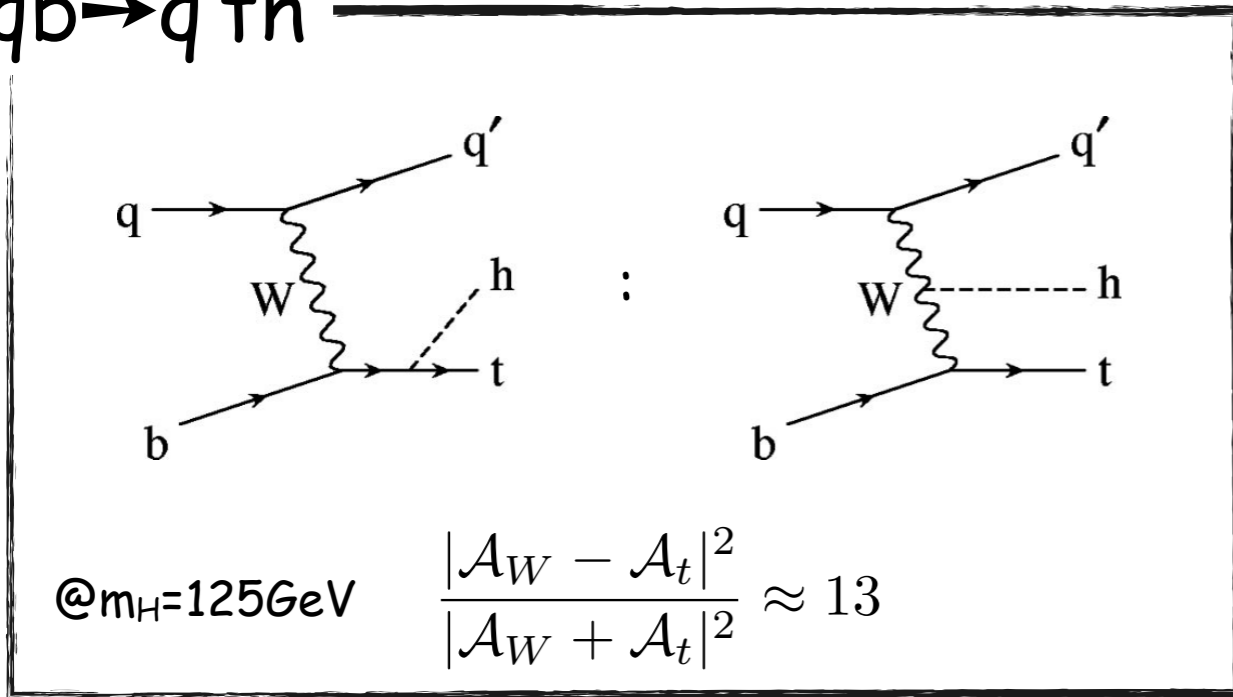
$h \rightarrow \gamma\gamma$



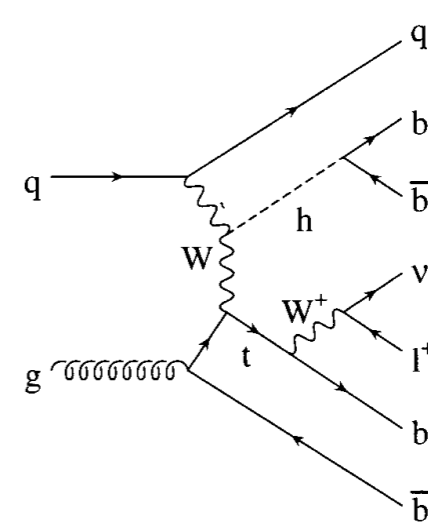
rare decay and one needs  
some largish luminosity to  
be sensitive to the sign of  $c$

see Azatov's  
talk

$qb \rightarrow q'th$



look at final state:  $4b + 1 \text{ fwd jet} + l^\pm + p^T$ .

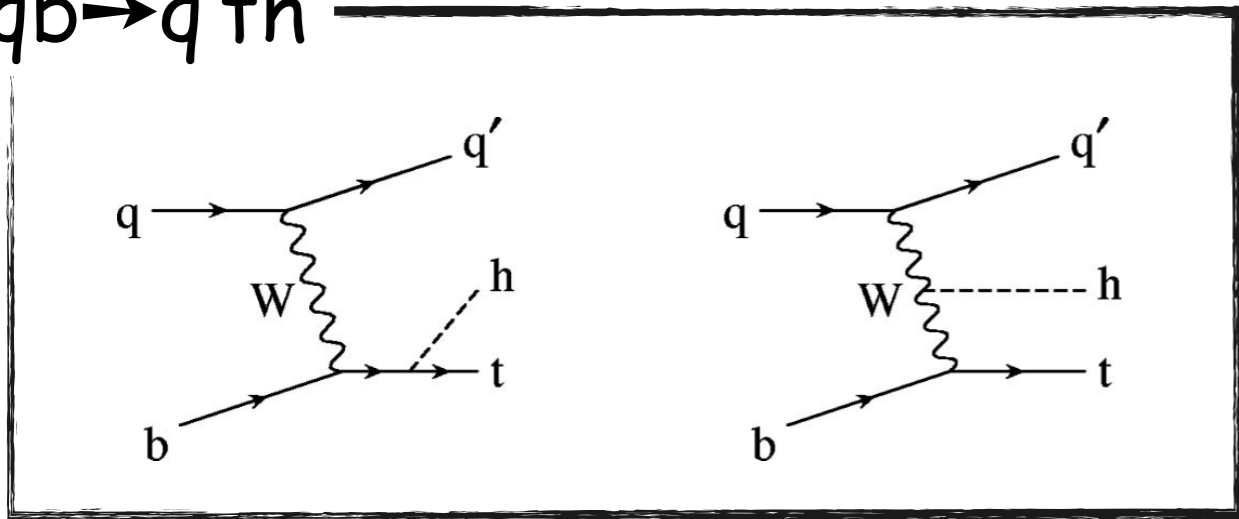


Maltoni, Stelzer, Willenbrock '01

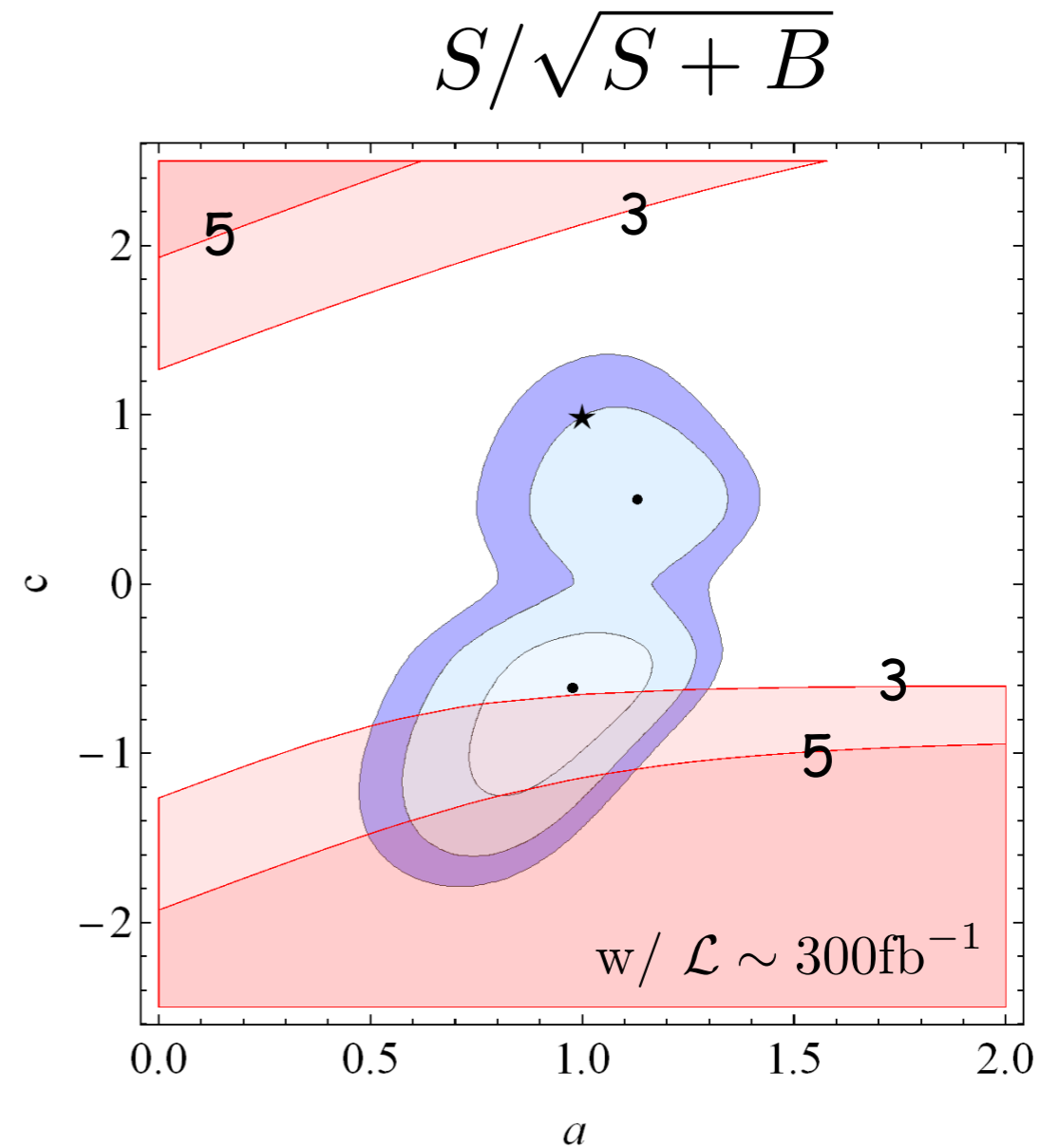
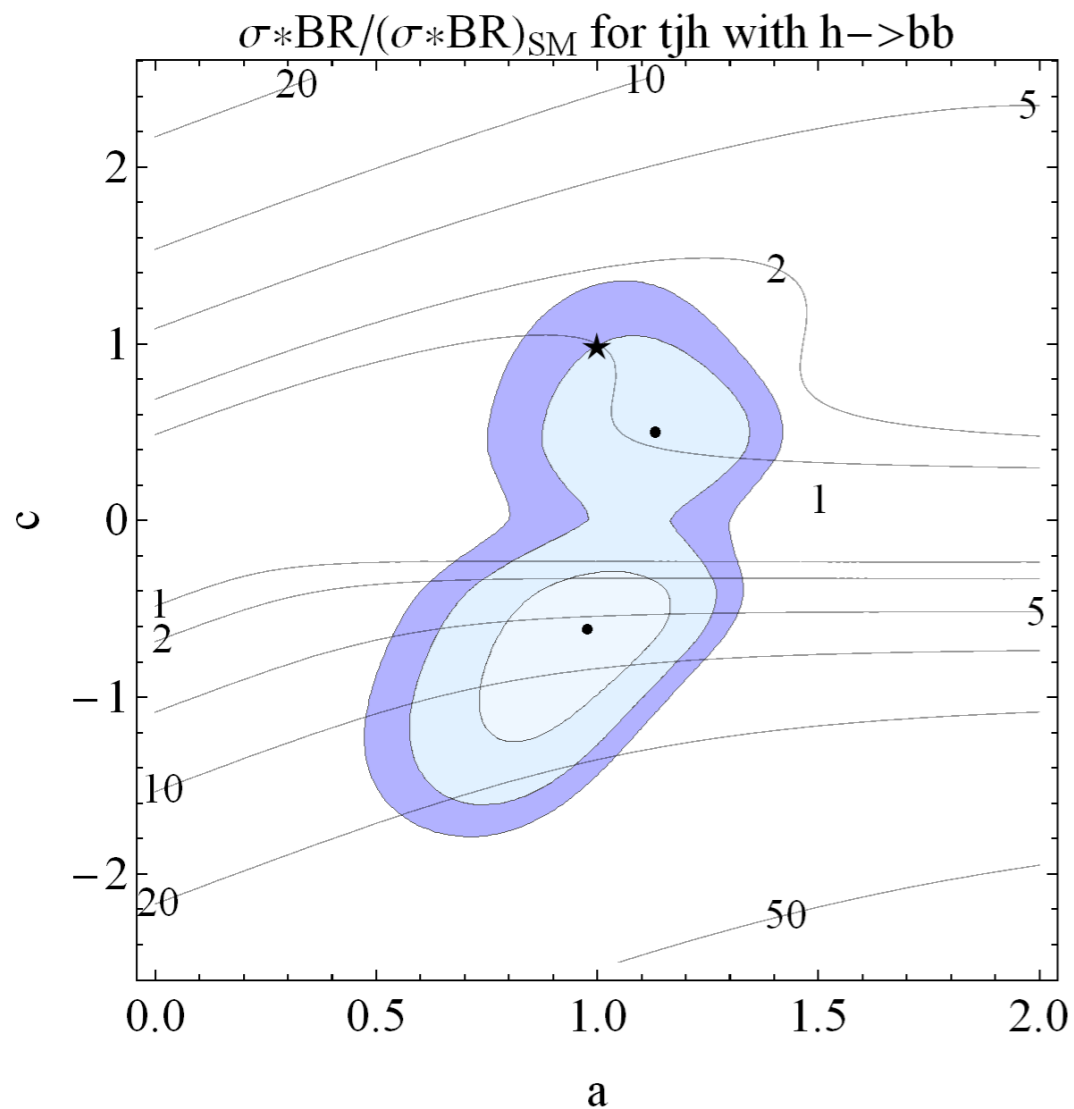
# Fermiophilia or Disfermiophilia?

Farina, Grojean, Maltoni,  
Salvioni, Thamm 'in progress

$qb \rightarrow q'th$

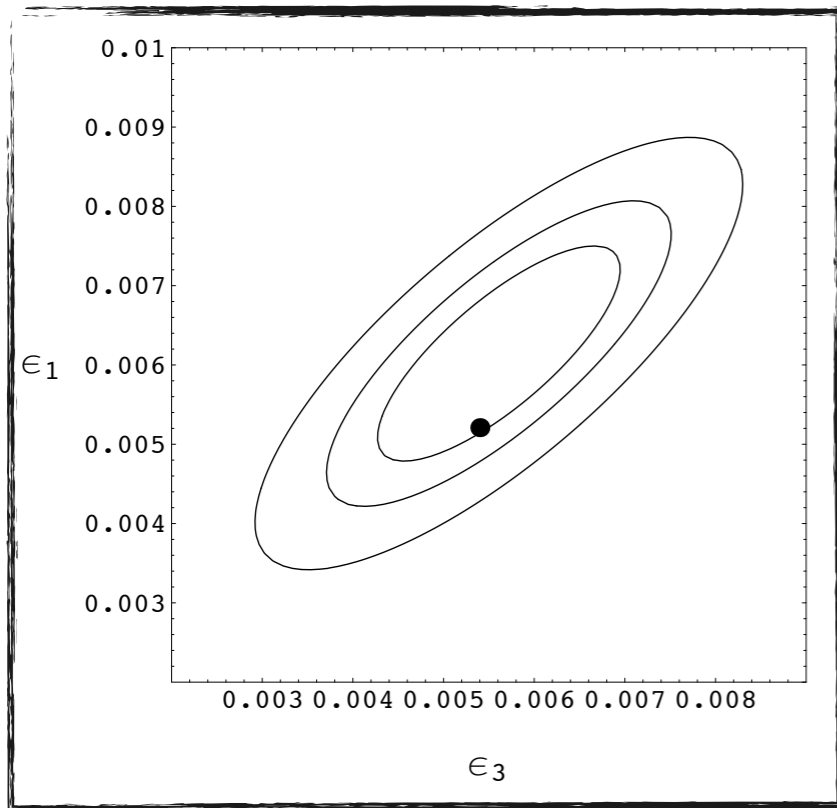


*the sign ambiguity will remain  
with us for a long time!*





# A tension between LHC and EW data?



EW fit strongly suggests custodial symmetry

$$\Sigma = e^{i\sigma^a \pi^a / v}$$

Goldstone of  
 $SU(2)_L \times SU(2)_R / SU(2)_V$

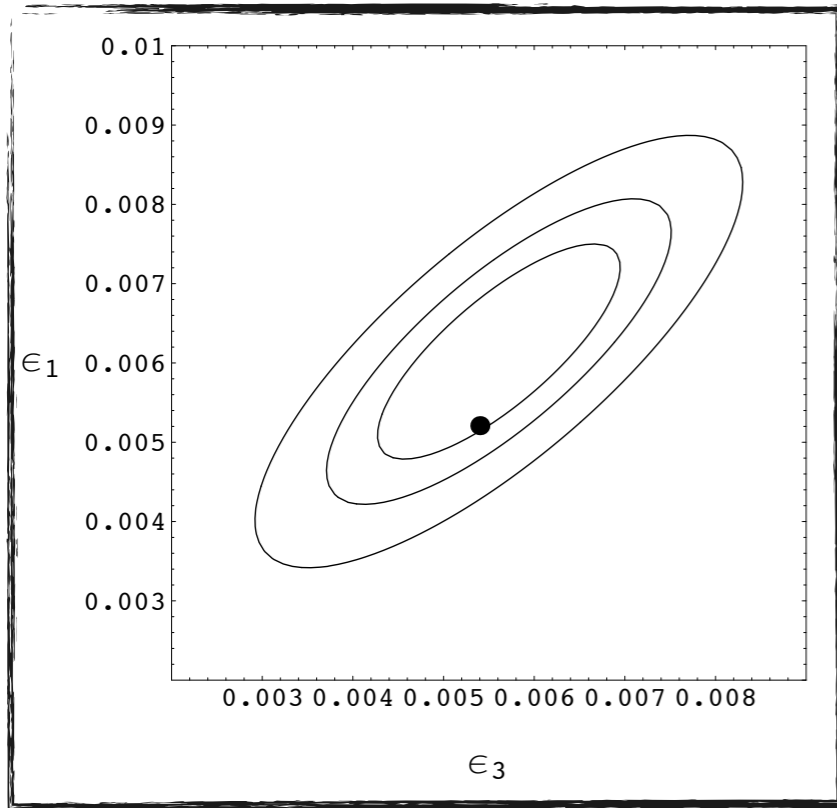
$$\frac{v^2}{4} \text{Tr} (D_\mu \Sigma^\dagger D^\mu \Sigma) \Rightarrow \rho = 1 \quad \text{ie} \quad \epsilon_1 = \hat{T} = 0 \quad \checkmark$$

$$\text{also} \Rightarrow \mu_{ZZ} = \mu_{WW}$$

$$\left( \mu_i = \frac{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i)}{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i) |_{\text{SM}}} \right)$$

$$\frac{v^2}{8} \text{Tr}^2 (\Sigma^\dagger D_\mu \Sigma \sigma^3) \Rightarrow \rho = 2 \quad \text{ie} \quad \epsilon_1 = \hat{T} = 1 \quad \text{strongly disfavored}$$

# A tension between LHC and EW data?



EW fit strongly suggests custodial symmetry

$$\Sigma = e^{i\sigma^a \pi^a / v}$$

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$$\text{also } \Rightarrow \mu_{ZZ} = \mu_{WW} \quad \times$$

but

Channel [Exp]	$\mu_{119.5} (\mu_{119.5}^L)$	$\mu_{124} (\mu_{124}^L)$	$\mu_{125} (\mu_{125}^L)$
$pp \rightarrow Z Z^* \rightarrow \ell^+ \ell^- \ell^+ \ell^-$ [ATLAS]	$-0.5^{+0.5??} (5.1)$	$1.6^{+1.4}_{-0.8} (4.7)$	$1.4^{+1.3}_{-0.8} (4.1)$
$pp \rightarrow W W^* \rightarrow \ell^+ \nu \ell^- \bar{\nu}$ [ATLAS]	$0.0^{+1.2}_{-1.3} (2.4)$	$0.1^{+0.7}_{-0.7} (1.6)$	$0.1^{+0.7}_{-0.6} (1.4)$

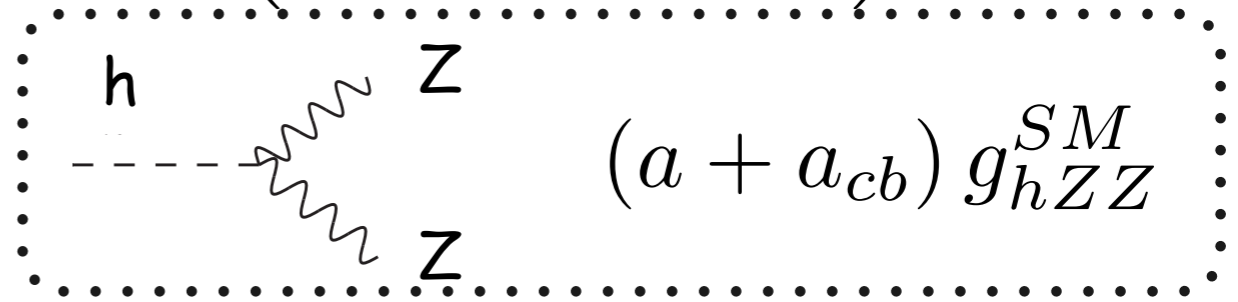
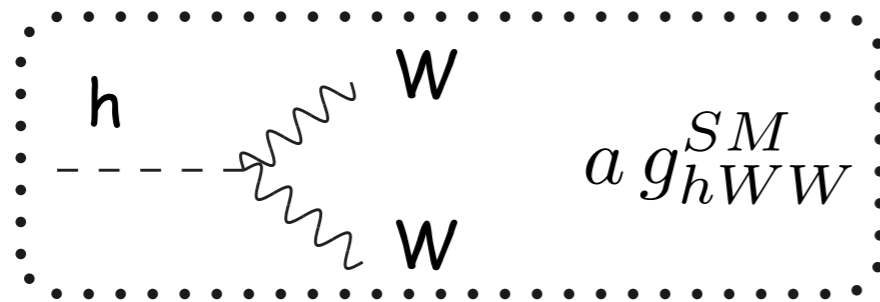
1. has LHC identified a violation of the custodial symmetry?
2. if yes, how to reconcile LHC data with EW data?

$$\frac{v^2}{8} \text{Tr}^2 (\Sigma^\dagger D_\mu \Sigma \sigma^3) \Rightarrow \rho = 2 \text{ ie } \epsilon_1 = \hat{T} = 1 \quad \text{strongly disfavored}$$

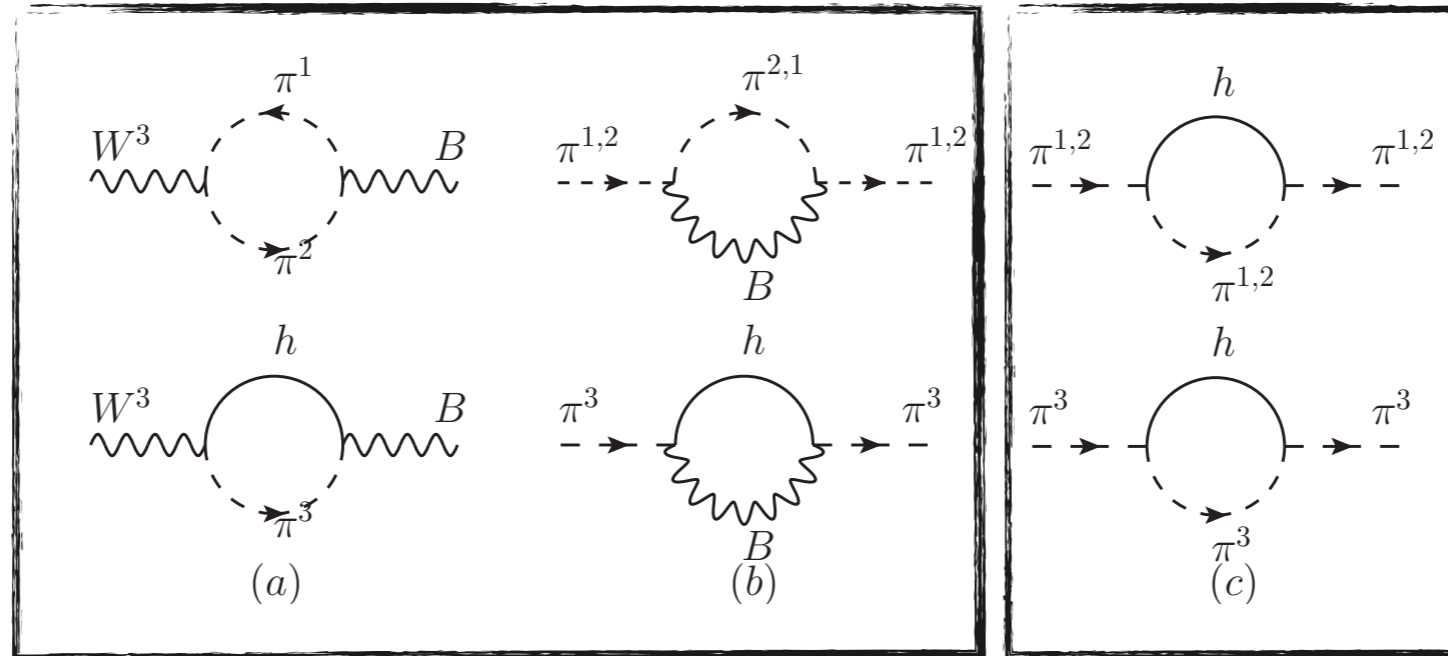
# DisZphilia or how to live with custodial breaking

Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left( \text{Tr} [\Sigma^\dagger D_\mu \Sigma \sigma^3] \right)^2 \left( t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$



EWPT



Log( $\Lambda$ ) UV sensitivity

$\Lambda^2$  UV sensitivity

$$\Delta\epsilon_1 = -\frac{3}{16\pi} \frac{\alpha(m_Z)}{\cos^2 \theta_W} \left[ 1 - (a + a_{cb})^2 + \left( \frac{g}{g'} \right)^2 (a^2 - (a + a_{cb})^2) \right] \log \left( \frac{\Lambda^2}{m_h^2} \right)$$

$$\Delta\epsilon_3 = +\frac{1}{48\pi} \frac{\alpha(m_Z)}{\sin^2 \theta_W} \left[ 1 - (a + a_{cb})^2 \right] \log \left( \frac{\Lambda^2}{m_h^2} \right)$$

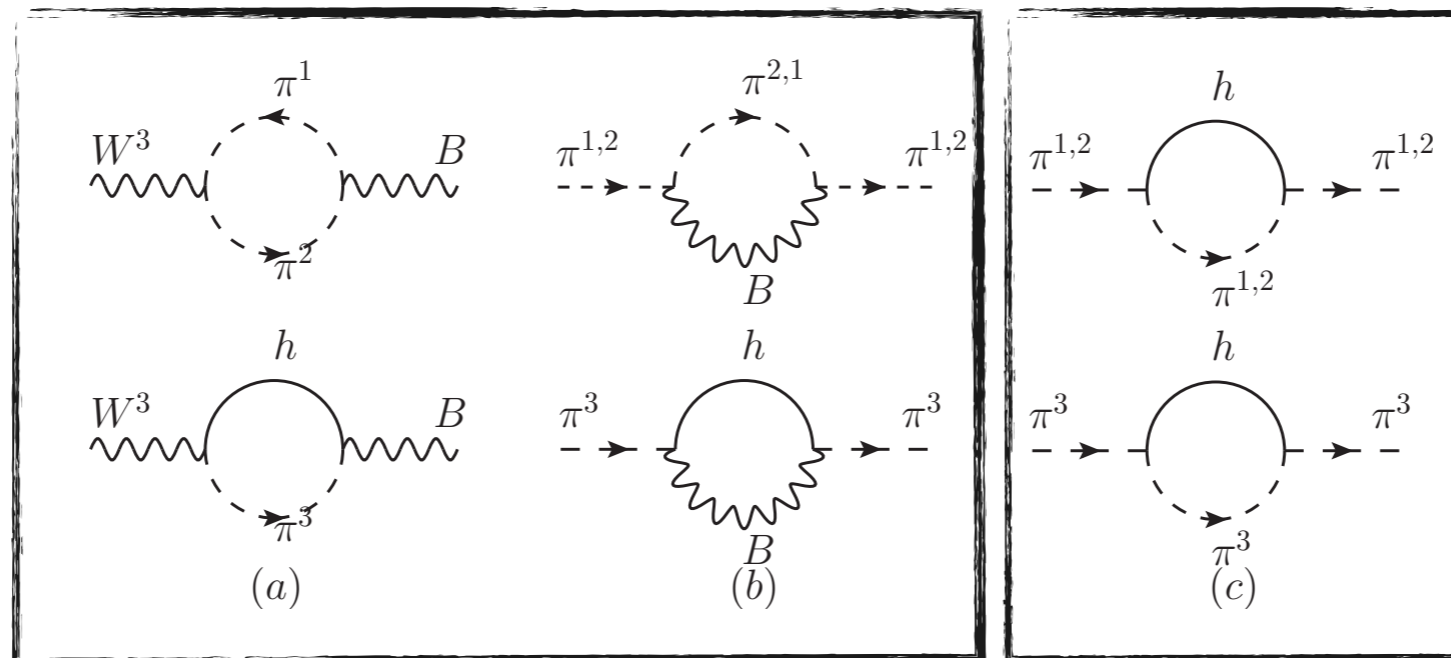
$$\Delta\epsilon_1 = \left( (a + a_{cb})^2 - a^2 \right) \frac{\Lambda^2}{16\pi^2 v^2}$$



# DisZphilia or how to live with custodial breaking

Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left( \text{Tr} \left[ \Sigma^\dagger D_\mu \Sigma \sigma^3 \right] \right)^2 \left( t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$



Log( $\Lambda$ ) UV sensitivity

$\Lambda^2$  UV sensitivity

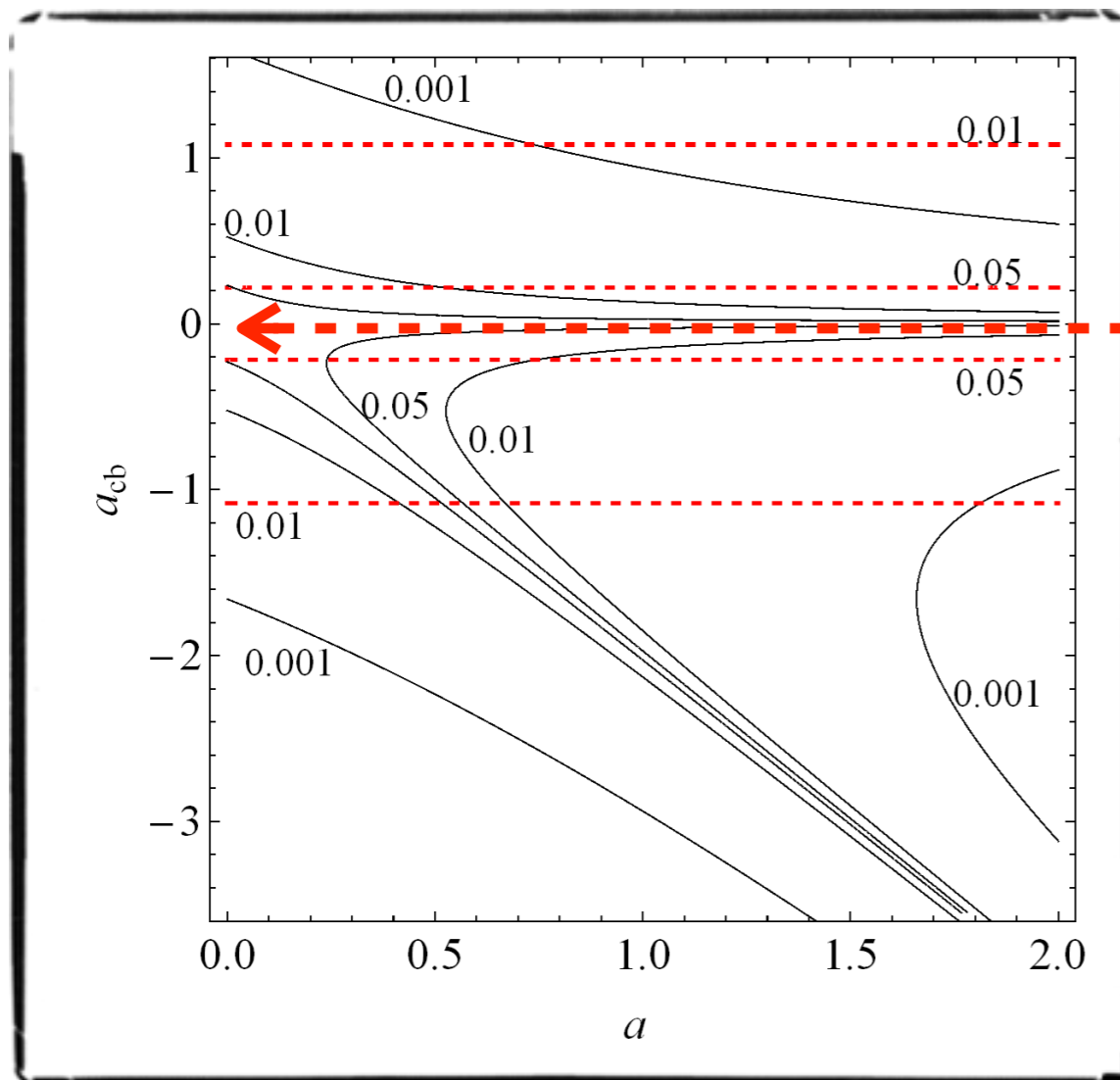
EWPT highly model-dependent  
 tuning between tree-level and loop contributions?  
 new light states?

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Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left( \text{Tr} \left[ \Sigma^\dagger D_\mu \Sigma \sigma^3 \right] \right)^2 \left( t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$

EWPT highly model-dependent  
tuning between tree-level and loop contributions?  
new light states?



*Amount of fine-tuning in EWPT*

← ———— → custodial invariance

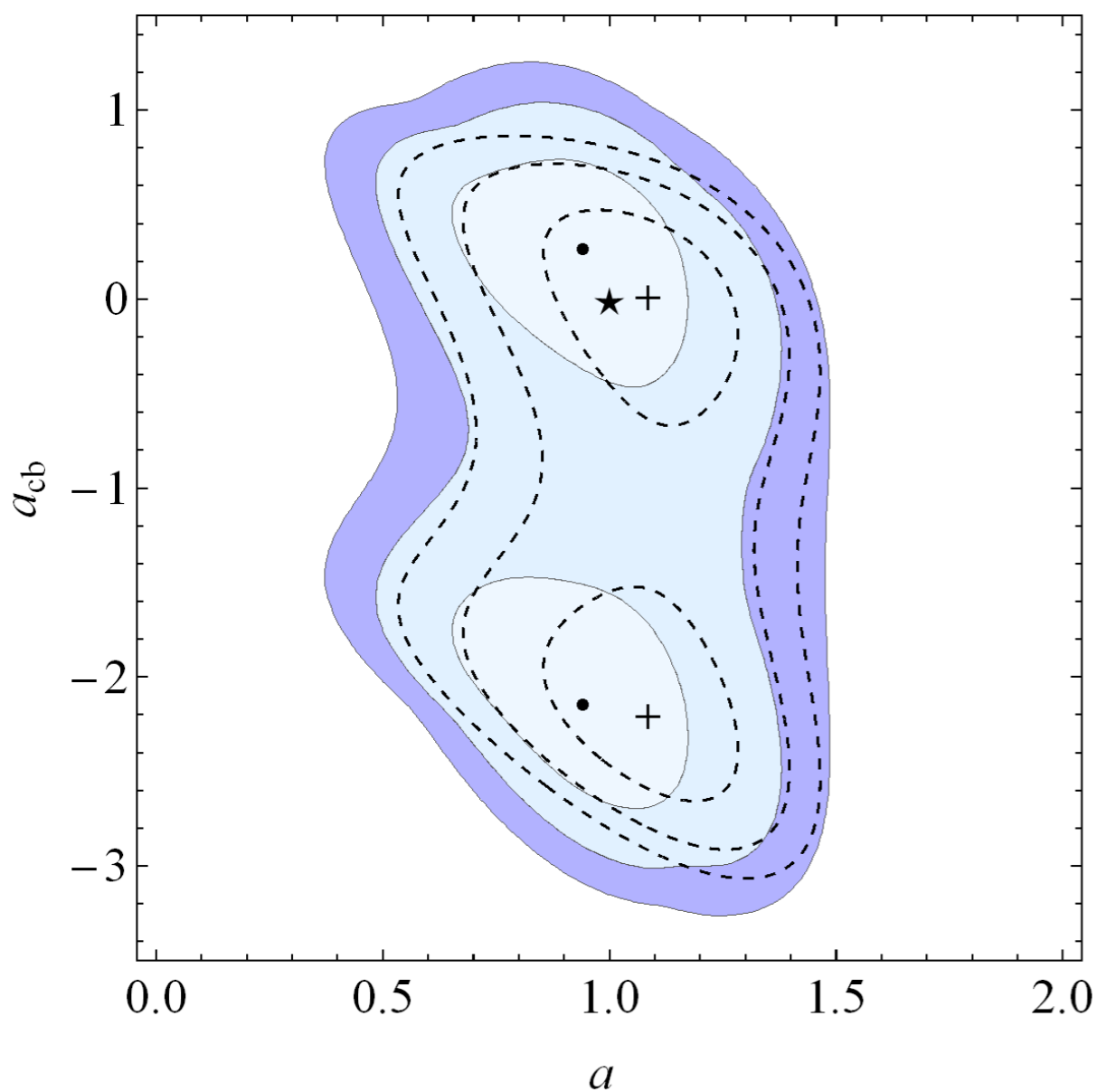
$\Lambda^2$  UV sensitivity  
ie  
could be as bad as the hierarchy problem

# DisZphilia or how to live with custodial breaking

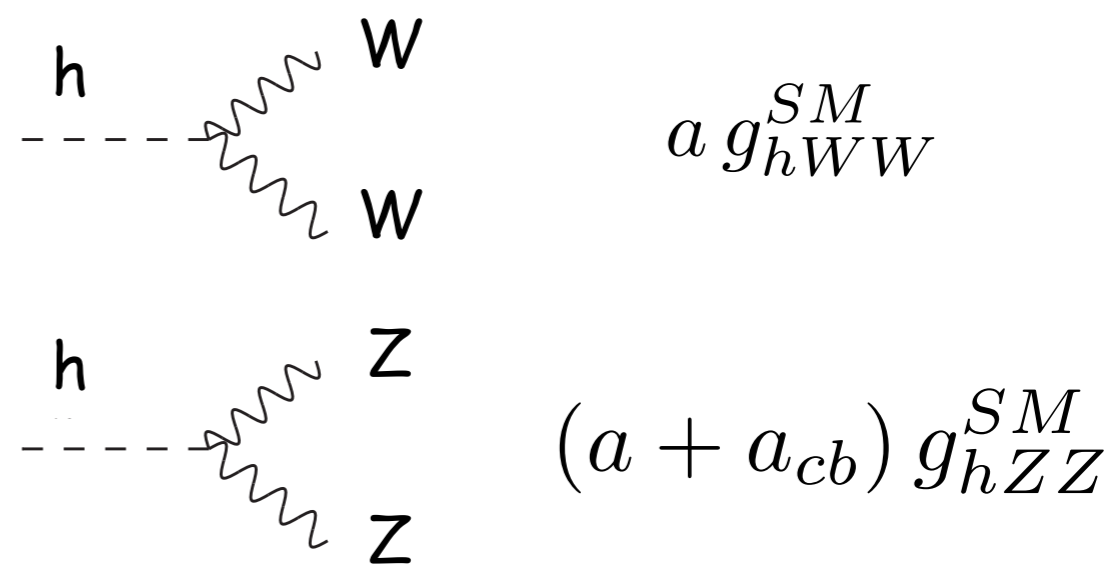
Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left( \text{Tr} [\Sigma^\dagger D_\mu \Sigma \sigma^3] \right)^2 \left( t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$

Fit to LHC data



-----  $c = 1$   
 ————— marginalization over  $c$



LHC data are symmetric under

$$(a + a_{cb}) \leftrightarrow -(a + a_{cb})$$

i.e.

$$(a, a_{cb}) \leftrightarrow (a, -(2a + a_{cb}))$$

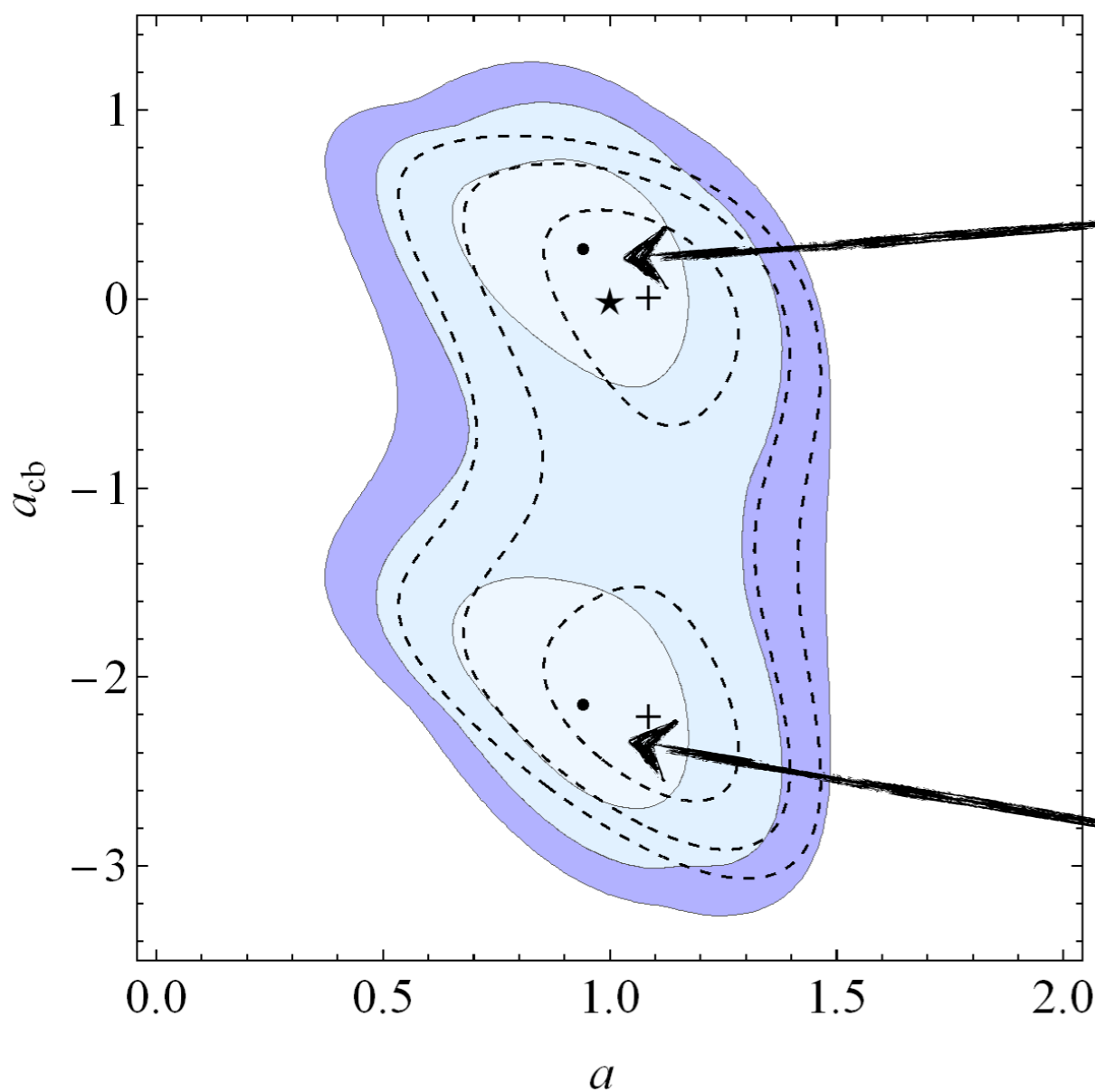


# DisZphilia or how to live with custodial breaking

Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left( \text{Tr} [\Sigma^\dagger D_\mu \Sigma \sigma^3] \right)^2 \left( t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$

Fit to LHC data



-----  $c = 1$   
 ————— marginalization over  $c$

$$2 \frac{h}{v} \left( m_W^2 W_\mu^+ W_\mu^- + \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right)$$

the two solutions can only be distinguished in the presence of interference with a single  $hZZ$  vertex

$$2 \frac{h}{v} \left( m_W^2 W_\mu^+ W_\mu^- - \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right)$$

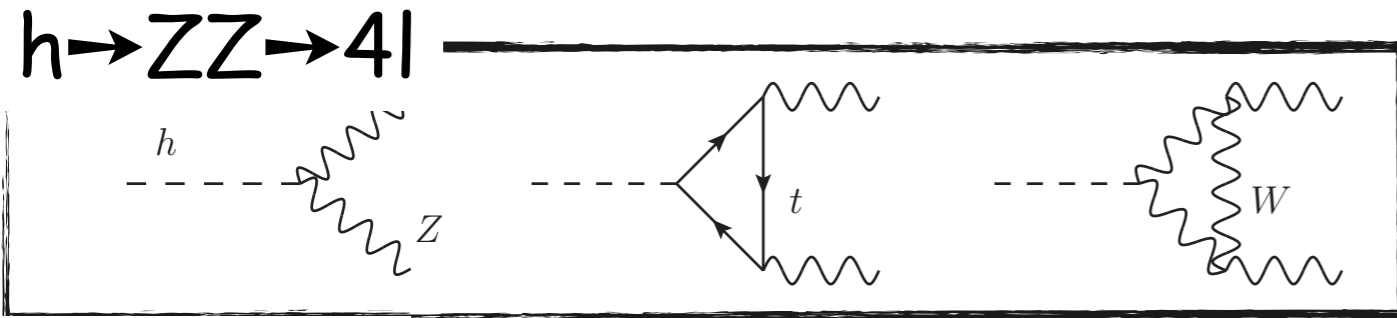
↑  
 "disZphilia"

# Zphilia or DisZphilia?

difficult!

Farina, Grojean, Salvioni '12

difference is physically relevant only in the presence of interference with single hZZ coupling

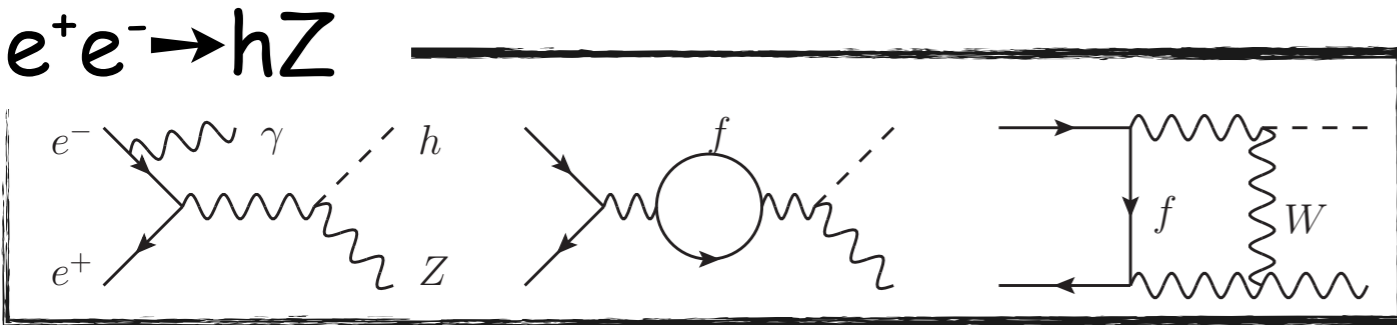


TH prediction

$$\Delta = \left| \frac{\Gamma_Z^+ - \Gamma_Z^-}{\Gamma_Z^+ + \Gamma_Z^-} \right| = \delta \approx 1\%$$

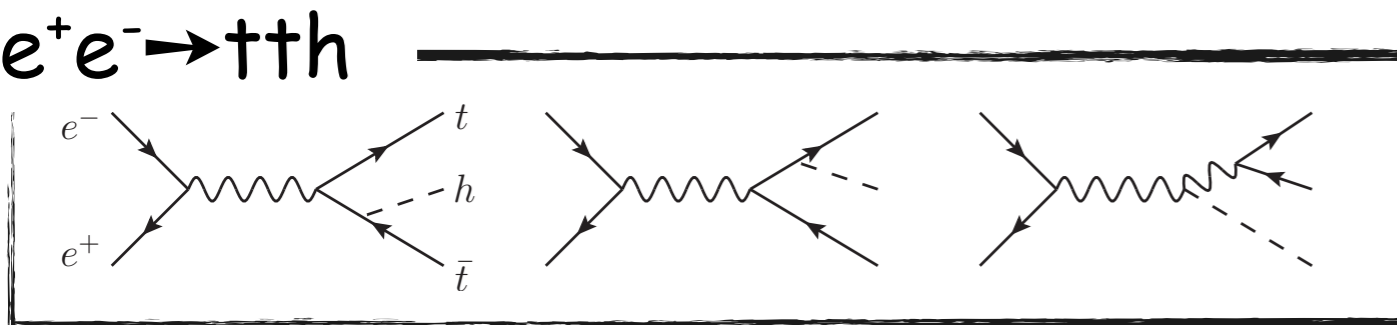
ILC ( $\sqrt{s}=800\text{GeV}$  and  $1\text{ab}^{-1}$ )

$\approx 1\%$



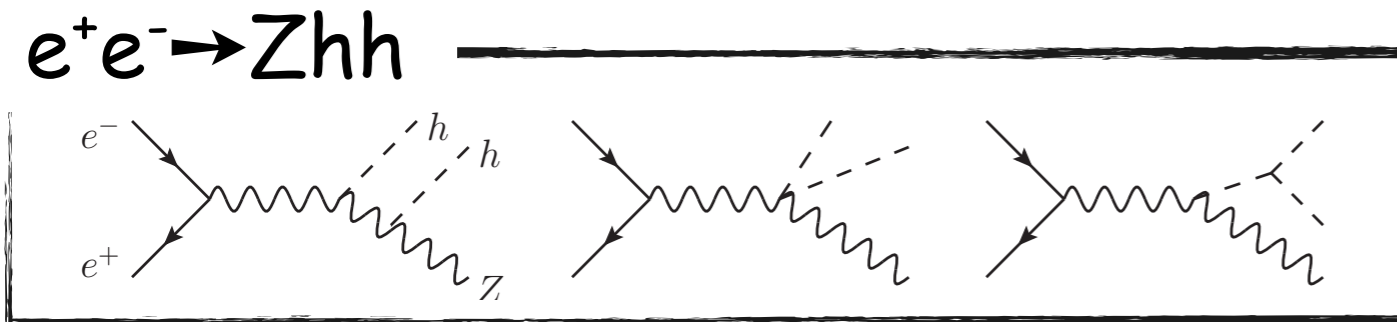
$$\Delta = \left| \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \right| \approx 15\%$$

$\approx 5\%$



$$\Delta = \left| \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \right| \lesssim 4\%$$

$\approx 10\%$



$$\Delta = \left| \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \right| \approx 50\%$$

$\approx 10\%$



# Signs of New Particles?

Espinosa, Grojean, Muhlleitner, Trott '12

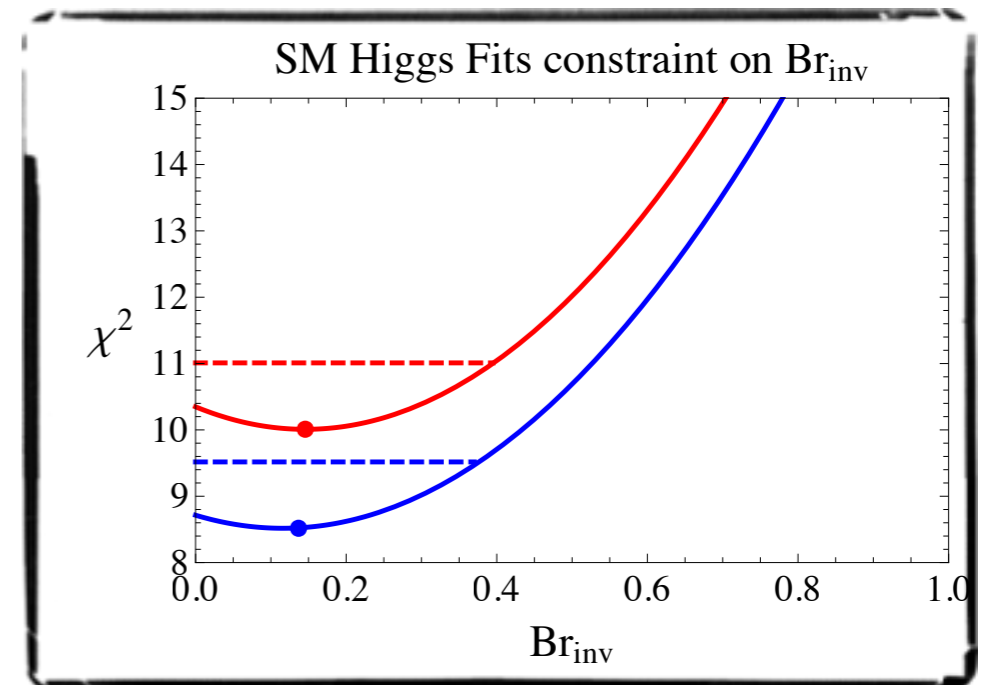
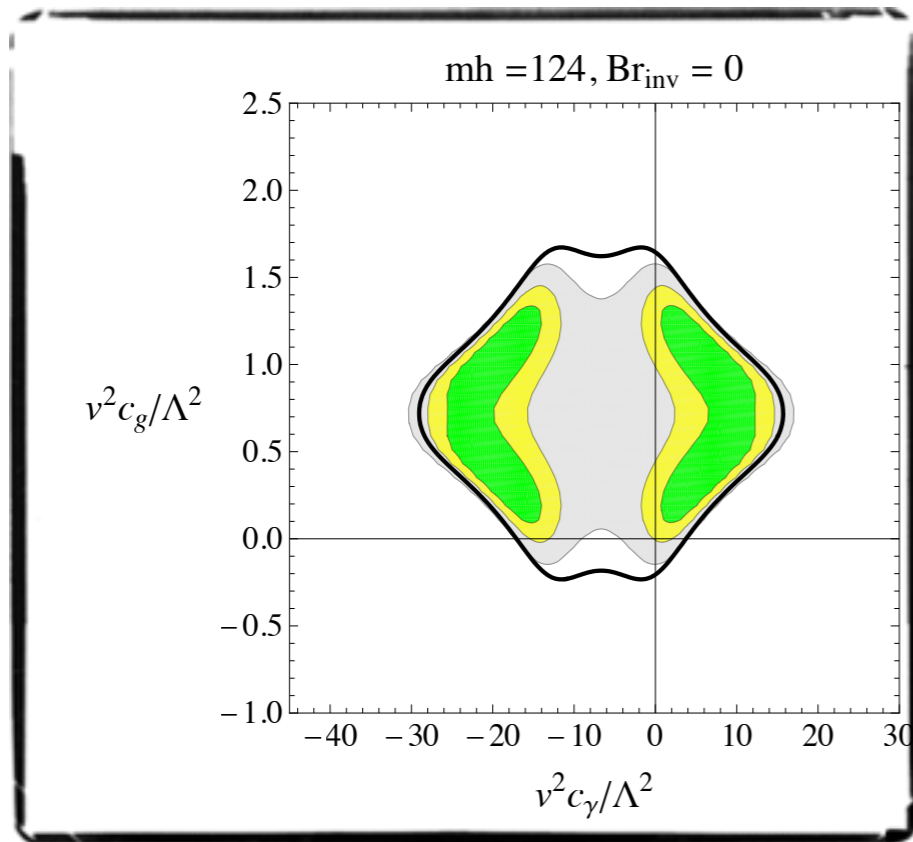
The Higgs can couple to new particles

charged under  
SM gauge group

neutral under  
SM gauge group

$$\mathcal{L} = -\frac{\tilde{c}_\gamma e^2}{32\pi^2 \Lambda^2} H^\dagger H F_{\mu\nu} F^{\mu\nu} - \frac{\tilde{c}_g g_s^2}{32\pi^2 \Lambda^2} H^\dagger H G_{\mu\nu}^a G^{a\mu\nu}$$

$$\text{Br}(h \rightarrow f) \equiv \frac{\Gamma(h \rightarrow f)}{\Gamma_{\text{SM}} + \Gamma_{\text{inv}}} = (1 - \text{Br}_{\text{inv}}) \times \text{Br}_{\text{SM}}(h \rightarrow f)$$



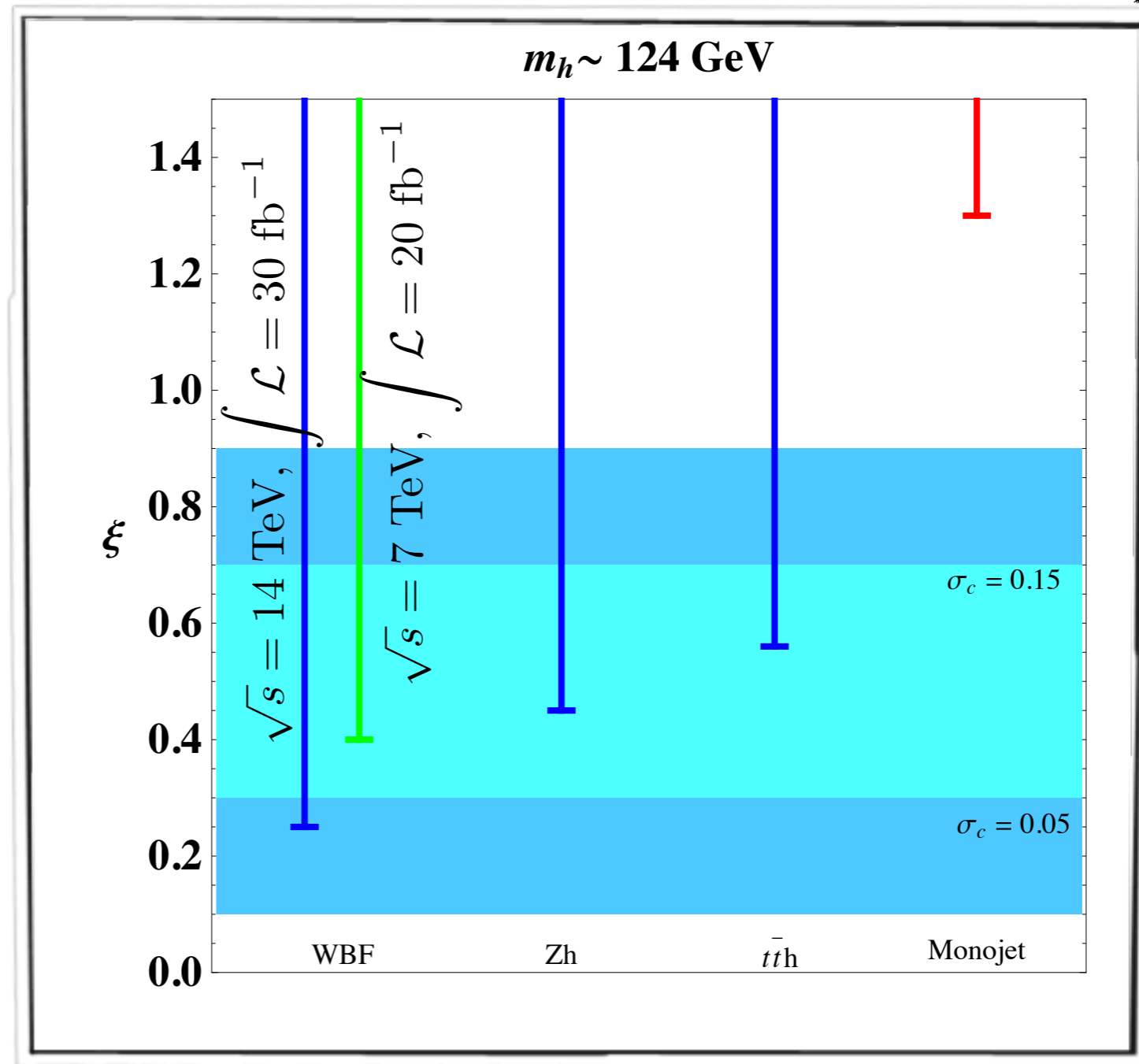


# Search for Invisible Decays with Visible Channels

Espinosa, Grojean, Muhlleitner, Trott '12

direct (vertical) vs indirect (horizontal) searches

Values of  
 $\xi = \sigma / \sigma_{SM} BR(h \rightarrow inv)$   
 for which a 95%CL limit can be imposed



# The Question of the next Decade(s)

What is this Higgs boson that might have been discovered at  $\sim 125\text{GeV}$ ?

*"Higgs = emergency tire of the SM"*

Altarelli @ Blois'10



[picture courtesy to Andreas Weiler]