



# Light Top Partners for the Light Higgs

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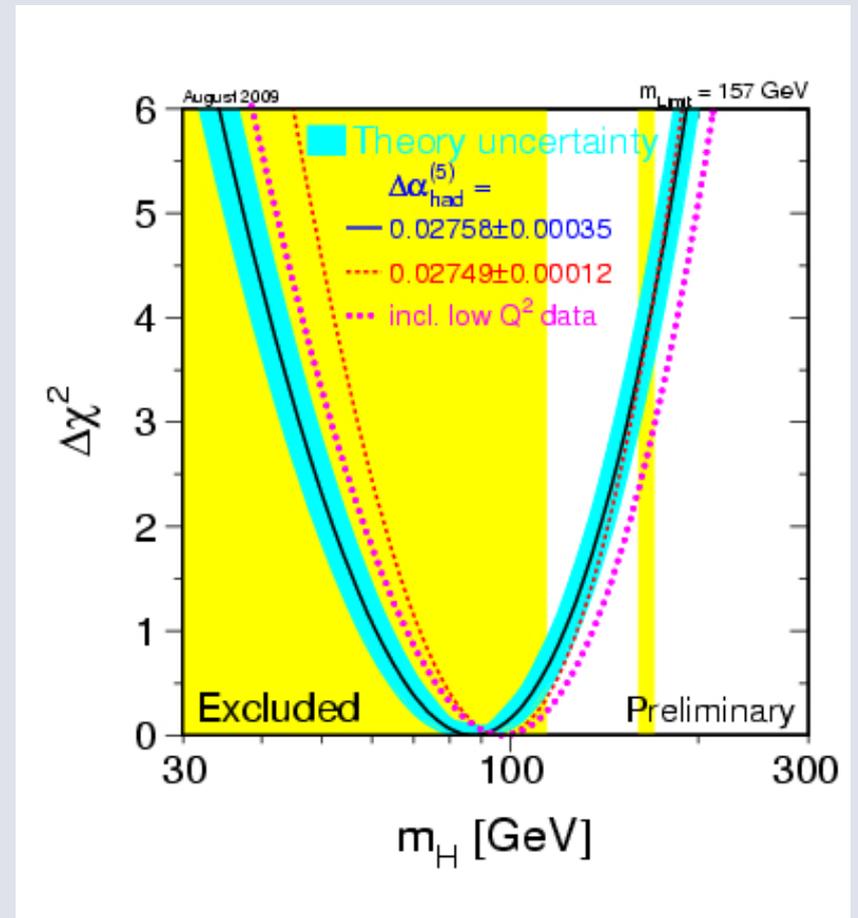
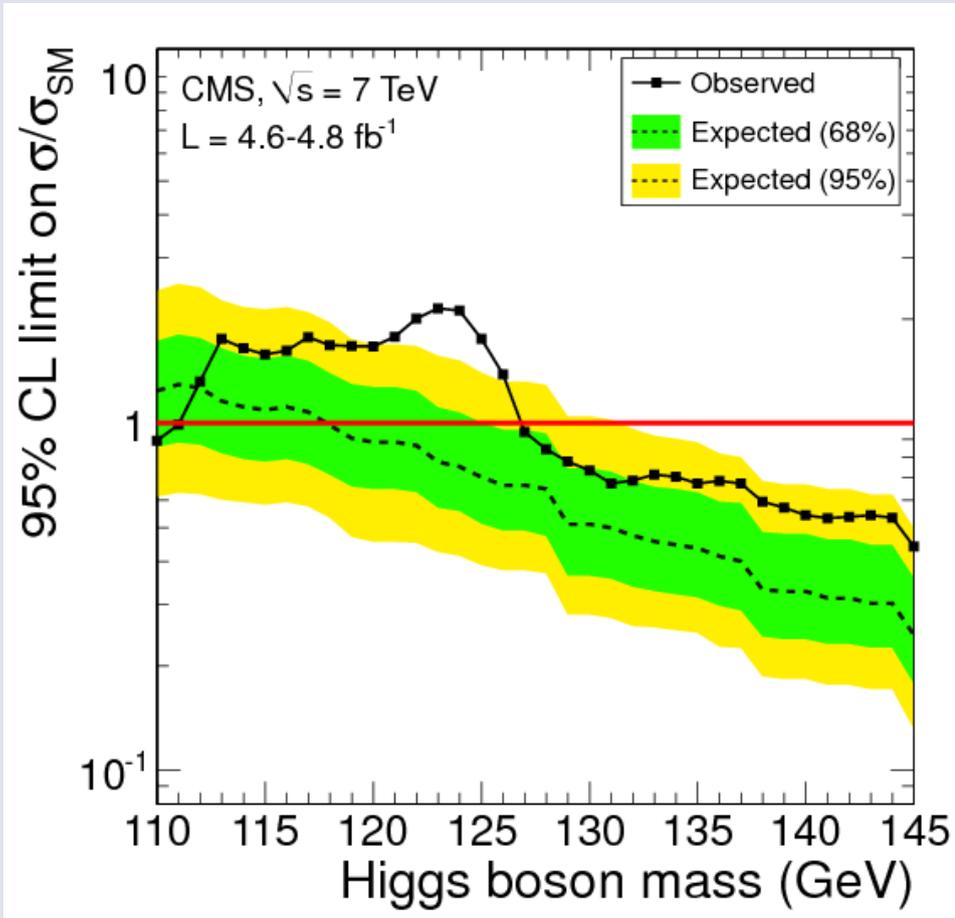
Based on the work O.M., G. Panico and A. Wulzer  
arXiv:1204.6333

# Hints of the light Higgs

direct

+

indirect



# Implications of the Light Higgs

Naturalness problem:

$$m_h^2 = m_{h0}^2 + \delta m_h^2(SM) + \delta m_h^2(UV)$$

$$\delta m_h(SM) \sim \Lambda \quad \text{- cutoff scale}$$

## Composite Higgs Scenario

- Compositeness:  
the finite size of the particles provides  
a new cutoff  $m_\rho$

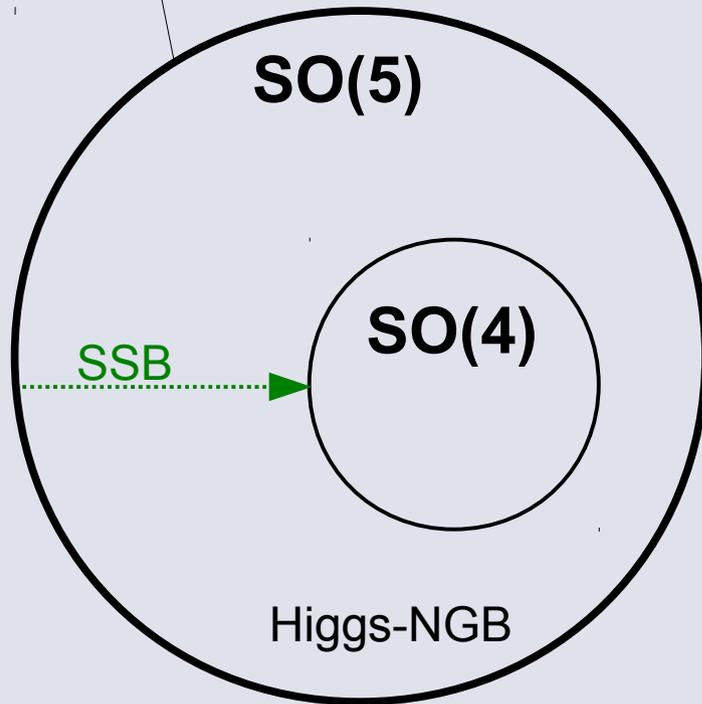
- Higgs is pNGB:  
lighter than other resonances

$$m_h \ll m_\rho$$



# General set-up

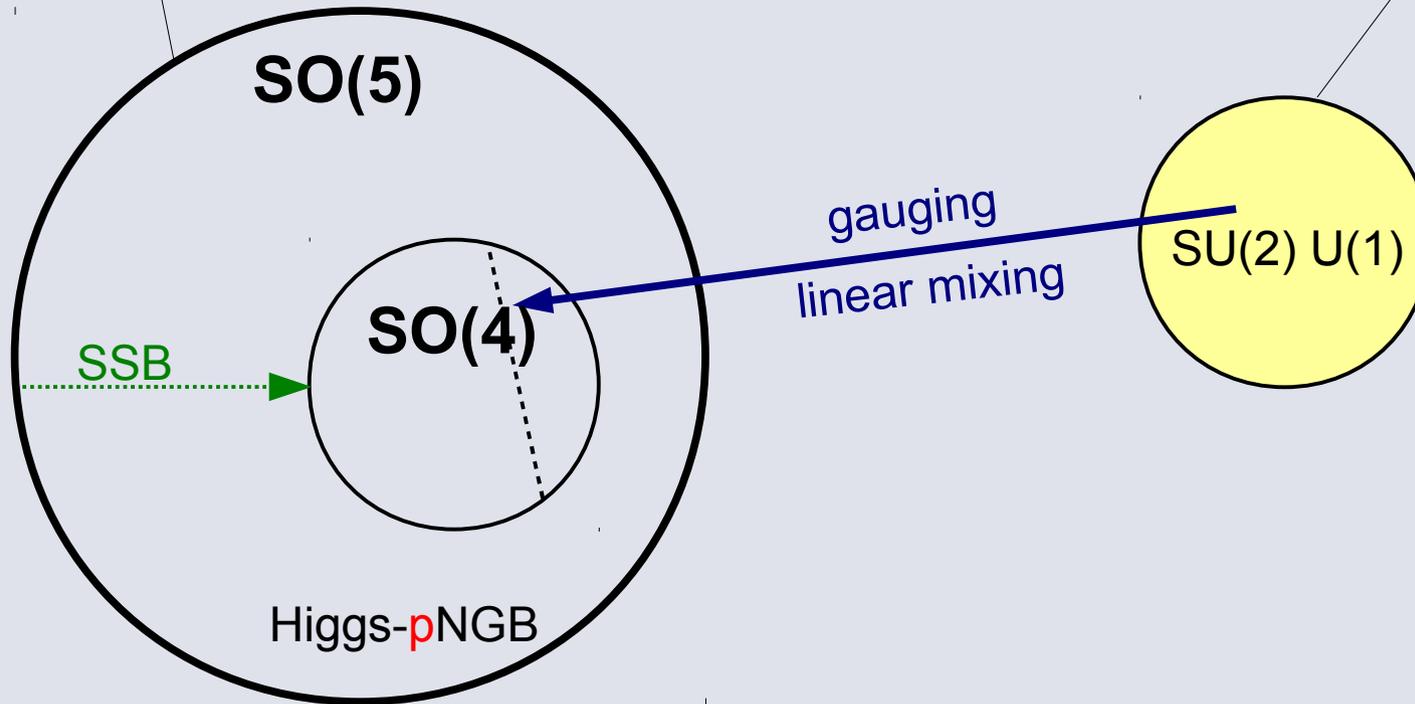
New strong sector



# General set-up

New strong sector

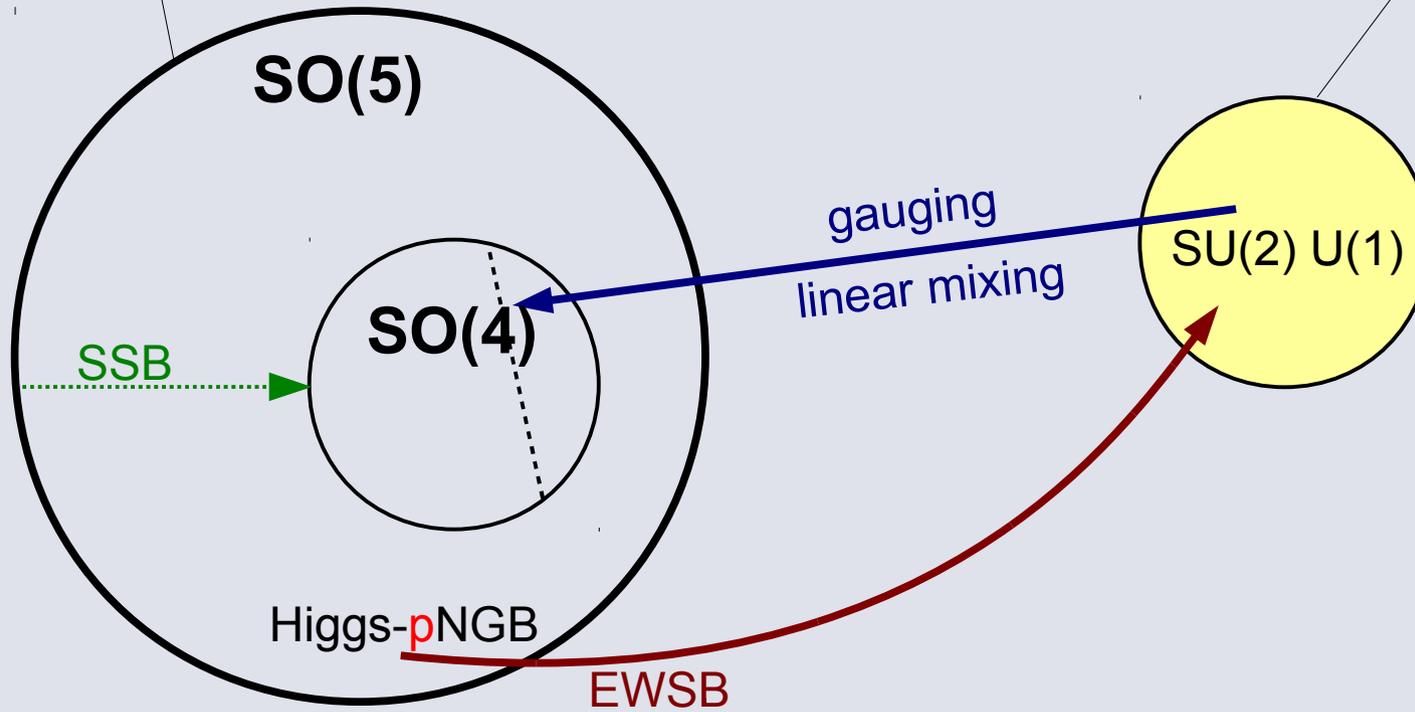
SM  
without scalars



# General set-up

New strong sector

SM  
without scalars



# General set-up

- ▶ Higgs is a NGB

$$h \rightarrow h + c$$

- ▶ Custodial symmetry + minimality

$$U = e^{i \frac{\sqrt{2}}{f_\pi} \Pi_{\hat{a}} T^{\hat{a}}} \quad T^{\hat{a}} \in SO(5)/SO(4)$$

- ▶ Partial compositeness

The diagram illustrates partial compositeness. It features a central equation:  $y_L (\bar{q}_L)^\alpha \Delta_{\alpha I}^L (\mathcal{O}_R)^I + y_R (\bar{t}_R) \Delta_I^R (\mathcal{O}_L)^I$ . Above the equation, a green double-headed arrow labeled "Composite operators" spans the width of the operator terms  $(\mathcal{O}_R)^I$  and  $(\mathcal{O}_L)^I$ . Below the equation, another green double-headed arrow labeled "Elementary quarks" spans the width of the quark terms  $(\bar{q}_L)^\alpha$  and  $(\bar{t}_R)$ . The arrows indicate the relationship between the elementary fields and the composite operators they mix with.

$$y_L (\bar{q}_L)^\alpha \Delta_{\alpha I}^L (\mathcal{O}_R)^I + y_R (\bar{t}_R) \Delta_I^R (\mathcal{O}_L)^I$$

Elementary quarks

- ▶ Higgs potential is IR saturated

# Top mass from Top Partners

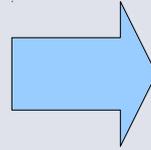
Top mass generation via interactions with the Top Partners

*linear mixing*

$$y_L f_\pi \bar{t}_L T_R + y_R f_\pi \bar{t}_R \tilde{T}_L$$

*proto-Yukawa*

$$Y_* h \bar{T} \tilde{T}$$



$$\begin{array}{ccccccc} t_L & & T_R & & \tilde{T}_L & & t_R \\ & \times & & \times & & \times & \\ & y_L & & Y_* & & y_R & \end{array}$$

# Top mass from Top Partners

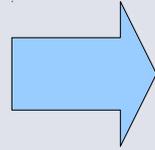
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$$y_t = Y_* \sin \varphi_L \sin \varphi_R$$

$$\sin \varphi_{L,R} = \frac{y_{L,R} f_\pi}{m_{T,\tilde{T}}}$$

degree of top compositeness

$$t_L^{\text{ph}} = \cos \varphi_L t_L + \sin \varphi_L T_L$$

# Top mass from Top Partners

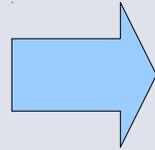
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$$y_t = Y_* \sin \varphi_L \sin \varphi_R$$

$$\sin \varphi_{L,R} = \frac{y_{L,R} f_\pi}{m_{T,\tilde{T}}}$$

the *smaller* is  $y$  the *smaller* are the masses of the top partners !

$$y^2 = \frac{y_t}{Y_*} \frac{m_T m_{\tilde{T}}}{f^2}$$

given that the source of the goldstone symmetry breaking are mixings one could expect:

$$m_h \sim y \sim m_{T,\tilde{T}}$$

# Scalar potential

Leading contribution to the potential does not lead to a successful EWSB

$$V^{(2)}(h) = \frac{N_c M_*^4}{16\pi^2 g_*^2} \left[ \frac{1}{2} c_L y_L^2 - c_R y_R^2 \right] \boxed{\sin^2(h/f_\pi)}$$

typical mass of  
fermionic  
resonances

typical coupling of  
fermionic  
resonances

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$$\text{Fine tuning: } \frac{1}{2} c_L y_L^2 = c_R y_R^2 (1 + \mathcal{O}(y^2/g_*^2))$$

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$$V^{(4)}(h) = \frac{N_c M_*^4}{16\pi^2 g_*^4} \left[ c_1^{(4)} y^4 \sin^2(h/f_\pi) + c_2^{(4)} y^4 \sin^2(h/f_\pi) \cos^2(h/f_\pi) \right]$$

$$\text{Fine tuning: } \sin^2(v/f_\pi) \ll 1$$

# Scalar potential

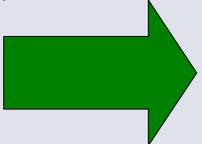
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$$m_H^2 \simeq \frac{2N_c y^4}{16\pi^2} f_\pi^2 \sin^2(2v/f_\pi)$$

# Scalar potential

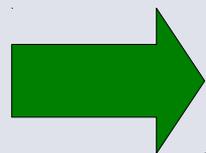
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Fine tuning:  $\frac{1}{2} c_L y_L^2 = c_R y_R^2 (1 + \mathcal{O}(y^2/g_*^2))$

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Fine tuning:  $\sin^2(v/f_\pi) \ll 1$



Source of the goldstones symmetry breaking<sup>4</sup>

$$m_H^2 \simeq \frac{2N_c y^4}{16\pi^2} f_\pi^2 \sin^2(2v/f_\pi)$$

# Light Higgs wants Light Top Partners

$$m_H \simeq 4\sqrt{3}m_t \frac{m_T m_{\tilde{T}}}{4\pi Y_* f_\pi^2}$$

more refined formula:

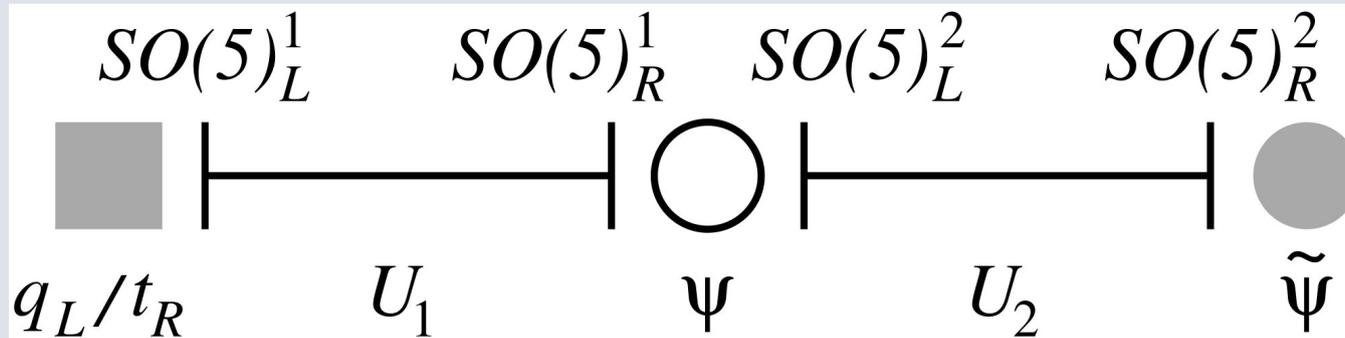
$$m_H \simeq \frac{\sqrt{N_c}}{\pi} \frac{\min(m_T, m_{\tilde{T}})}{f_\pi} m_t$$

$$f_\pi = 500 \text{ GeV}$$

$$m_f = 700 \text{ GeV}$$

accessible at the LHC !

# Example: DCHM3



- 3-site model
- Higgs boson has a triple goldstone symmetry protection
- 2 layers of composite resonances

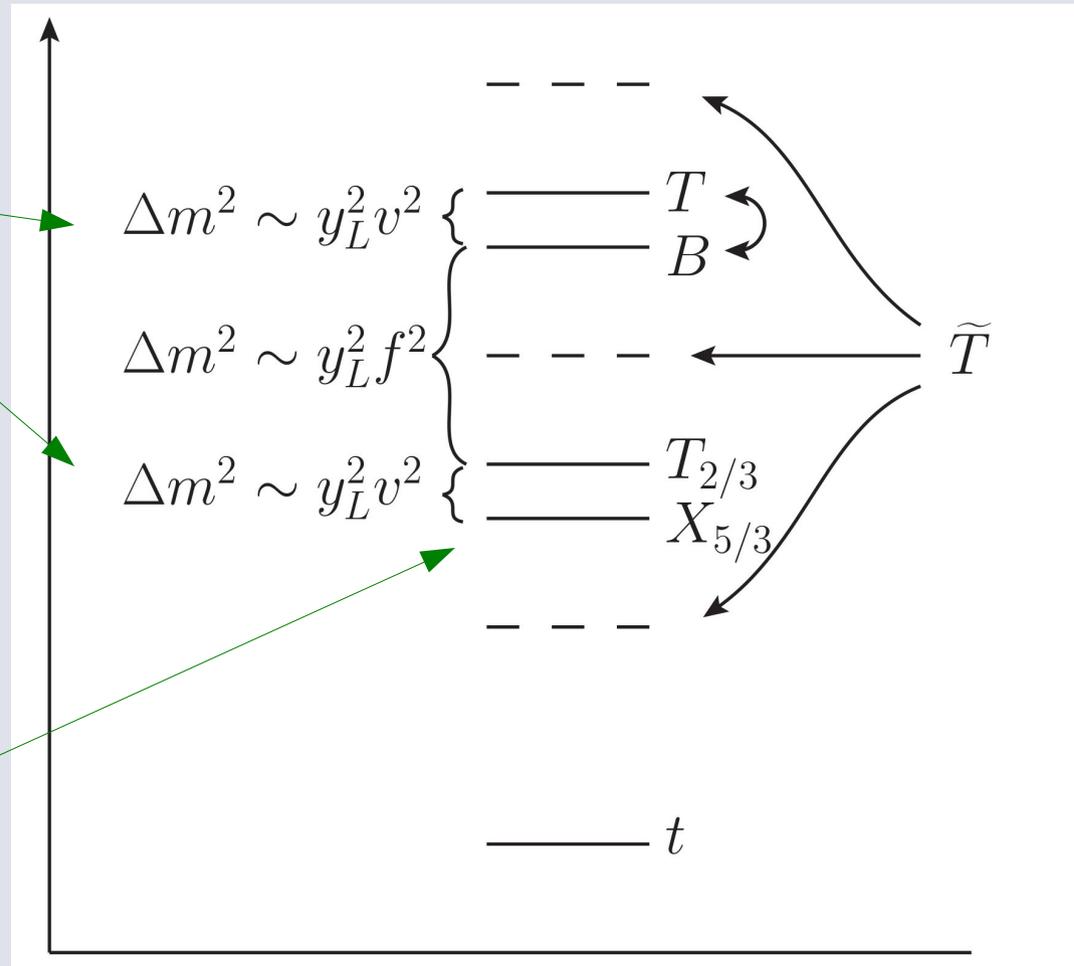
$$\psi, \tilde{\psi} \in \mathbf{5} \text{ of } SO(5)$$

- Higgs potential is calculable at one loop

# Example: DCHM3

## Simplified mass spectrum

approx. doublet structure remains after EWSB

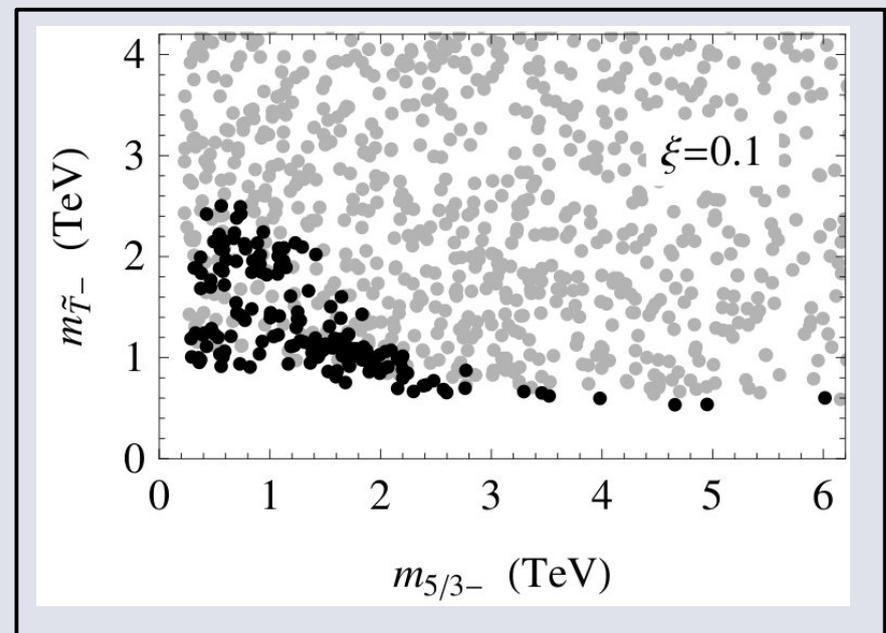
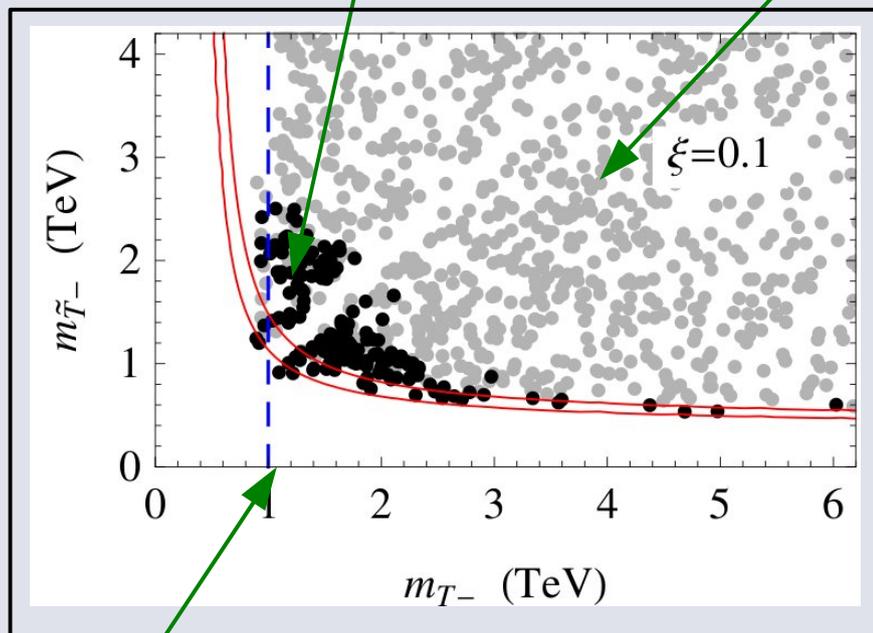


lightest state in the bidoublet has exotic charge 5/3

# Example: DCHM3

$$m_H = 115..130 \text{ GeV} \quad m_H > 130 \text{ GeV}$$

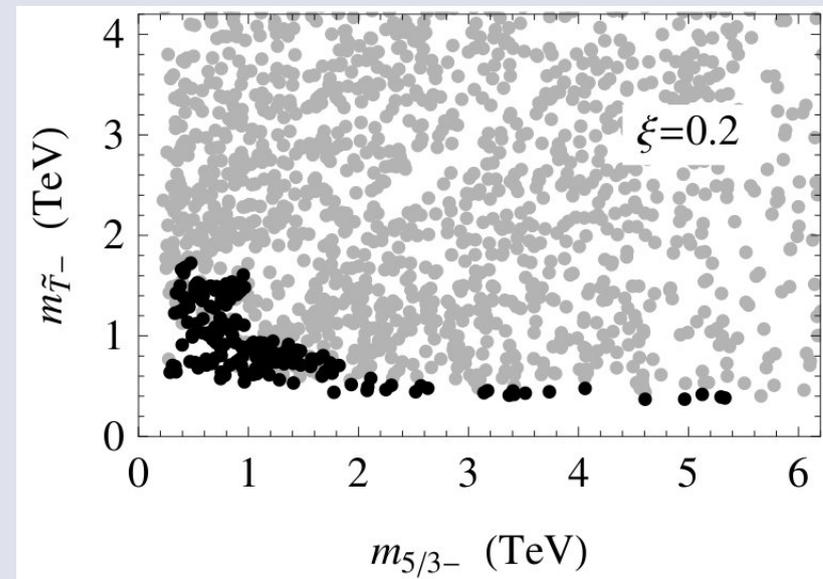
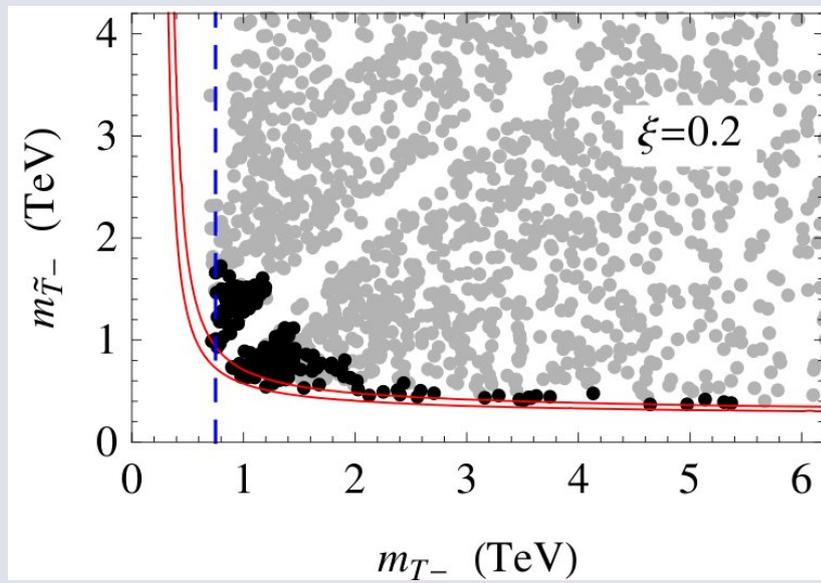
$$\xi = \left( \frac{v}{f} \right)^2$$



min mass needed  
for correct m top

Light Higgs always requires Light Partners

# Example: DCHM3

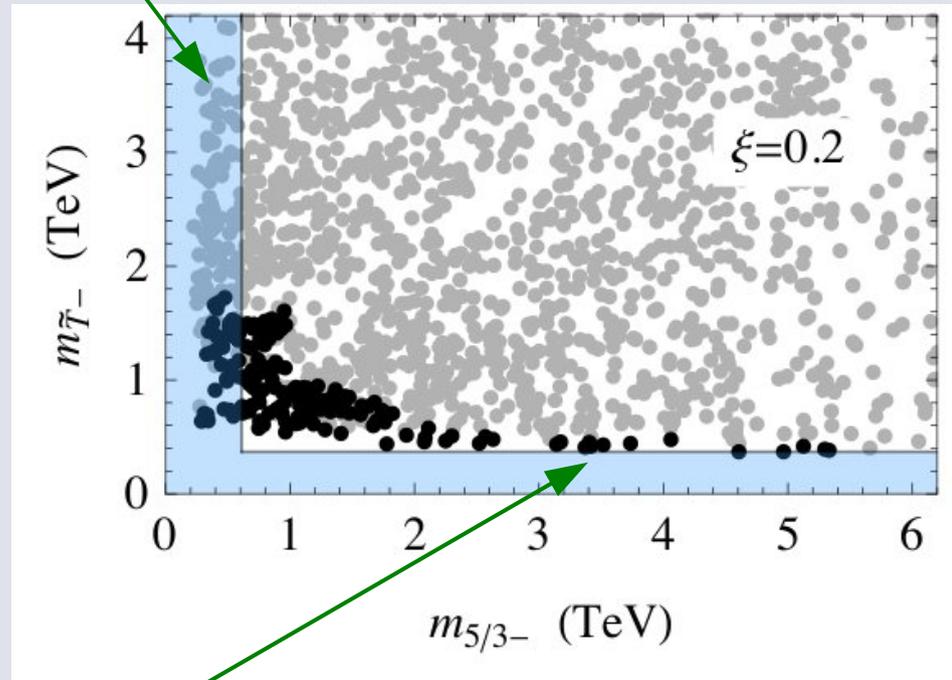


# Example: DCHM3

Bounds from direct searches

$b'b' \rightarrow WtWt \rightarrow$  same-sign dileptons and trileptons

$\rightarrow m_{5/3} > 611 \text{ GeV}$

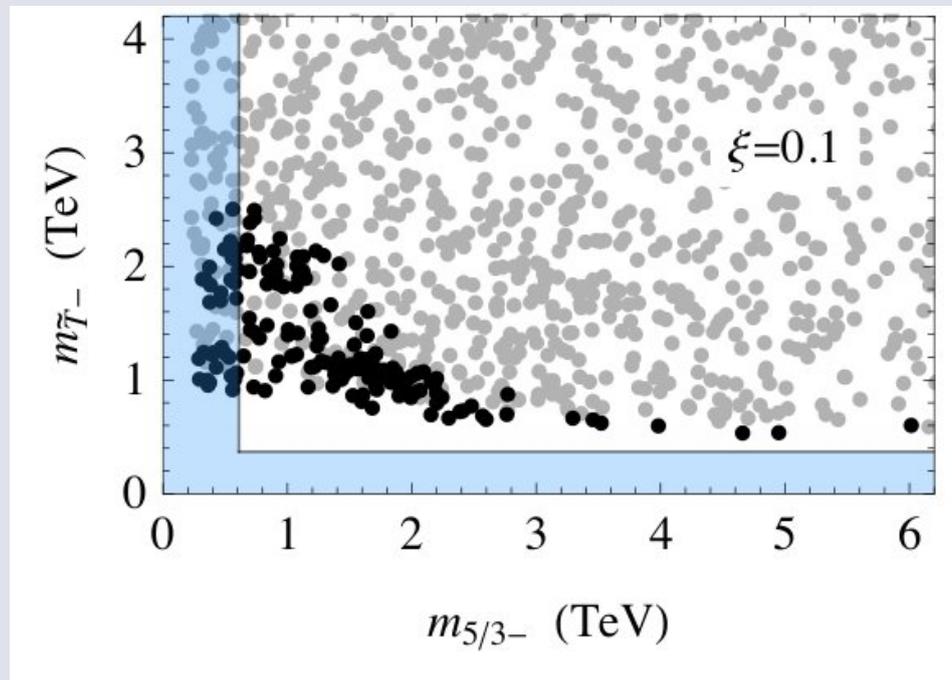


$t't' \rightarrow WbWb \rightarrow$  dileptons

$\rightarrow m_{\tilde{T}} > 370 \text{ GeV}$

# Example: DCHM3

Bounds from direct searches



# Conclusions

- In a broad class of models where the Higgs boson is realized as pNGB there is a structural correlation between the Higgs mass and the masses of composite fermionic resonances
- This correlation typically requires the presence of the light resonances below 1 TeV if the Higgs is light
- Present experimental bounds on the fermionic resonances together with the light Higgs mass put significant constraints on the allowed parameter space of considered type of models