

PLANCK 2012

FROM THE PLANCK SCALE
TO THE ELECTROWEAK SCALE

28 MAY - 1 JUNE 2012, WARSAW, POLAND

Implications of CPV in Charm for Weak Scale NP

Jernej F. Kamenik



Institut "Jožef Stefan"



Univerza v Ljubljani

Fakulteta za matematiko in fiziko

28/05/2012, Warsaw

Outline

- Significance of CPV in Charm within and beyond SM
 - Quantify (parametrize) theory expectations of direct CPV in charm decays
- Δa_{CP} implications for weak scale NP
 - EFT & models
- new insights into CPV in $SU(3)_Q$ breaking NP models

Why CP Violation in Charm?

- CPV in charm provides a unique probe of New Physics (NP)
 - sensitive to NP in the up sector
 - SM charm physics is CP conserving to first approximation (2 generation dominance, no hard GIM breaking)
- Common lore "any signal for CPV would be NP":
 - In the SM, CPV in mixing enters at $O(V_{cb}V_{ub}/V_{cs}V_{us}) \sim 10^{-3}$
 - In the SM, direct CPV enters at $O([V_{cb}V_{ub}/V_{cs}V_{us}] \alpha_s/\pi) \sim 10^{-4}$ (in singly Cabibbo suppressed decays)

Why CP Violation in Charm?

- CPV in charm provides a unique probe of New Physics (NP)
 - sensitive to NP in the up sector
 - SM charm physics is CP conserving to first approximation (2 generation dominance, no hard GIM breaking)
- Common lore "any signal for CPV would be NP":
 - In the SM, CPV in mixing enters at $O(V_{cb}V_{ub}/V_{cs}V_{us}) \sim 10^{-3}$
 - In the SM, direct CPV enters at $O([V_{cb}V_{ub}/V_{cs}V_{us}] \alpha_s/\pi) \sim 10^{-4}$ (in singly Cabibbo suppressed decays)

Experimental observables

- **CPV in decays (direct CPV)**

- Time-integrated CPV decay asymmetries to CP eigenstates

$$a_f \equiv \frac{\Gamma(D^0 \rightarrow f) - \Gamma(\bar{D}^0 \rightarrow f)}{\Gamma(D^0 \rightarrow f) + \Gamma(\bar{D}^0 \rightarrow f)}.$$

- Focus on K^+K^- and $\pi^+\pi^-$ final states: $\Delta a_{CP} \equiv a_{K^+K^-} - a_{\pi^+\pi^-}$

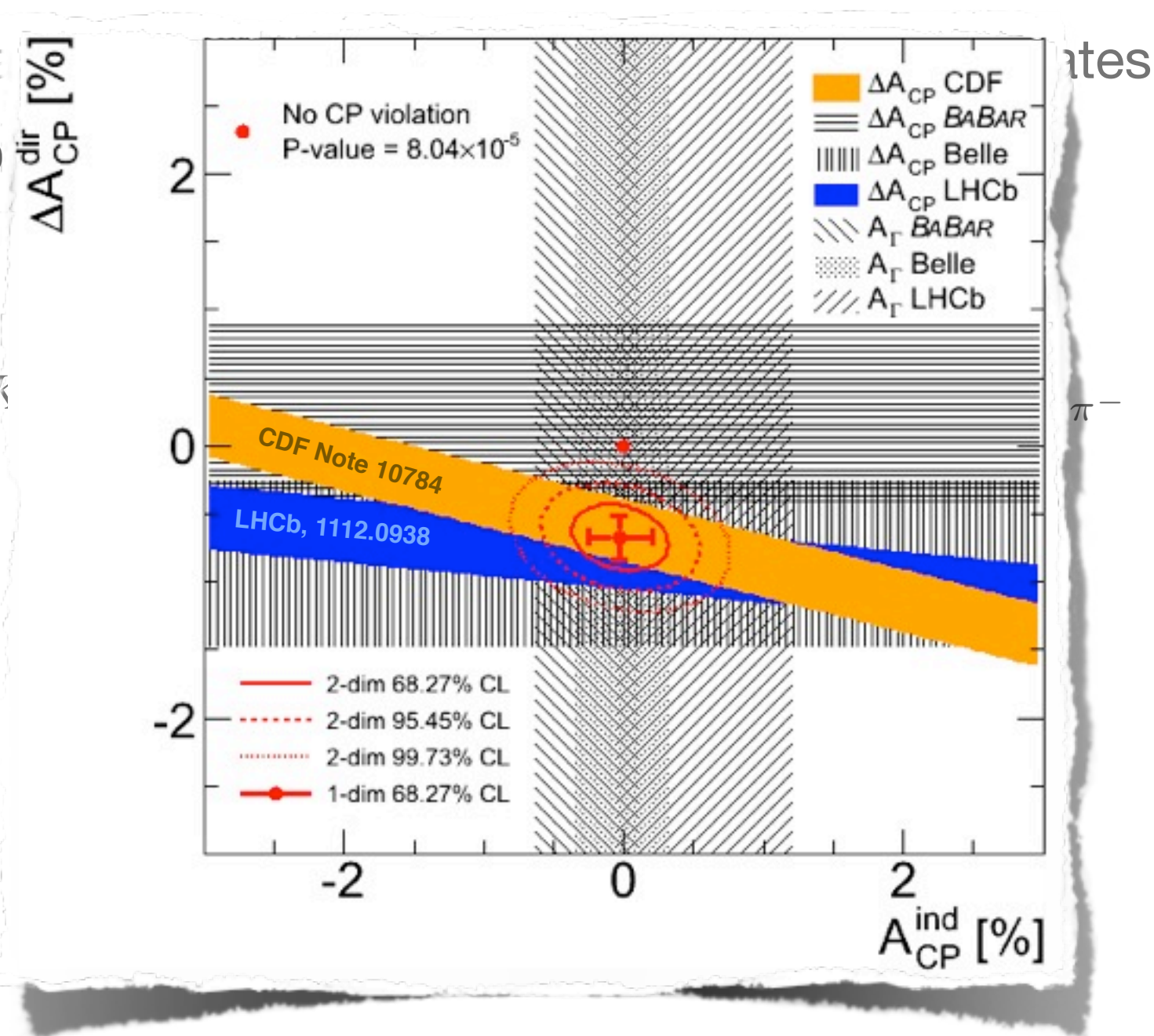
Experimental observables

- CPV in decays (direct CPV)

- Time-integrated

$$a_f \equiv \frac{\Gamma(D^0 \rightarrow f)}{\Gamma(D^0 \rightarrow \bar{f})}$$

- Focus on K^+K^-



Experimental observables

- **CPV in decays (direct CPV)**

- Time-integrated CPV decay asymmetries to CP eigenstates

$$a_f \equiv \frac{\Gamma(D^0 \rightarrow f) - \Gamma(\bar{D}^0 \rightarrow f)}{\Gamma(D^0 \rightarrow f) + \Gamma(\bar{D}^0 \rightarrow f)}.$$

- Focus on K^+K^- and $\pi^+\pi^-$ final states: $\Delta a_{CP} \equiv a_{K^+K^-} - a_{\pi^+\pi^-}$

$$\Delta a_{CP}^{\text{World}} = -(0.67 \pm 0.16)\% \quad (\sim 4\sigma \text{ from } 0)$$

Experimental observables

- **CPV in decays (direct CPV)**

- Time-integrated CPV decay asymmetries to CP eigenstates

$$a_f \equiv \frac{\Gamma(D^0 \rightarrow f) - \Gamma(\bar{D}^0 \rightarrow f)}{\Gamma(D^0 \rightarrow f) + \Gamma(\bar{D}^0 \rightarrow f)}.$$

- Focus on K^+K^- and $\pi^+\pi^-$ final states: $\Delta a_{CP} \equiv a_{K^+K^-} - a_{\pi^+\pi^-}$

$$\Delta a_{CP}^{\text{World}} = -(0.67 \pm 0.16)\% \quad (\sim 4\sigma \text{ from } 0)$$

- ...beyond natural expectation within the SM

- but not possible to prove from first principles, and **SM-like explanation cannot be excluded**

Grossman et al., hep-ph/0609178

Cheng & Chiang, 1201.0785

Franco, Mishima & Silvestrini, 1203.3131

Li, Lu & Yu, 1203.3120

Golden & Grinstein Phys. Lett. B 222 (1989)

Brod, Kagan & Zupan 1111.5000

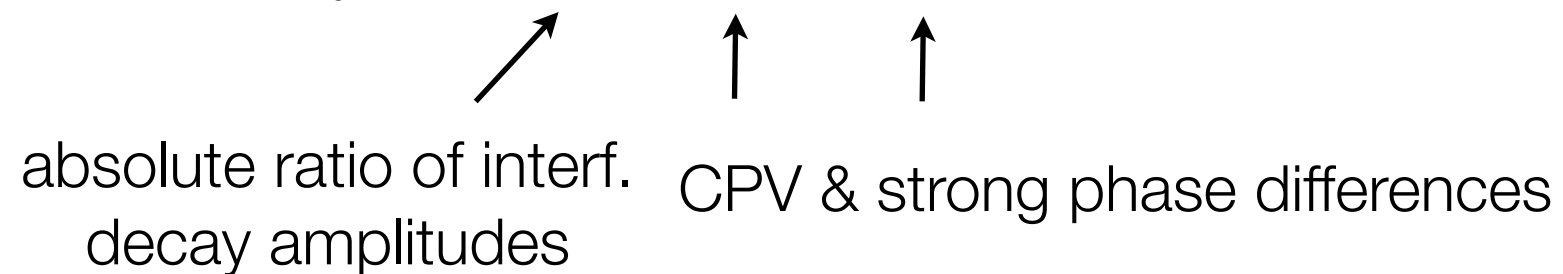
Brod, Grossman, Kagan & Zupan 1203.6659

Feldmann, Nandi & Soni

1202.3795


Quantifying theoretical predictions

- SM expectations for $a_f^{\text{dir.}} \simeq 2r_f \sin \phi_f \sin \delta_f$ $f = K, \pi$


absolute ratio of interf.
decay amplitudes CPV & strong phase differences

Quantifying theoretical predictions

- SM expectations for $a_f^{\text{dir.}} \simeq 2r_f \sin \phi_f \sin \delta_f$ $f = K, \pi$



 absolute ratio of interf. decay amplitudes CPV & strong phase differences

- in terms of weak decay amplitudes $\lambda_q \equiv V_{cq}^* V_{uq}$, $\lambda_d + \lambda_s + \lambda_b = 0$


$$A_K = \lambda_s A_K^{(1)} + \lambda_b A_K^{(2)}$$

$$A_\pi = \lambda_d A_\pi^{(1)} + \lambda_b A_\pi^{(2)}$$

- naive expectation $A_f^{(2)} / A_f^{(1)} \sim \frac{\alpha_s(m_c)}{\pi}$
 (assuming $m_c \gg \Lambda_{QCD}$)

Quantifying theoretical predictions

- SM expectations for $a_f^{\text{dir.}} \simeq 2r_f \sin \phi_f \sin \delta_f$ $f = K, \pi$



 absolute ratio of interf. decay amplitudes CPV & strong phase differences

- in terms of weak decay amplitudes $\lambda_q \equiv V_{cq}^* V_{uq}$, $\lambda_d + \lambda_s + \lambda_b = 0$

$$A_K = \lambda_s A_K^{(1)} + \lambda_b A_K^{(2)}$$

$$A_\pi = \lambda_d A_\pi^{(1)} + \lambda_b A_\pi^{(2)}$$




$$r_f \propto \xi = |\lambda_b / \lambda_s| \simeq |\lambda_b / \lambda_d| \approx 0.0007$$

- naive expectation $A_f^{(2)} / A_f^{(1)} \sim \frac{\alpha_s(m_c)}{\pi}$ $\phi_K^{\text{SM}} = \arg(\lambda_b / \lambda_s)$
 (assuming $m_c \gg \Lambda_{QCD}$) $\approx -\arg(\lambda_b / \lambda_d) = -\phi_\pi^{\text{SM}}$
 $\approx 70^\circ$

Quantifying theoretical predictions

- SM expectations for $a_f^{\text{dir.}} \simeq 2r_f \sin \phi_f \sin \delta_f$ $f = K, \pi$



 absolute ratio of interf. decay amplitudes CPV & strong phase differences

- in terms of weak decay amplitudes $\lambda_q \equiv V_{cq}^* V_{uq}$, $\lambda_d + \lambda_s + \lambda_b = 0$

$$A_K = \lambda_s A_K^{(1)} + \lambda_b A_K^{(2)}$$

$$A_\pi = \lambda_d A_\pi^{(1)} + \lambda_b A_\pi^{(2)}$$



$$r_f \propto \xi = |\lambda_b / \lambda_s| \simeq |\lambda_b / \lambda_d| \approx 0.0007$$

- naive expectation $A_f^{(2)} / A_f^{(1)} \sim \frac{\alpha_s(m_c)}{\pi}$ $\phi_K^{\text{SM}} = \arg(\lambda_b / \lambda_s)$
 (assuming $m_c \gg \Lambda_{QCD}$) $\approx -\arg(\lambda_b / \lambda_d) = -\phi_\pi^{\text{SM}}$
 $\approx 70^\circ$

$$\Delta a_{CP} \approx (0.13\%) \text{Im}(\Delta R^{\text{SM}}),$$

$$\Delta R^{\text{SM}} \equiv \frac{A_K^{(2)}}{A_K^{(1)}} + \frac{A_\pi^{(2)}}{A_\pi^{(1)}}$$

Quantifying theoretical predictions

- SM expectations for $a_f^{\text{dir.}} \simeq 2r_f \sin \phi_f \sin \delta_f$ $f = K, \pi$



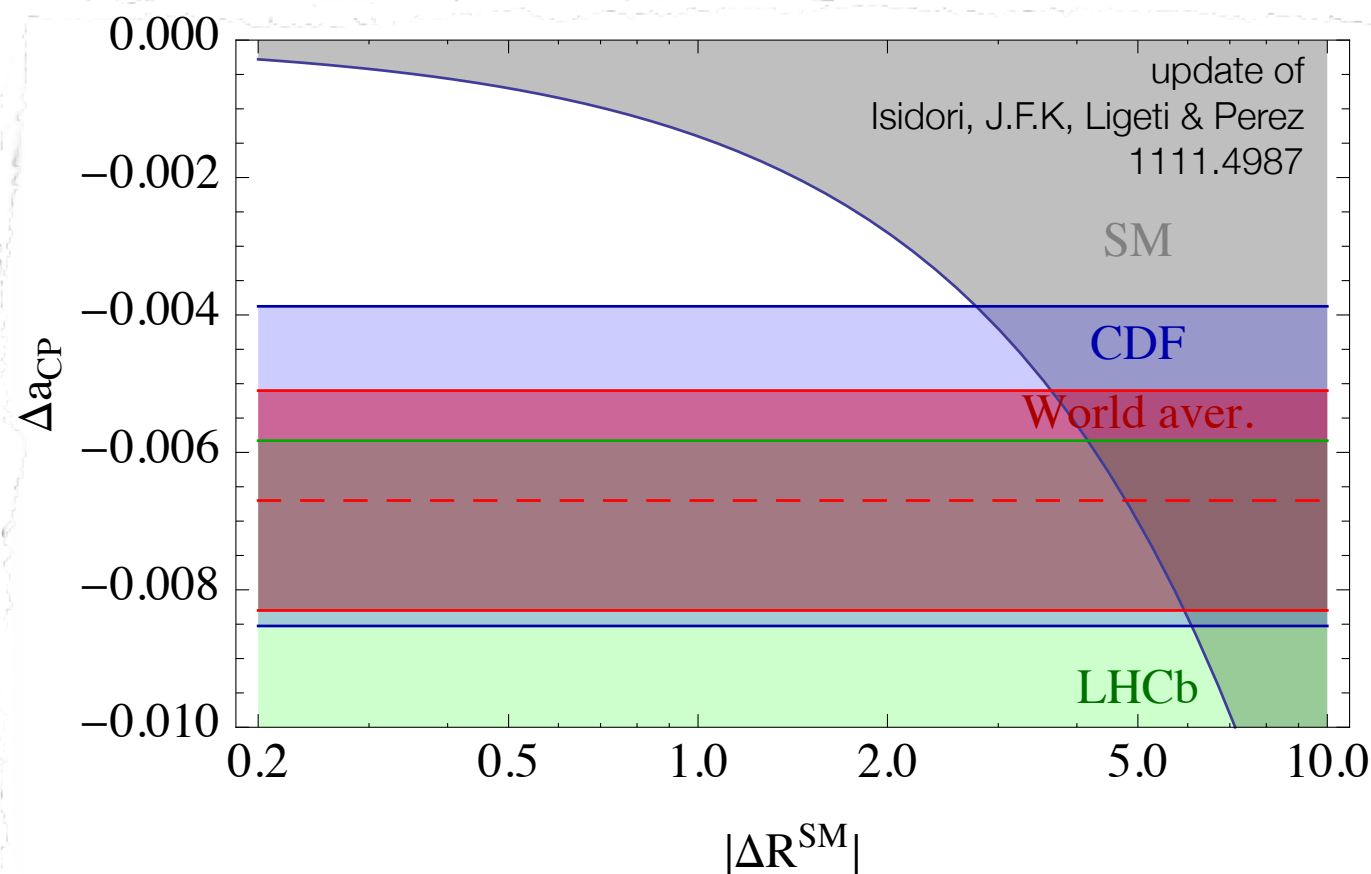
- in terms

$$A_K = \lambda$$

$$A_\pi = \lambda$$

- naive ex

(assuming



differences

$$\lambda_d + \lambda_s + \lambda_b = 0$$

$$s| \simeq |\lambda_b/\lambda_d| \approx 0.0007$$

$$/\lambda_s)$$

$$\lambda_b/\lambda_d) = -\phi_\pi^{\text{SM}}$$

$$\Delta a_{CP} \approx (0.13\%) \text{Im}(\Delta R^{\text{SM}}),$$

$$\Delta R^{\text{SM}} \equiv \frac{A_K^{(2)}}{A_K^{(1)}} + \frac{A_\pi^{(2)}}{A_\pi^{(1)}}$$

Quantifying theoretical predictions

- SM expectations for $a_f^{\text{dir.}} \simeq 2r_f \sin \phi_f \sin \delta_f$ $f = K, \pi$



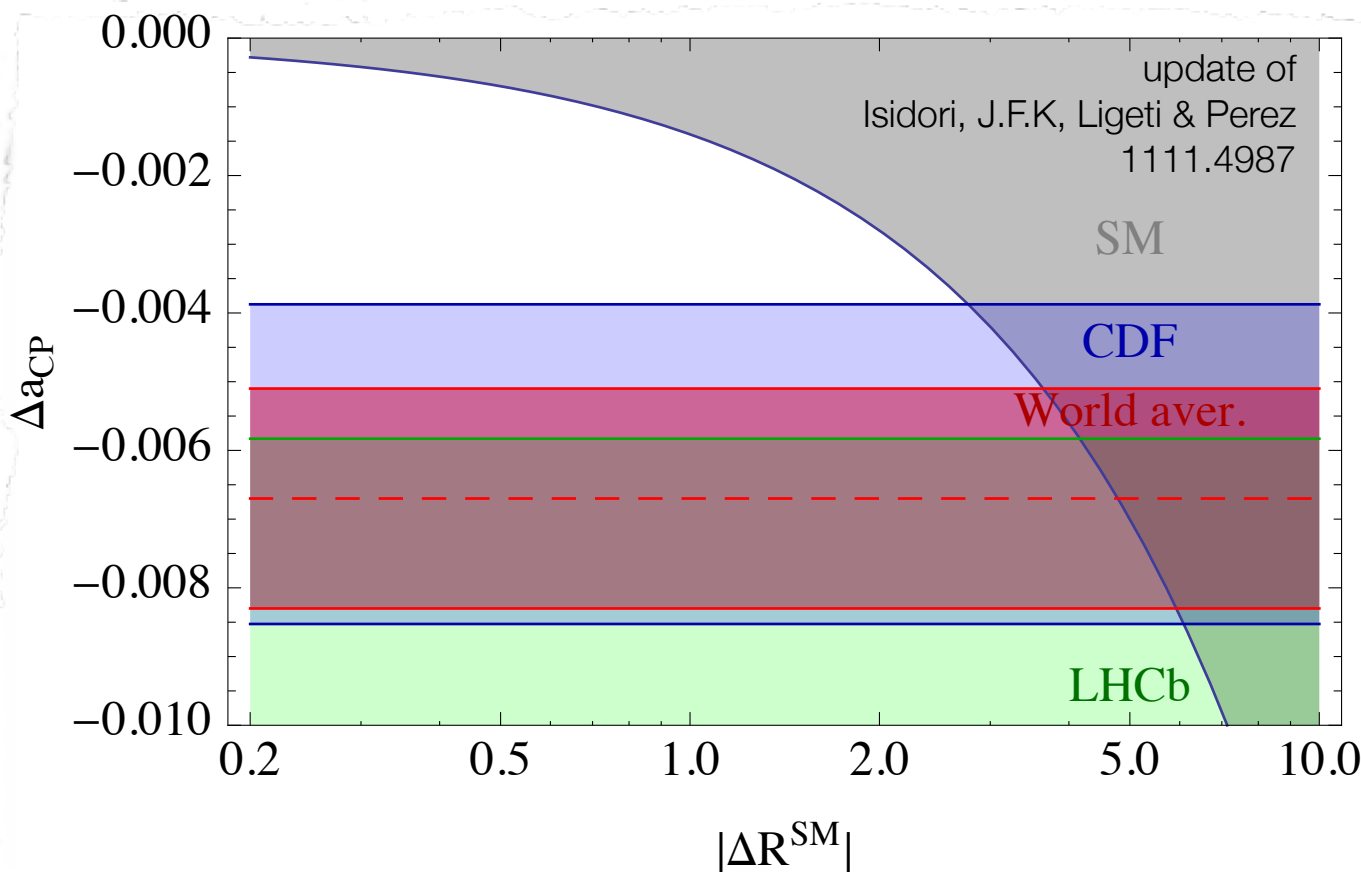
- in terms

$$A_K = \lambda$$

$$A_\pi = \lambda$$

- naive ex

(assuming



differences

$$\lambda_d + \lambda_s + \lambda_b = 0$$

$$s| \simeq |\lambda_b/\lambda_d| \approx 0.0007$$

$$/\lambda_s)$$

$$\lambda_b/\lambda_d) = -\phi_\pi^{\text{SM}}$$

$$\Delta a_{CP} \approx (0.13\%) \text{Im}(\Delta R^{\text{SM}}),$$

$$\Delta R^{\text{SM}} \equiv \frac{A_K^{(2)}}{A_K^{(1)}} + \frac{A_\pi^{(2)}}{A_\pi^{(1)}}$$

0(4-6) values of $|\Delta R^{\text{SM}}|$ needed

Δa_{CP} and New Physics

- Assume SM does not saturate the experimental value

- Parametrize NP contributions in EFT normalized to the effective SM scale

$$\mathcal{H}_{|\Delta c|=1}^{\text{eff-NP}} = \frac{G_F}{\sqrt{2}} \sum_i C_i^{\text{NP}} Q_i$$

$$Q_1^q = (\bar{u}q)_{V-A} (\bar{q}c)_{V-A}$$

$$Q_2^q = (\bar{u}_\alpha q_\beta)_{V-A} (\bar{q}_\beta c_\alpha)_{V-A},$$

$$Q_5^q = (\bar{u}c)_{V-A} (\bar{q}q)_{V+A},$$

$$Q_6^q = (\bar{u}_\alpha c_\beta)_{V-A} (\bar{q}_\beta q_\alpha)_{V+A},$$

$$Q_7 = -\frac{e}{8\pi^2} m_c \bar{u} \sigma_{\mu\nu} (1 + \gamma_5) F^{\mu\nu} c,$$

$$Q_8 = -\frac{g_s}{8\pi^2} m_c \bar{u} \sigma_{\mu\nu} (1 + \gamma_5) T^a G_a^{\mu\nu} c,$$

+ Ops. with $V \leftrightarrow A$

x 5 $q\bar{q}$ flavor structures

- most general dim 6 Hamiltonian at $\mu < m_{W,t}$

Δa_{CP} and New Physics

- Assume SM does not saturate the experimental value
- Parametrize NP contributions in EFT normalized to the effective SM scale

$$\Delta a_{CP} \approx (0.13\%) \text{Im}(\Delta R^{\text{SM}}) + 9 \sum_i \text{Im}(C_i^{\text{NP}}) \text{Im}(\Delta R_i^{\text{NP}}) \quad \Delta R_i^{\text{NP}} \equiv \frac{G_F}{\sqrt{2}} \sum_{f=\pi, K} \frac{\langle Q_i \rangle_f}{A_f^{(1)}}$$

- for $\text{Im}(C_i^{\text{NP}}) = \frac{v^2}{\Lambda^2} : \frac{(10 \text{ TeV})^2}{\Lambda^2} = \frac{(0.61 \pm 0.17) - 0.12 \text{Im}(\Delta R^{\text{SM}})}{\text{Im}(\Delta R^{\text{NP}})}$

Are such contributions allowed by other flavor constraints?

Δa_{CP} and New Physics

Isidori, J.F.K, Ligeti & Perez
1111.4987

- In EFT can be estimated via “weak mixing” of operators

$$T \left\{ \mathcal{H}_{|\Delta c|=1}^{\text{eff-NP}}(0), \mathcal{H}^{\text{SM}}(x) \right\}$$

- Important constraints expected from **D- \bar{D} mixing** and direct CPV in $K^0 \rightarrow \pi^+ \pi^-$ (**ϵ'/ϵ**)

- Quadratic NP contributions

- either chirally suppressed...
- ...or highly UV sensitive

$\propto \alpha m_c^2 / v^2$ $\propto \Lambda^2$

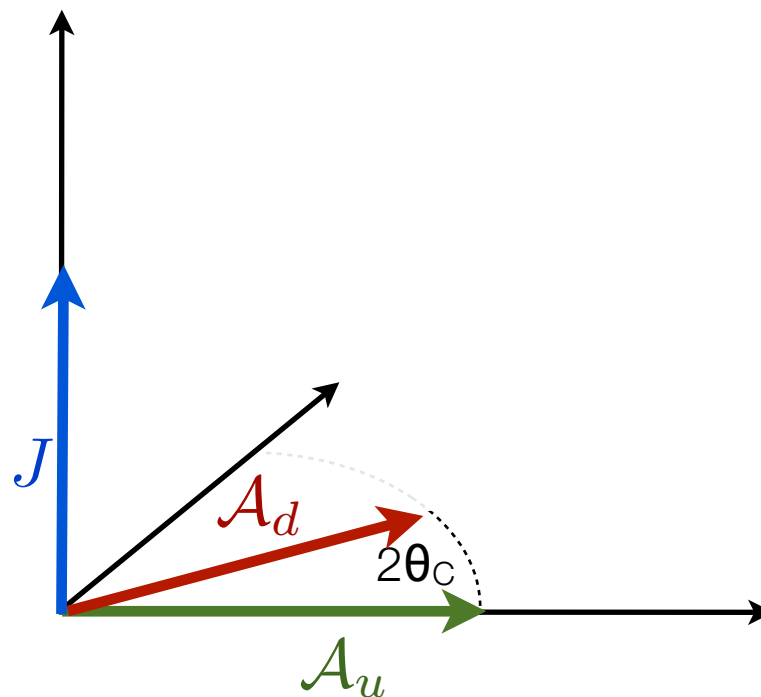
}

On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$

On Universality of CPV in $SU(3)_Q$ breaking NP

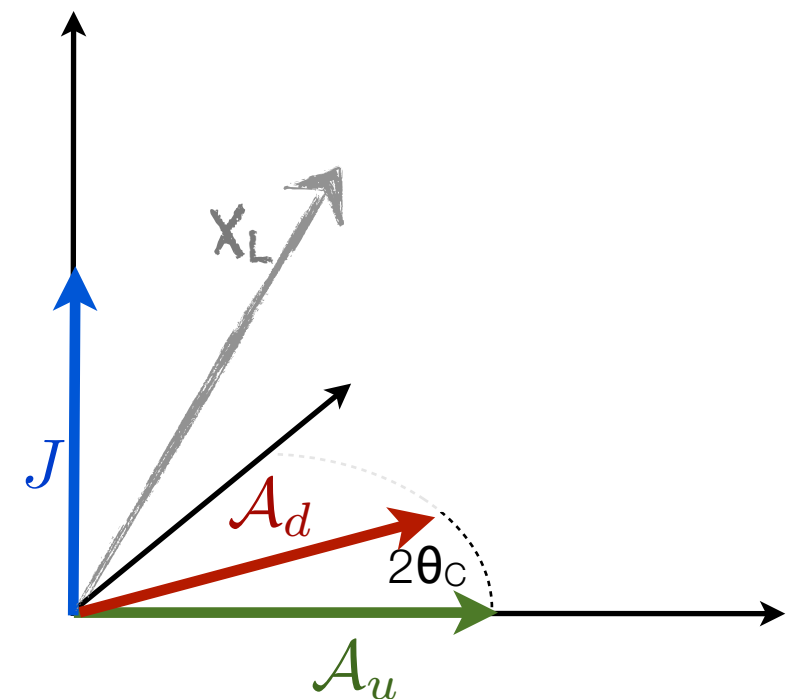
- SM quark flavor symmetry $\mathcal{G}_F = \boxed{SU(3)_Q} \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$
- in the 2-gen limit single source of CPV: $J \equiv i[\mathcal{A}_u, \mathcal{A}_d]$ Gedalia, Mannelli & Perez
1002.0778, 1003.3869
- invariant under $SO(2)$ rotations between up-down mass bases



On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = \boxed{SU(3)_Q} \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$
- in the 2-gen limit single source of CPV: $J \equiv i[\mathcal{A}_u, \mathcal{A}_d]$ Gedalia, Mannelli & Perez
1002.0778, 1003.3869
- invariant under $SO(2)$ rotations between up-down mass bases
- $SU(2)_Q$ breaking NP $\mathcal{O}_L = \left[(X_L)^{ij} \bar{Q}_i \gamma^\mu Q_j \right] L_\mu$

$$\text{Im}(X_L^u)_{12} = \text{Im}(X_L^d)_{12} \propto \text{Tr}(X_L \cdot J) .$$



On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = \boxed{SU(3)_Q} \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$
- SM 3-gen case characterized by $SU(3)/SU(2)$ breaking pattern by $Y_{b,t}$
Kagan et al., 0903.1794
- 3-gen X_L can be decomposed under $SU(2)_Q$, constrained separately
(barring cancelations)
- SM breaking of residual $SU(2)_Q$ suppressed by m_c/m_t , m_s/m_b , θ_{13} , θ_{23}
(charm and kaon sectors dominated by 2-gen physics) see also talk by Barbieri

On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{t}\bar{\text{r}}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{t}\bar{\text{r}}}$
- **Implication (1):** direct correspondence between Δa_{CP} and ε'/ε
(no weak loop suppression)
- constraint on $SU(3)_Q$ breaking NP contributions

$$\Delta a_{CP}^{\text{NP}} \lesssim 4 \times 10^{-4}$$

On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$
- **Implication (2):** bounds on degeneracy in SUSY alignment models
 - SUSY effects in flavor \sim masses, splittings (degeneracy); mixing angles
(e.g. $\bar{q}_i \tilde{q}_j \tilde{g}$)

On Universality of CPV in $SU(3)_Q$ breaking NP

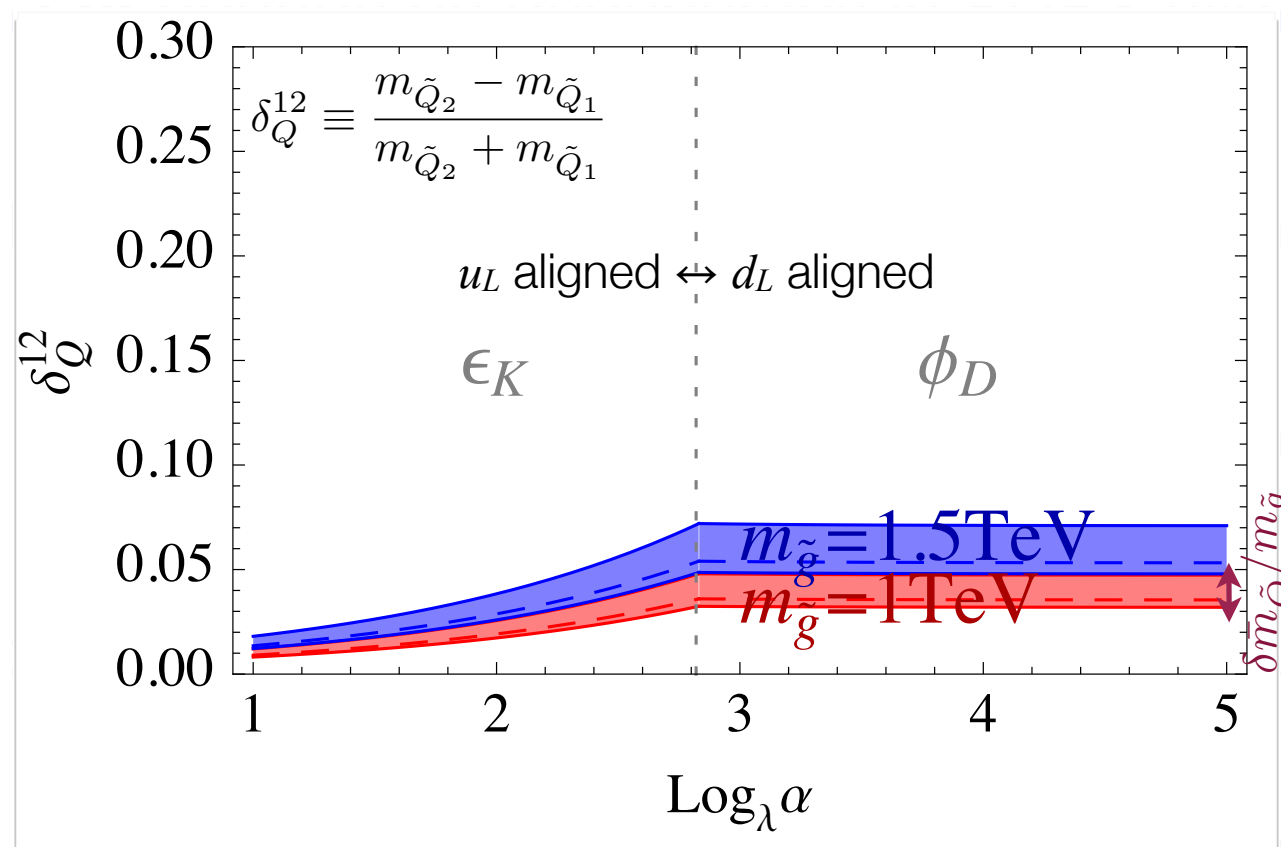
- SM quark flavor symmetry $\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$
- **Implication (2):** bounds on degeneracy in SUSY alignment models
 - SUSY effects in flavor \sim masses, splittings (degeneracy); mixing angles
(e.g. $\bar{q}_i \tilde{q}_j \tilde{g}$)
 - D & K mixing said to imply that **alignment not viable w/o degeneracy**
(Nir & Seiberg, hep-ph/9304307) Blum et al., 0903.2118
 - based on assumption of \sim maximal CPV in K & D mixing
 - not actually attainable in alignment due to CPV universality

On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = \boxed{SU(3)_Q} \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$

• Implications

- SUSY e
- D & K m
- base
- not a



• Implications for flavor models

(e.g. $\bar{q}_i \tilde{q}_j \tilde{g}$)

e w/o degeneracy

Blum et al., 0903.2118

D mixing

universality

On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = \boxed{SU(3)_Q} \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$

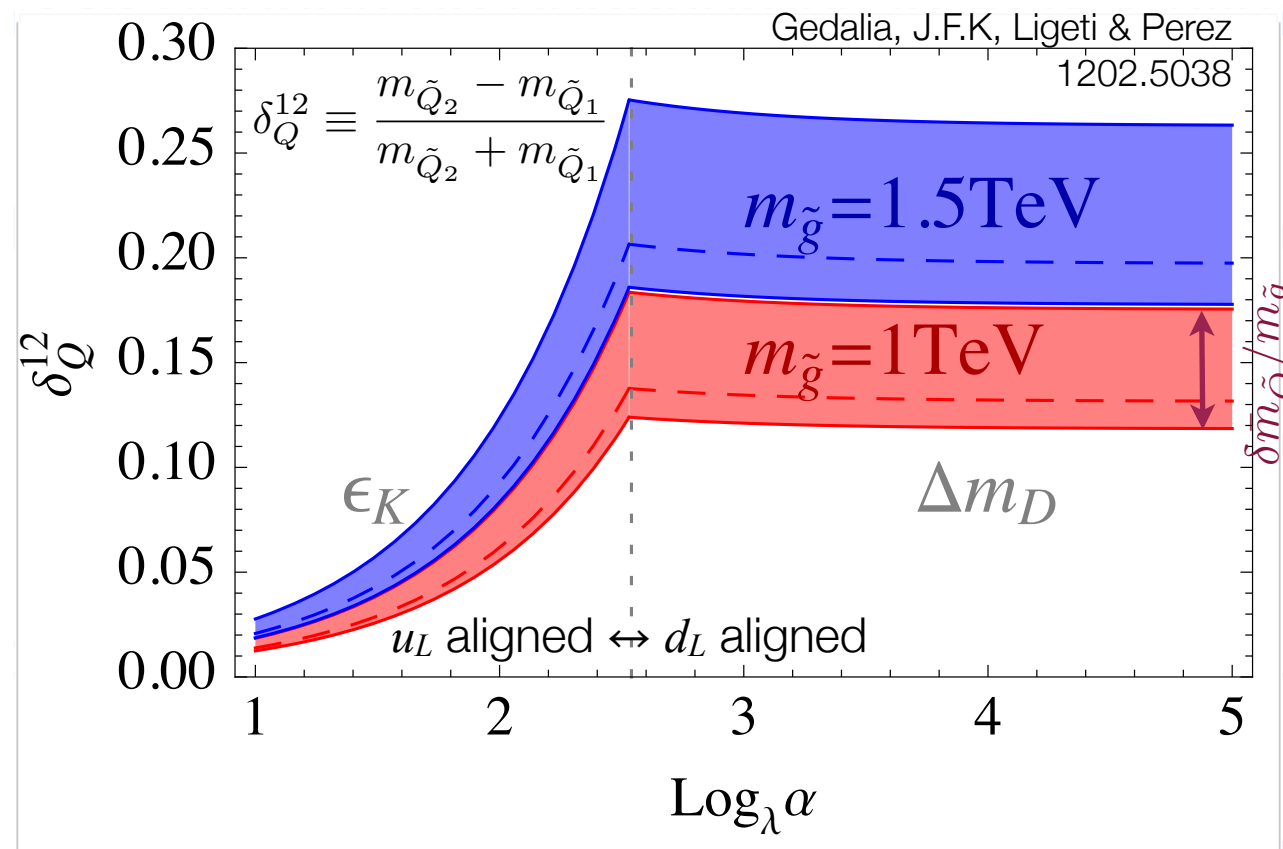
• Implications

- SUSY e

- D & K m

- base

- not a



• Implications for flavor models

eracy); mixing angles
(e.g. $\bar{q}_i \tilde{q}_j \tilde{g}$)

e w/o degeneracy

Blum et al., 0903.2118

D mixing

universality

On Universality of CPV in $SU(3)_Q$ breaking NP

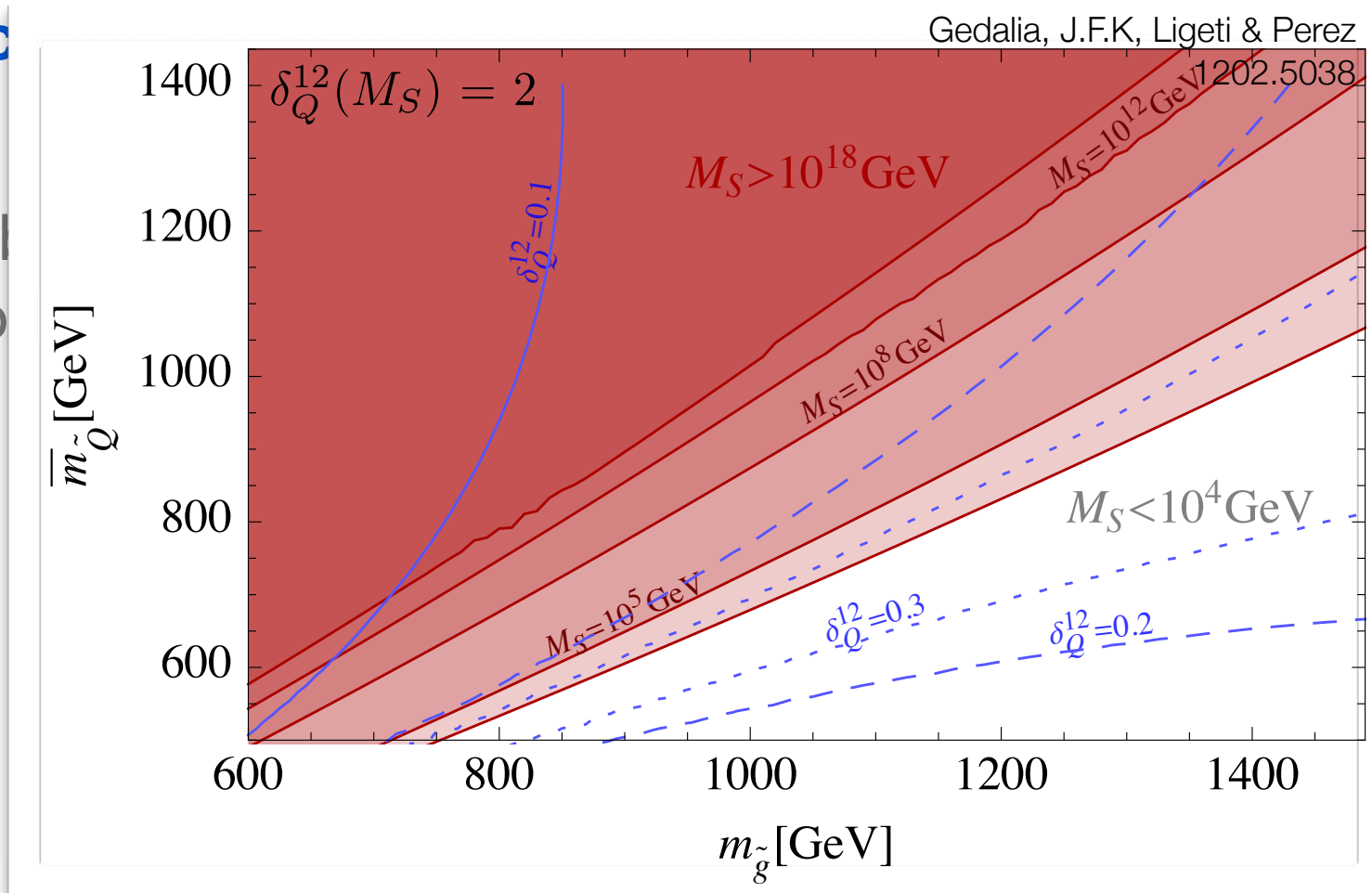
- SM quark flavor symmetry $\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$
- **Implication (2):** bounds on degeneracy in SUSY alignment models
 - viable SUSY spectra can be generated from complete anarchy at moderate mediation scales (M_S) (SUSY QCD RGE)

On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor symmetry $\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$

• Implications

- viable models



• viable models

• complete anarchy at

On Universality of CPV in $SU(3)_Q$ breaking NP

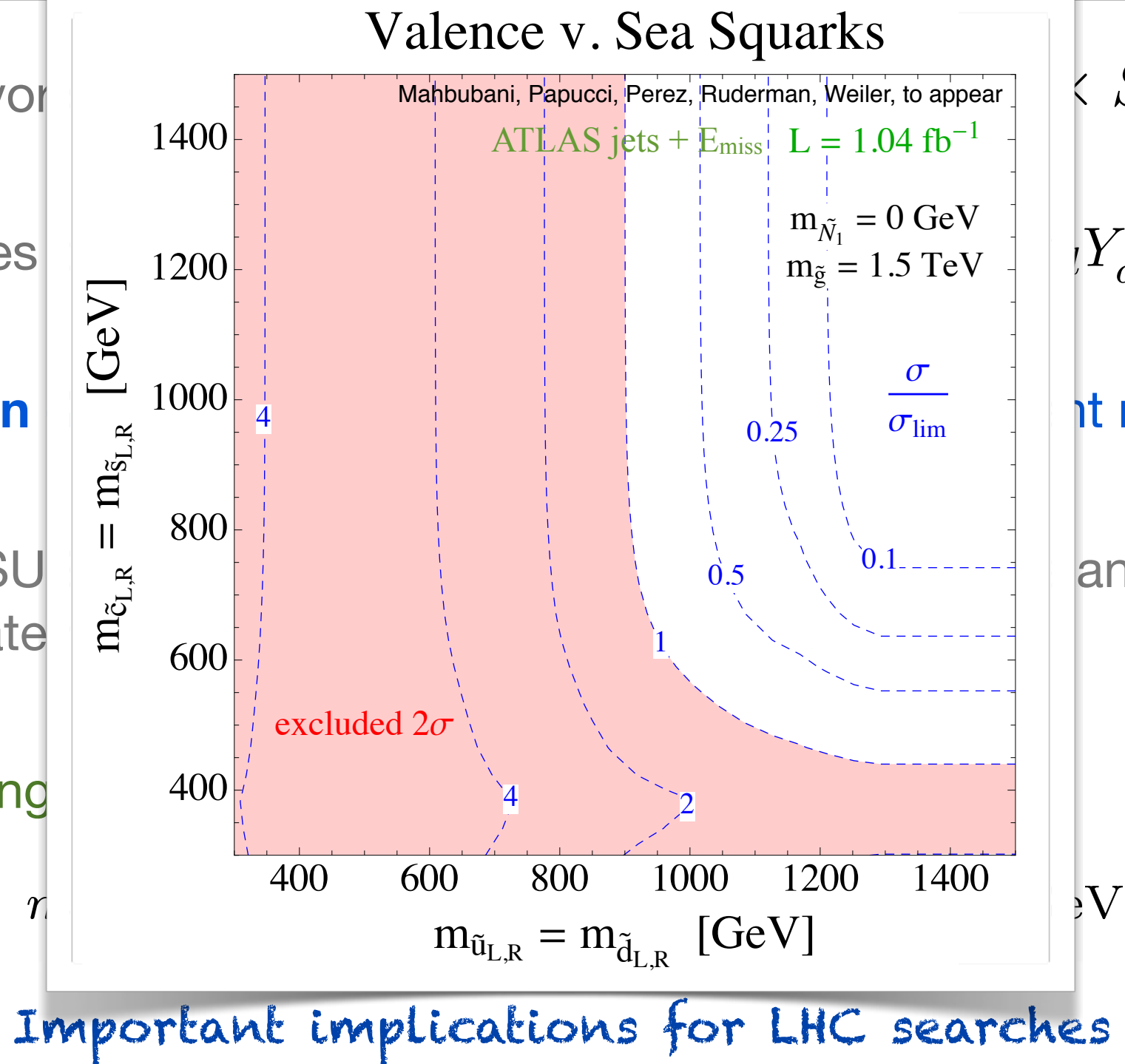
- SM quark flavor symmetry $\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$
- two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}}$, $\mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}}$
- **Implication (2):** bounds on degeneracy in SUSY alignment models
- viable SUSY spectra can be generated from complete anarchy at moderate mediation scales (SUSY QCD RGE)
- surprising mass hierarchies still viable, e.g.

$$m_{\tilde{g}} = 1.3 \text{ TeV}, m_{\tilde{Q}_1} = 550 \text{ GeV}, m_{\tilde{Q}_2} = 950 \text{ GeV}$$

Important implications for LHC searches

On Universality of CPV in $SU(3)_Q$ breaking NP

- SM quark flavor
- two sources
- **Implication**
- viable SU moderate
- surprising



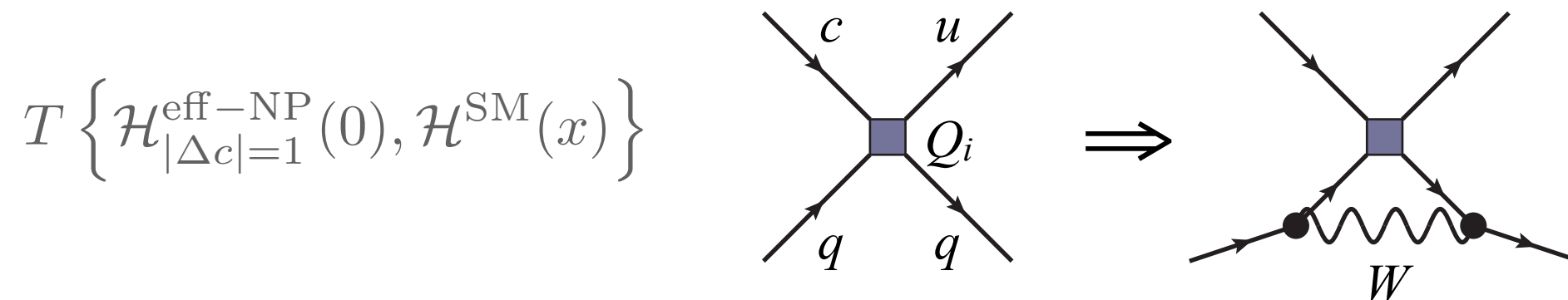
see talks by Mahbubani on Wednesday
& Perez on Friday

}

Δa_{CP} and New Physics

Isidori, J.F.K, Ligeti & Perez
1111.4987

- In EFT can be estimated via “weak mixing” of operators



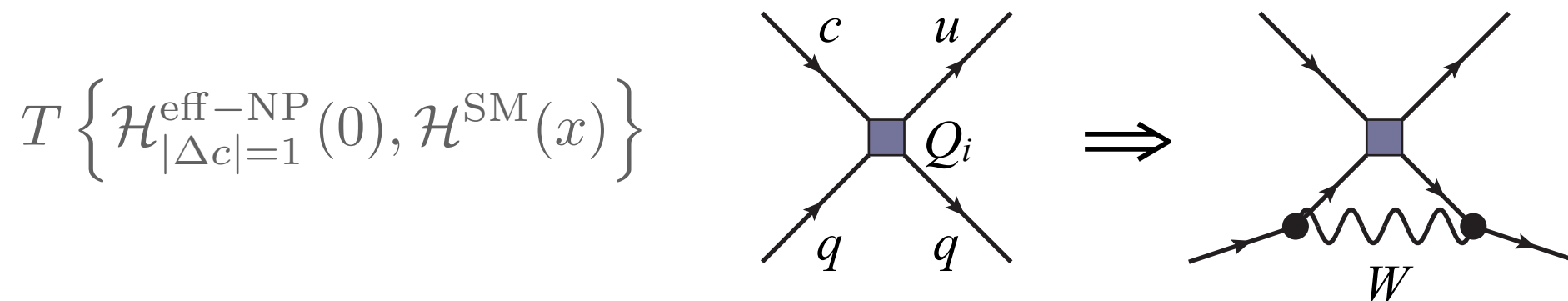
- Important constraints expected from **D- \bar{D} mixing** and direct CPV in $K^0 \rightarrow \pi^+ \pi^-$ (**ϵ'/ϵ**)
 - LL 4q operators: **excluded**
 - LR 4q operators: **ajar** - potentially visible effects in D- \bar{D} and/or ϵ'/ϵ

Model example:
Hochberg, Nir, 1112.5268

Δa_{CP} and New Physics

Isidori, J.F.K, Ligeti & Perez
1111.4987

- In EFT can be estimated via “weak mixing” of operators



- Important constraints expected from **D- \bar{D} mixing** and direct CPV in $K^0 \rightarrow \pi^+ \pi^-$ (**ϵ'/ϵ**)
 - LL 4q operators: **excluded**
 - LR 4q operators: **ajar** - potentially visible effects in D- \bar{D} and/or ϵ'/ϵ
 - RR 4q operators: unconstrained in EFT - UV sensitive contributions?

Dipole operators only weakly constrained (edm's)

Δa_{CP} in (enter favorite NP model name)

Before LHCb result,

Δa_{CP} in charm not on top of NP theorists expectations

Δa_{CP} in (enter favorite NP model name)

Before LHCb result,

Δa_{CP} in charm not on top of NP theorists expectations

In last 6 months, situation has improved considerably

Δa_{CP} in SUSY Models

Grossman, Kagan & Nir, hep-ph/0609178
Giudice, Isidori & Paradisi, 1201.6204
Hiller, Hochberg, Nir, 1204.1046

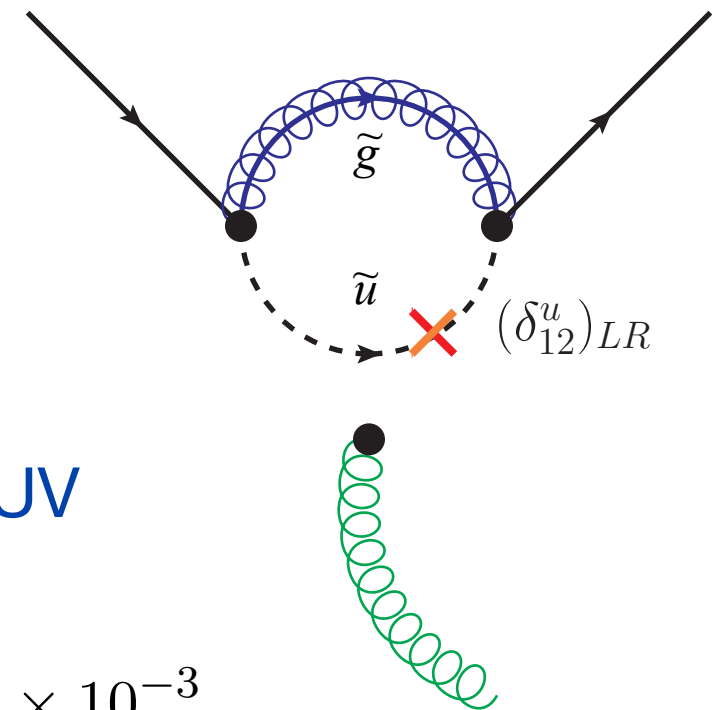
- Left-right up-type squark mixing contributions

$$|\Delta a_{CP}^{\text{SUSY}}| \approx 0.6\% \left(\frac{|\text{Im}(\delta_{12}^u)_{LR}|}{10^{-3}} \right) \left(\frac{\text{TeV}}{\tilde{m}} \right)$$

- contributions to $\Delta F=2$ helicity suppressed

- requires large trilinear (A) terms, non-trivial flavor in UV

$$\text{Im}(\delta_{12}^u)_{LR} \approx \frac{\text{Im}(A) \theta_{12} m_c}{\tilde{m}} \approx \left(\frac{\text{Im}(A)}{3} \right) \left(\frac{\theta_{12}}{0.3} \right) \left(\frac{\text{TeV}}{\tilde{m}} \right) 0.5 \times 10^{-3}$$

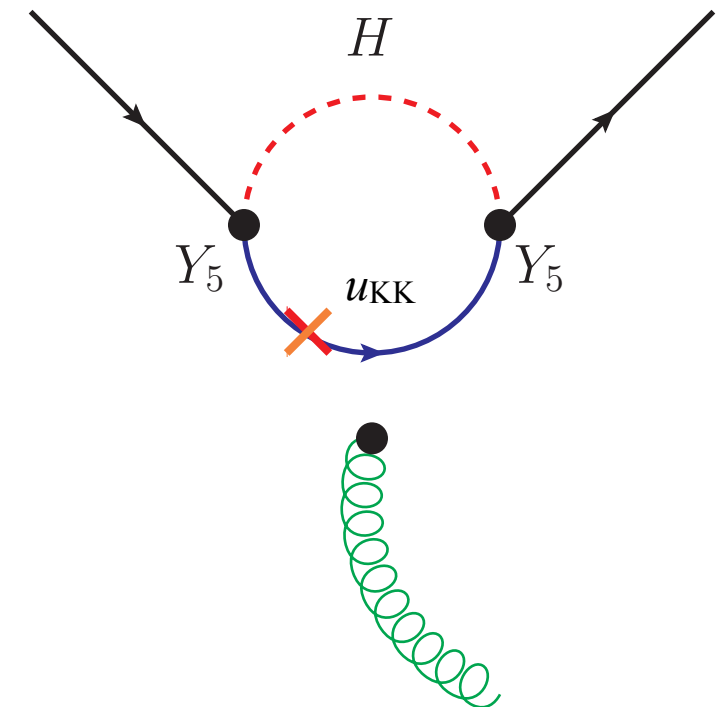


Δa_{CP} in Warped Extra-Dim. Models

- Anarchic flavor with bulk Higgs

$$|\Delta a_{CP}^{\text{chromo}}|_{\text{RS}} \simeq 0.6\% \times \left(\frac{Y_5}{6}\right)^2 \left(\frac{3 \text{ TeV}}{m_{KK}}\right)^2$$

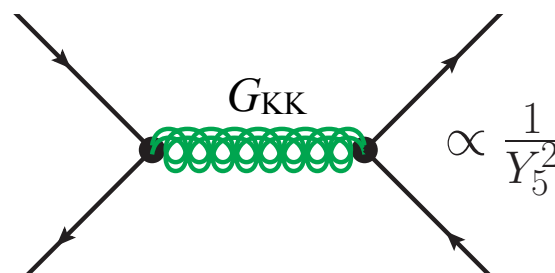
Delaunay, J.F.K., Perez & Randall
in progress



- requires very large 5D Yukawas

- helps to avoid D- \bar{D} mixing constraints

Gedalia et al., 0906.1879



- implies low UV cut-off $\frac{1}{2} \lesssim Y_5 \lesssim \frac{4\pi}{\sqrt{N_{KK}}}$

Agashe, Azatov & Zhu, 0810.1016
Csaki et al., 0907.0474

- Can be mapped to 4D partial compositeness models

Conclusions

- The observed size of CPV in charm decays is borderline
 - larger than naive SM expectations...

Conclusions

- The observed size of CPV in charm decays is borderline
 - larger than naive SM expectations...however, SM explanation cannot be excluded from first principles
- More experimental observables could clarify the picture

Grossman, Kagan & Zupan, 1204.3557
Isidori & J.F.K., 1205.3164

Conclusions

- The observed size of CPV in charm decays is borderline
 - larger than naive SM expectations...however, SM explanation cannot be excluded from first principles
 - More experimental observables could clarify the picture Grossman, Kagan & Zupan, 1204.3557
Isidori & J.F.K., 1205.3164
- If NP, points towards new flavor structures in U_R sector at the TeV scale
 - QCD dipoles most plausible candidates, can be generated in motivated UV completions
- Universality of CPV in $SU(3)_Q$ breaking NP - **only tiny effects in Δa_{CP} allowed**
 - *SUSY corollary*: quark-squark alignment models less constrained than previously thought - **important implications for LHC searches**

Backup

CPV in $|\Delta c|=2$

- **CPV in Mixing** $|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle$

$$m \equiv \frac{m_1 + m_2}{2},$$

$$x \equiv \frac{m_2 - m_1}{\Gamma},$$

$$\Gamma \equiv \frac{\Gamma_1 + \Gamma_2}{2},$$

$$y \equiv \frac{\Gamma_2 - \Gamma_1}{2\Gamma}.$$

- Experimentally accessible mixing quantities:

- x, y (CP conserving) Cannot be estimated accurately within SM
NP contributions are predictable
- flavor specific time-dependent CPV decay asymmetries [sensitive to q/p]

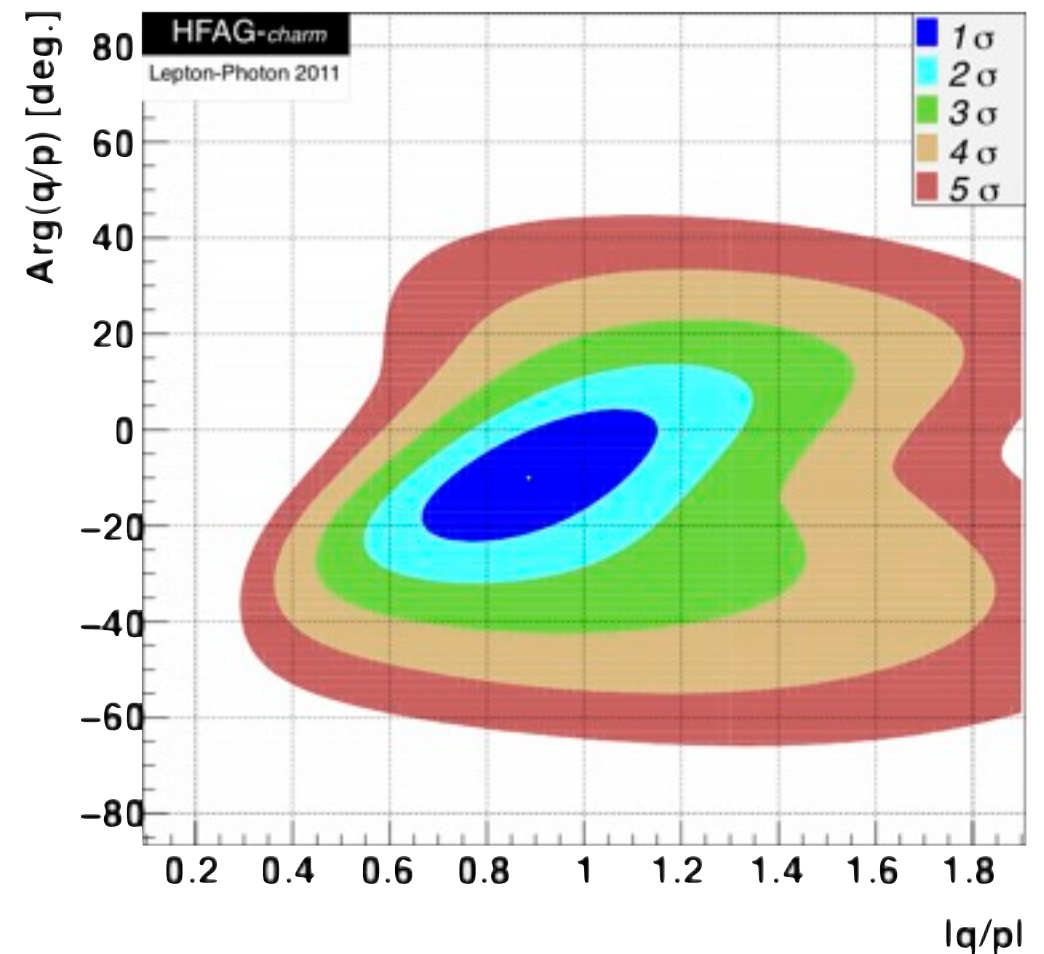
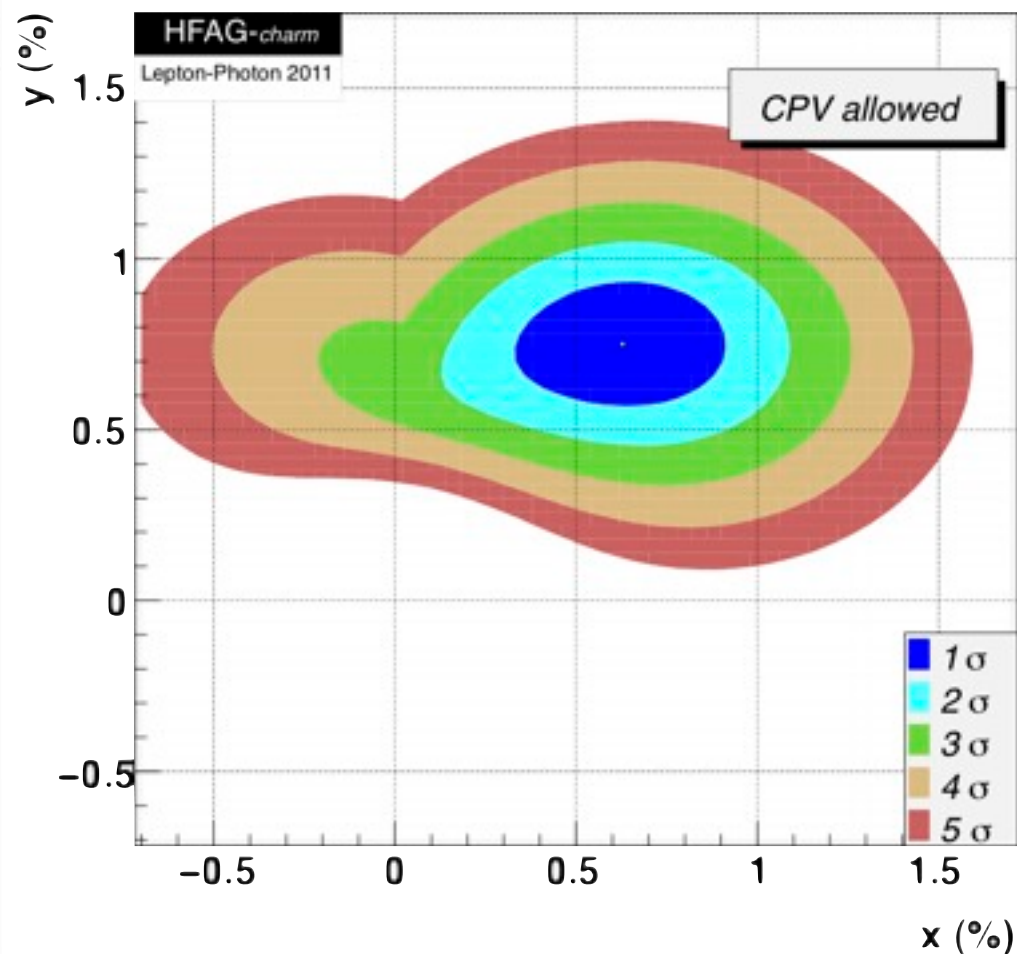
$$a_f(t) \equiv \frac{\Gamma(D^0(t) \rightarrow f) - \Gamma(\bar{D}^0(t) \rightarrow f)}{\Gamma(D^0(t) \rightarrow f) + \Gamma(\bar{D}^0(t) \rightarrow f)},$$

CPV in $|\Delta c|=2$

- **CPV in Mixing** $|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle$

$$m \equiv \frac{m_1 + m_2}{2},$$
$$x \equiv \frac{m_2 - m_1}{\Gamma},$$

$$\Gamma \equiv \frac{\Gamma_1 + \Gamma_2}{2},$$
$$y \equiv \frac{\Gamma_2 - \Gamma_1}{2\Gamma}.$$



CPV in $|\Delta c|=2$

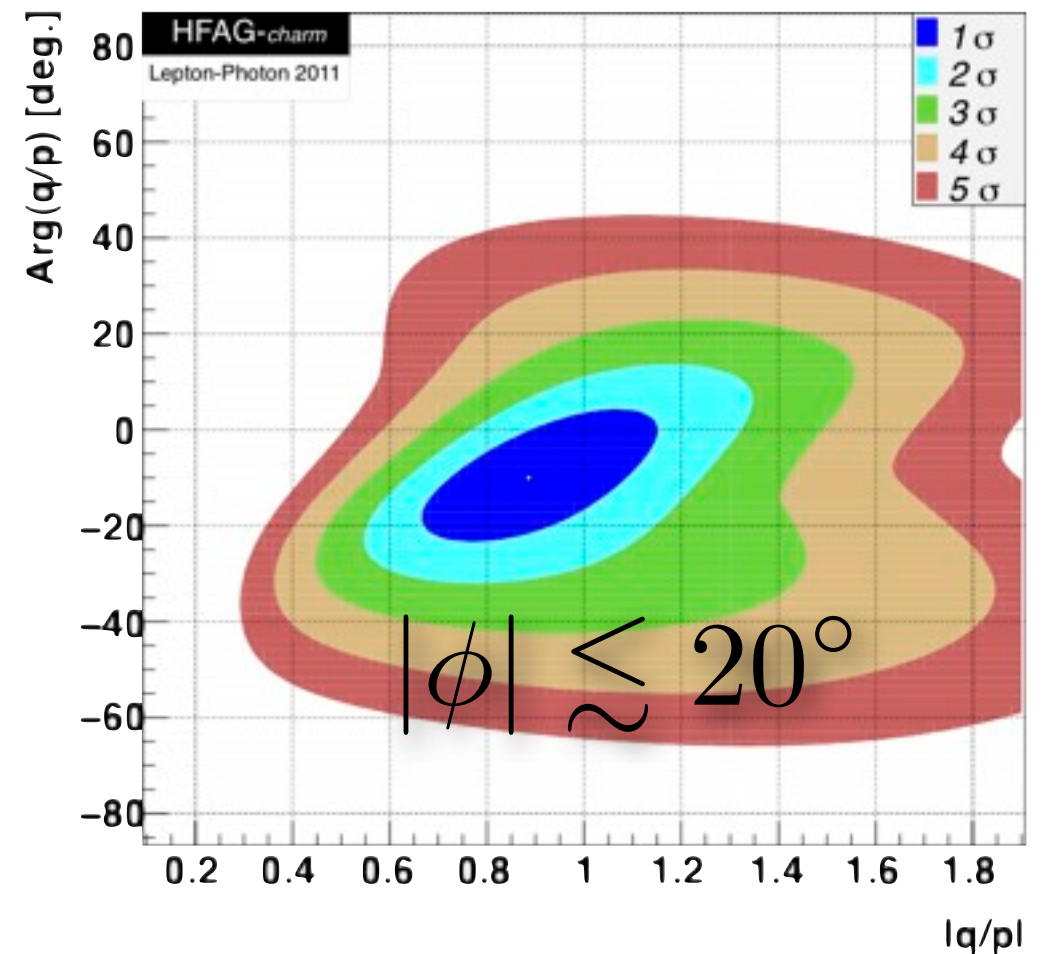
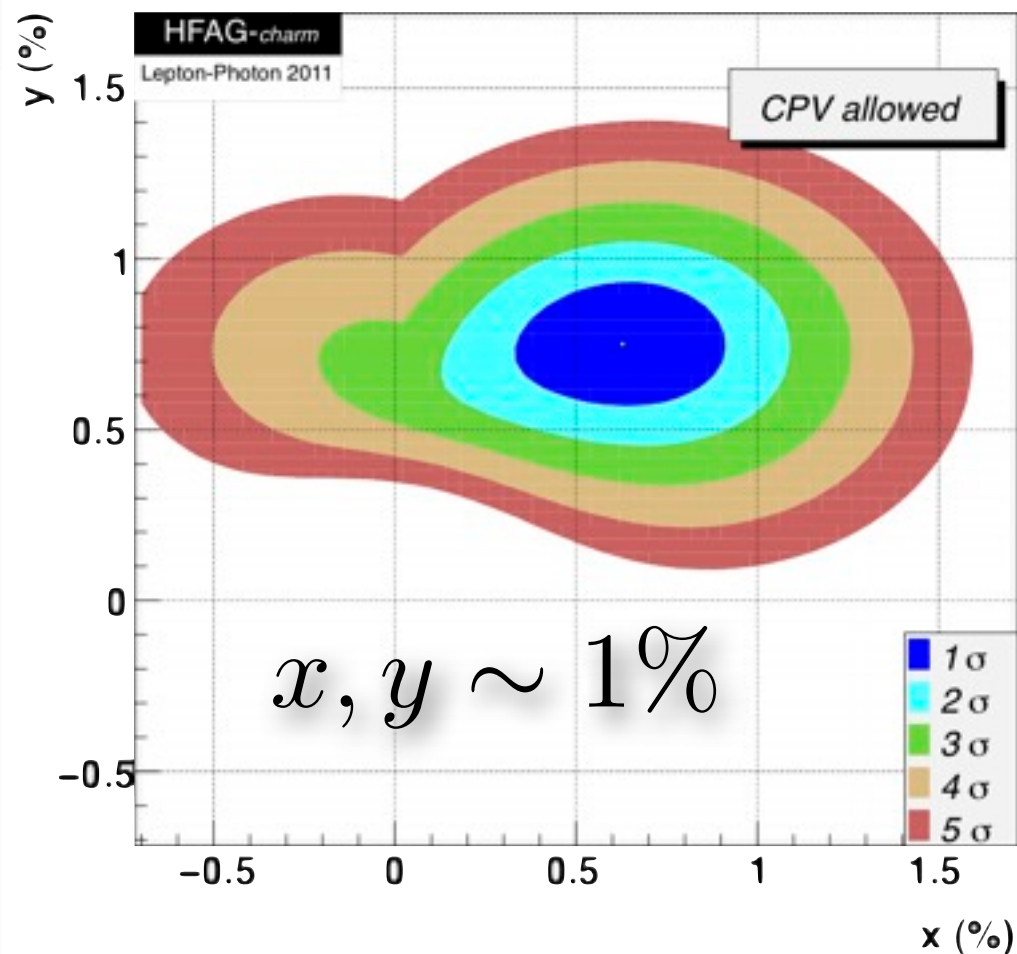
- **CPV in Mixing** $|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle$

$$m \equiv \frac{m_1 + m_2}{2},$$

$$x \equiv \frac{m_2 - m_1}{\Gamma},$$

$$\Gamma \equiv \frac{\Gamma_1 + \Gamma_2}{2},$$

$$y \equiv \frac{\Gamma_2 - \Gamma_1}{2\Gamma}.$$



CPV in $|\Delta c|=2$

- CPV in Mixing

Isidori, Nir & Perez 1002.0900

Operator	Bounds on Λ (TeV)		Bounds on c_{ij} ($\Lambda = 1$ TeV)		Observables
	Re	Im	Re	Im	
$(\bar{s}_L \gamma^\mu d_L)^2$	9.8×10^2	1.6×10^4	9.0×10^{-7}	3.4×10^{-9}	$\Delta m_K; \varepsilon_K$
$(\bar{s}_R d_L)(\bar{s}_L d_R)$	1.8×10^4	3.2×10^5	6.9×10^{-9}	2.6×10^{-11}	$\Delta m_K; \varepsilon_K$
$(\bar{c}_L \gamma^\mu u_L)^2$	1.2×10^3	2.9×10^3	5.6×10^{-7}	1.0×10^{-7}	$\Delta m_D; q/p , \phi_D$
$(\bar{c}_R u_L)(\bar{c}_L u_R)$	6.2×10^3	1.5×10^4	5.7×10^{-8}	1.1×10^{-8}	$\Delta m_D; q/p , \phi_D$
$(\bar{b}_L \gamma^\mu d_L)^2$	5.1×10^2	9.3×10^2	3.3×10^{-6}	1.0×10^{-6}	$\Delta m_{B_d}; S_{B_d \rightarrow \psi K}$
$(\bar{b}_R d_L)(\bar{b}_L d_R)$	1.9×10^3	3.6×10^3	5.6×10^{-7}	1.7×10^{-7}	$\Delta m_{B_d}; S_{B_d \rightarrow \psi K}$
$(\bar{b}_L \gamma^\mu s_L)^2$	1.1×10^2	1.1×10^2	7.6×10^{-5}	7.6×10^{-5}	Δm_{B_s}
$(\bar{b}_R s_L)(\bar{b}_L s_R)$	3.7×10^2	3.7×10^2	1.3×10^{-5}	1.3×10^{-5}	Δm_{B_s}

$$x, y \sim 1\%$$

$$|\phi| \lesssim 20^\circ$$

Imply significant constraints on CPV NP contributions, second only to kaon sector

Generic Implications for Experiment

- correlations with EDM's, rare top & down-type quark processes

Giudice, Isidori & Paradisi, 1201.6204

very model dependent

Hochberg & Nir, 1112.5268

Altmannshofer et al., 1202.2866

- NP explanations of Δa_{CP} via chromo-magnetic dipole operators
 - generically predict EM dipoles - rare radiative charm decays

$$D^0 \rightarrow X \gamma$$

$$D^0 \rightarrow X e^+ e^-$$

Delaunay, J.F.K., Perez & Randall
in progress

Expected NP rates few orders below SM LD contributions

- possibility to access CPV observables
- NP explanations of Δa_{CP} via $\Delta I=3/2$ contributions

Isidori & J.F.K., 1205.3164

- SM contributions to $A_f^{(2)}$ purely $\Delta I=1/2$

Grossman, Kagan & Zupan, 1204.3557

No CPV expected in pure $\Delta I=3/2$ decays

Generic Implications for Experiment

- NP explanations of Δa_{CP} via chromo-magnetic dipole operators

Grossman, Kagan & Nir, hep-ph/0609178

Giudice, Isidori & Paradisi, 1201.6204

$$|\Delta a_{CP}^{NP}| \approx -1.8 |\text{Im}[C_8^{NP}(m_c)]| ,$$

(estimate of matrix element in QCD fact.)

- generically predict EM dipoles

$$|\text{Im}[C_7^{NP}(m_c)]| \approx |\text{Im}[C_8^{NP}(m_c)]| \approx 0.4 \times 10^{-2} . \quad \text{(QCD RGE evolution with TeV NP)}$$

Isidori & J.F.K., 1205.3164

- possibility to access CPV observables in $D^0 \rightarrow \pi\pi\gamma, KK\gamma$

- in SM CPV expected similar as in $D^0 \rightarrow \pi\pi, KK$
- large strong phases natural for LD SM contributions

$$|a_{(\rho,\omega)\gamma}|^{\max} = 0.04(1) \left| \frac{\text{Im}[C_7(m_c)]}{0.4 \times 10^{-2}} \right| \left[\frac{10^{-5}}{\mathcal{B}(D \rightarrow (\rho,\omega)\gamma)} \right]^{1/2} \lesssim 10\% .$$

(smaller effects also in $D^0 \rightarrow KK\gamma$ with m_{KK} around Φ mass)

Generic Implications for Experiment

- NP explanations of Δa_{CP} via $\Delta I=3/2$ contributions

Grossman, Kagan & Zupan, 1204.3557

- SM contributions to $A_K^{(d)}, A_\pi^{(s)}$ purely $\Delta I=1/2$

No CPV expected in pure $\Delta I=3/2$ decays

$$\Gamma(D^+ \rightarrow \pi^+ \pi^0) - \Gamma(D^- \rightarrow \pi^- \pi^0) = 0 \quad (\text{up to small isospin breaking})$$

- nonzero difference would point towards CPV $\Delta I=3/2$ NP contributions
- decay amplitude sum-rules even in presence isospin breaking

$$\frac{1}{\sqrt{2}} |A_{\pi^+ \pi^-} - \bar{A}_{\pi^- \pi^+}| \neq |A_{\pi^0 \pi^0} - \bar{A}_{\pi^0 \pi^0}|, \quad \rightarrow \quad \text{signal of } \Delta I=3/2 \text{ CPV NP}$$

- experimentally accessible with time-dependent measurements
(also Dalitz plot analyses in $D \rightarrow 3\pi$, $D \rightarrow KK\pi$)

Δa_{CP} and 4th Generation

- 3-gen CKM non-unitarity and b' penguins

Feldmann, Nandi & Soni
1202.3795

$$\Delta a_{CP} \propto 4 \operatorname{Im} \left[\frac{\lambda_{b'}}{\lambda_d - \lambda_s} \right] \simeq \frac{2 \sin \theta_{14} \sin \theta_{24} \sin(\delta_{14} - \delta_{24})}{\sin \theta_{12}}$$

- No parametric enhancement allowed due to existing $\Delta F=2$ CPV bounds

Nandi & Soni, 1011.6091
Buras et al., 1002.2126

- Effects comparable to SM still allowed

- Similar conclusions for generic mixing with vector-like quarks

Grossman, Kagan & Nir
hep-ph/0609178
Altmannshofer et al.
1202.2866