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Substructure enhancement in Dark Matter indirect detection with γ-rays



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Dark matter is a cornerstone of modern cosmology... But we don't know what it's made of!

Motivation I

★ If DM is a WIMP, indirect detection is a promising possibility to learn about its nature and properties:

In regions of high DM density in the Universe, DM can annihilate emitting photons, positrons, antiprotons or neutrinos.



- ★ If DM is a WIMP (cold relic), standard structure formation tells you that you should expect DM to clump on all scales down to the freestreaming scale.
- ★ Clumping means enhanced annihilation rates for indirect detection!



What are the implications on the limits? Theoretical uncertainty?

 Here we analyze in detail the galactic signal, which is subject to less uncertainty than the extragalactic one.

Motivation II

★ To derive limits, we use the isotropic diffuse component in the sky measured by Fermi-LAT:

The isotropic diffuse component represents roughly 25% of the total flux (for $|b| > 10^\circ$).





[Abdo, et al., 1002.3603]



- Gamma-ray emission from DM annihilations
 Which direction in the sky?
- Transport of final state electrons and positrons
 Effect of diffusion on the gamma-ray emission
- ★ Galactic substructure: Minimal halo mass and mass function index
- Results: fluxes towards the galactic anticenter, and high latitudes
 Flux enhancement due to substructure (Boost factor)
- ★ Constraints on DM annihilation cross-sections

★ Conclusions

Gamma-rays from Dark Matter I

★ Dark matter annihilation can emit photons in many ways:



Gamma-rays from Dark Matter II

- ★ Dark matter annihilation can take place in our galaxy or outside. Here we concentrate on the galactic contribution only.
- The differential gamma-ray flux from DM annihilation within our galaxy is given by
 Direct output from Pythia

PROMPT

IC

$$= \frac{1}{4\pi} \frac{\langle \sigma v \rangle}{2} r_{\odot} \frac{\rho_{\odot}^2}{M_{\chi}^2} \frac{\mathrm{d}N_{\gamma}}{\mathrm{d}E_{\gamma}} \int \mathrm{d}\Omega \int_{\mathrm{los}} \frac{\mathrm{d}s}{r_{\odot}} \left(\frac{\rho(r)}{\rho_{\odot}}\right)^2$$

Components of the IRF

$$\frac{\mathrm{d}\Phi}{\mathrm{d}E_{\gamma}} = \frac{1}{4\pi} r_{\odot} \int \mathrm{d}\Omega \int_{\mathrm{los}} \frac{\mathrm{d}s}{r_{\odot}} \int_{m_{e}}^{M_{\chi}} \mathrm{d}E \mathcal{N}_{e}(r,E) \sum_{i} \mathcal{P}_{i}(E_{\gamma},E,r)$$

Electron density, calculated from the transport equation

Differential photon power emitted from IC scattering

Where we use an NFW density profile for our MW:

 $\frac{\mathrm{d}\Phi}{\mathrm{d}E_{\gamma}}$

$$\rho(r) = \frac{\rho_s}{r/r_s(1+r/r_s)^2}$$

 $r_s = 20.2 \ {\rm kpc}, \ \ \rho_\odot = 0.395 \ {\rm GeV/cm^3}, \ \ r_\odot = 8.29 \ {\rm kpc}$

Direction in the sky

★ When constraining DM annihilation cross-sections with the IGRB, it is customary to calculate the gamma-ray flux in the direction where it is minimal.



galactic anticenter (b=0°, l=180°) when the DM halo is smooth.

★ Here we argue that the direction of the highest latitudes (galactic poles, b=90°, l=0°) can also be used for the following reason:

The galactic diffuse component is dominated in this direction by proton-gas emission and γ -ray sources, which are subject to little uncertainty!



Residual flux at the level of the IGRB!

(also: the presence of substructure makes the signal more isotropic...)



Transport of galactic electrons

★ The diffusion-loss equation for electrons in steady state is given by



The diffusion-loss equation can be solved analytically in the absence of boundary conditions, and if energy losses are independent of position (true for the CMB!).

 $\mathcal{N}_e(\vec{x}, E) = \frac{1}{b(E)} \int_{E_s = E}^{E_s = \infty} dE_s \int d^3 \vec{x}_s \, G_e(\vec{x}_s, E_s \to \vec{x}, E) \, Q(\vec{x}_s, E_s)$

with the Green's function $G_e(\vec{x}_s, E_s \to \vec{x}, E) = \frac{1}{(4\pi K_0 \tilde{\tau})^{3/2}} \exp\left(-\frac{|\vec{x} - \vec{x}_s|^2}{4K_0 \tilde{\tau}}\right)$ given by

★ The assumption of no-diffusion corresponds to the limit

$$G_e(\vec{x}_s, E_s \to \vec{x}, E) \to \delta^3(\vec{x}_s - \vec{x})$$

in which case the gamma-ray spectrum from IC scattering is given by

$$\frac{\mathrm{d}N_{\gamma}^{\mathrm{IC}}}{\mathrm{d}E_{\gamma}}(r) = \int_{m_e}^{M_{\chi}} \mathrm{d}E \ \frac{\sum_i \mathcal{P}_i(E_{\gamma}, E, r)}{\sum_i b_i(E, r)} \int_E^{M_{\chi}} \mathrm{d}E_s \frac{\mathrm{d}N_e}{\mathrm{d}E_s}$$

Galactic substructure I

★ It is expected on theoretical grounds and confirmed in N-body simulations that DM forms clumps on a wide range of scales.

$$\rho_{\rm tot}(r) = \rho_{\rm sm}(r) + \rho_{\rm sub}(r)$$

$$\rho_{\rm sub}(r) = \frac{\rho_{\rm tot}(r)}{1 + r/r_b} \frac{r}{r_b}$$
Anti-biased distribution of subt

Bias radius

Anti-biased distribution of subhalos (see Via Lactea II simulation)

★ Knowing the distribution of clumps in our MW is of crucial importance to estimate the flux from DM annihilations. We use the formalism of probability functions:

$$\frac{dN_{\rm cl}(r,M_{\rm cl})}{dVdM_{\rm cl}} = N_{\rm cl} \frac{d\mathcal{P}_M(M_{\rm cl})}{dM_{\rm cl}} \frac{d\mathcal{P}_V(r)}{dV}$$

Mass function index!

Mass distribution function:

$$\frac{\mathrm{d}\mathcal{P}_M(M_{\mathrm{cl}})}{\mathrm{d}M_{\mathrm{cl}}}(M_{\mathrm{cl}}) = K_m \left(\frac{M_{\mathrm{cl}}}{M_{\odot}}\right)^{-\alpha_m}$$

 $\int_{0}^{R_{\rm vir}} \frac{\mathrm{d}\mathcal{P}_V(r)}{\mathrm{d}V} \mathrm{d}V = 1$ $\int_{M_{\rm min}}^{M_{\rm max}} \frac{\mathrm{d}\mathcal{P}_M(M_{\rm cl})}{\mathrm{d}M_{\rm cl}} dM_{\rm cl} = 1$

Minimal subhalo mass

Spatial distribution function: $\frac{\mathrm{d}\mathcal{P}_V(r)}{\mathrm{d}V} = \frac{\rho_{\mathrm{sub}}(r)}{M_{\mathrm{sub}}^{\mathrm{tot}}}$ Anti-biased!

Galactic substructure II

- ★ The mass function index and the minimal halo mass are the two most crucial parameters.
- ★ The minimal halo mass depends on the precise interactions of the DM particle with the SM, as it derives from the kinetic decoupling temperature. Here we consider $M_{\min} \in (10^{-11} M_{\odot}, 10^{-4} M_{\odot})$
- ★ The mass function index can be accessed in N-body simulations (VLII, Aquarius), but their resolution is still very far from M_{min} . The latest simulations find $\alpha_m = 1.9$ whereas the Press-Schechter theory (and extended versions) on the smallest scales predict $\alpha_m = 2$.

Here we choose to vary
$$lpha_m \in (1.9,2)$$

α_m	$M_{\rm min} = 10^{-11} M_{\odot}$	$M_{ m min} = 10^{-4} M_{\odot}$
	$f_{ m sub}^{ m tot}=0.699$	$f_{ m sub}^{ m tot} = 0.467$
2	$N_{ m sub}^{ m tot}=2.66 imes10^{21}$	$N_{ m sub}^{ m tot}=2.66 imes10^{14}$
	$r_b = 35.08 \ \mathrm{kpc}$	$r_b = 117.63 \ \mathrm{kpc}$
	$f_{ m sub}^{ m tot}=0.187$	$f_{ m sub}^{ m tot}=0.181$
1.9	$N_{ m sub}^{ m tot}=3.06 imes10^{19}$	$N_{ m sub}^{ m tot} = 1.54 imes 10^{13}$
	$r_b = 557.11 \text{ kpc}$	$r_b = 582.30 \ \rm kpc$

Gamma-ray fluxes on Earth

The gamma-ray flux on Earth from DM annihilations in our Galaxy can be calculated to be:



 The boost factor gives the flux enhancement due to the presence of substructure in our Galaxy.

Boost factor

Boost $\equiv \frac{\mathrm{d}\Phi^{\mathrm{sub}}/\mathrm{d}E_{\gamma} + \mathrm{d}\Phi^{\mathrm{smooth}}/\mathrm{d}E_{\gamma}}{\mathrm{d}\Phi^{\mathrm{nosub}}/\mathrm{d}E_{\gamma}}$ HLGAC 10 $\alpha_m = 2$ $\alpha_m = 2$ 10Boost $\alpha_m = 1.9$ $\alpha_m = 1.9$ 1 1 0.110 100 1000 0.110 1 1001000 1 E_{γ} [GeV] E_{γ} [GeV]

 $M_{\min} \in (10^{-11} M_{\odot}, 10^{-4} M_{\odot})$

Steve Blanchet, PLANCK2012, Warsaw, 30/05/12

Boost





★ The signal from the highest latitudes can be constrained by the IGRB measurement.

Conclusions

- ★ Diffusion of final state electrons/positrons plays a marginal role both for the galactic anticenter and the poles.
- ★ We have taken into account DM galactic substructure, in agreement with recent N-body simulations. The two most relevant parameters are the mass function index, and the minimal subhalo mass.
- ★ We found that substructure can boost the signal by up to a factor of 20. With the most pessimistic assumptions, the boost is as low as 20%.
- ★ We extracted exclusion limits for DM annihilation cross-sections, and found our limits for optimistic choices of the mass function index to be competitive with the most stringent to date.



Galactic diffuse emission



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Emission spectrum



Transport of galactic electrons



Galactic substructure

The gamma-ray flux from annihilations in galactic substructure can then be calculated with

$$\frac{\mathrm{d}\Phi}{\mathrm{d}E_{\gamma}} = \frac{1}{4\pi} \frac{\langle \sigma v \rangle}{2} r_{\odot} \frac{\rho_{\odot}^{2}}{M_{\chi}^{2}} \frac{\mathrm{d}N_{\gamma}}{\mathrm{d}E_{\gamma}} \int \mathrm{d}\Omega \int_{\mathrm{los}} \frac{\mathrm{d}s}{r_{\odot}} \frac{\mathrm{d}\mathcal{P}_{V}}{\mathrm{d}V}(r) \int_{M_{\mathrm{min}}}^{M_{\mathrm{max}}} \mathrm{d}M_{\mathrm{cl}} \underbrace{\boldsymbol{\xi}(M_{\mathrm{cl}}, r)}_{\mathrm{d}M_{\mathrm{cl}}} \frac{\mathrm{d}\mathcal{P}_{M}(M_{\mathrm{cl}})}{\mathrm{d}M_{\mathrm{cl}}} \\ \underbrace{\boldsymbol{\xi}(M_{\mathrm{cl}}, \vec{x}_{s}) \equiv \int_{V_{\mathrm{cl}}} \mathrm{d}V \left(\frac{\rho_{\mathrm{cl}}(M_{\mathrm{cl}}, \vec{x}_{s})}{\rho_{\odot}}\right)^{2}}_{\mathrm{annihilation volume}}$$

★ The most likely origin for the IGRB is from blazars. Assuming that they make most of it, we obtain more stringent constraints:

Exclusion limits

