Anwendung komplementärer Methoden, Ansätze und Higgsoff Anystossigkamansamer Ziele Warped Extra Dimension

Breiter Bereich der physikalischen Skalen; Mattonsden bei edrigsten zu den höchsten Energier Grant (EFT4LHC) Zeptometer Scale: Exploring the Origins of Johannes Gutenberg University Mainz Flavor and Electroweak Symmetry Breaking TRIGA Ionenfallen LHC **IceCube** MAMI Planck 2012 Energie [GeV] Warsaw, Poland, 31 May 2012 10<sup>-18</sup> 10<sup>-13</sup> 10<sup>0</sup> 104 1014 ultrakalte Antiprotonen Hadronen Higgs, SUSY, Medetra, de Casagrande, F. Goertz, Struktuch, Maxtrad-Dimensionen Higgs, SUSY, Neutrinos



# Higgs and flavor physics as indirect BSM probes

The hierarchy problem (mechanism of EWSB) and the origin of flavor are two big, unsolved mysteries of fundamental physics

- connected to deep questions such as the origin of mass of elementary particles, the stability of the electroweak scale, the matter-antimatter asymmetry in the Universe, the origin of fermion generations, and the reason for the hierarchies observed in the spectrum of fermion masses and mixing angles
- in SM, **flavor physics is connected to EWSB** via the Higgs Yukawa interactions

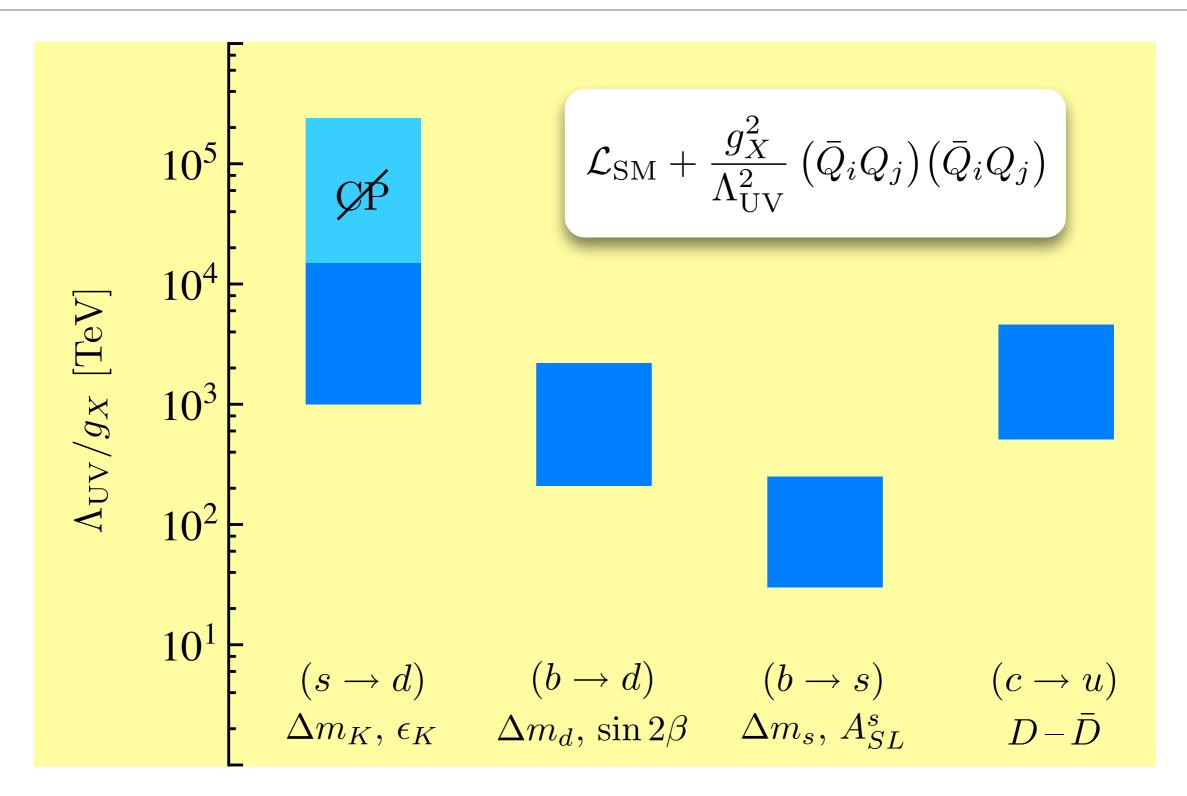
Higgs and flavor physics provide unique opportunities to probe the structure of electroweak interactions at the quantum level, thereby offering sensitive probes of physics beyond the SM

# Higgs and flavor physics as indirect BSM probes

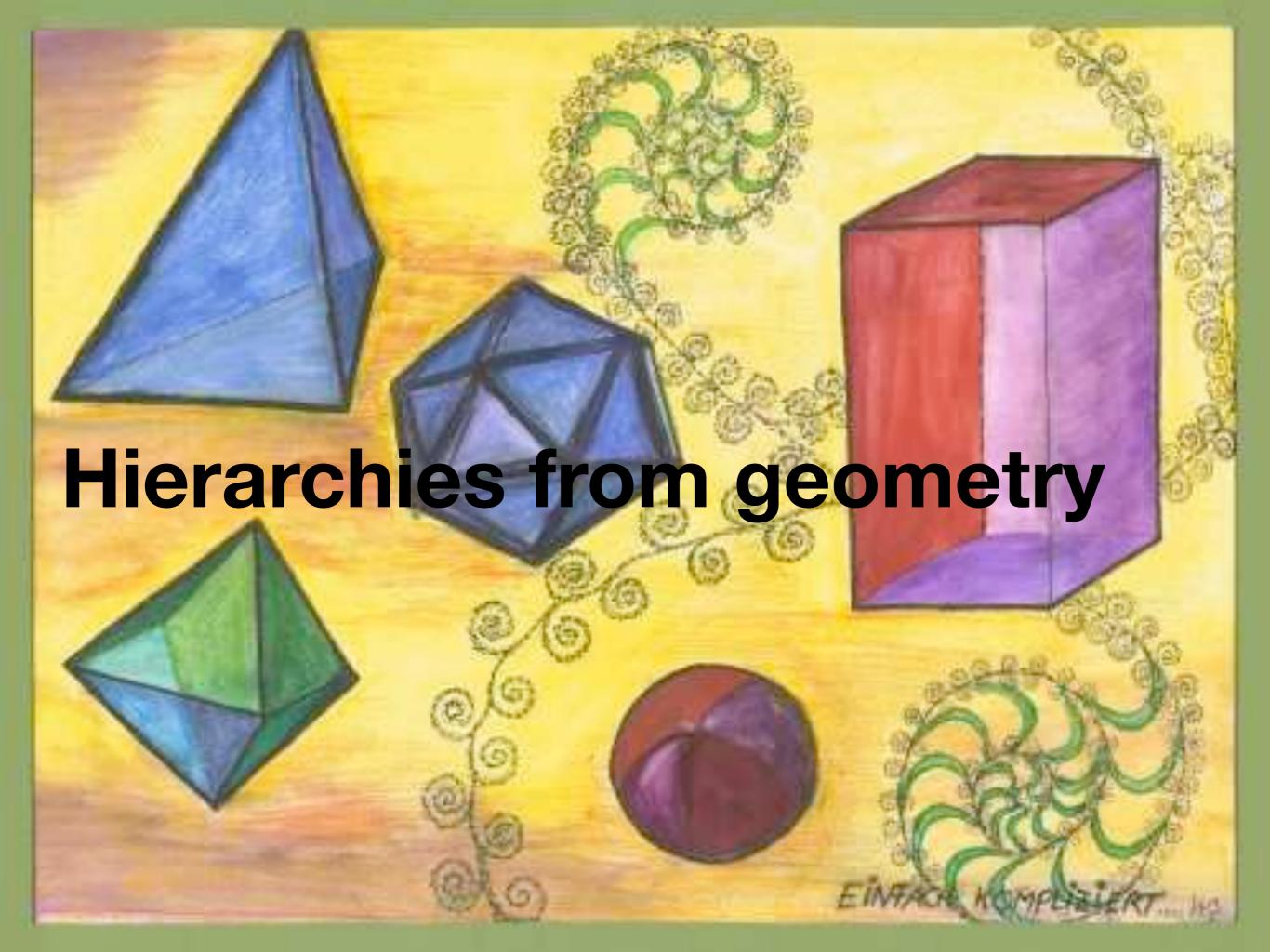
Possible solutions to flavor problem explaining  $\Lambda_{Higgs} << \Lambda_{flavor}$ :

- (i)  $\Lambda_{UV} >> 1 \text{ TeV}$ : Higgs fine tuned, new particles too heavy for LHC
- (ii)  $\Lambda_{\rm UV} \approx 1~{\rm TeV}$ : quark flavor-mixing protected by a **flavor symmetry**

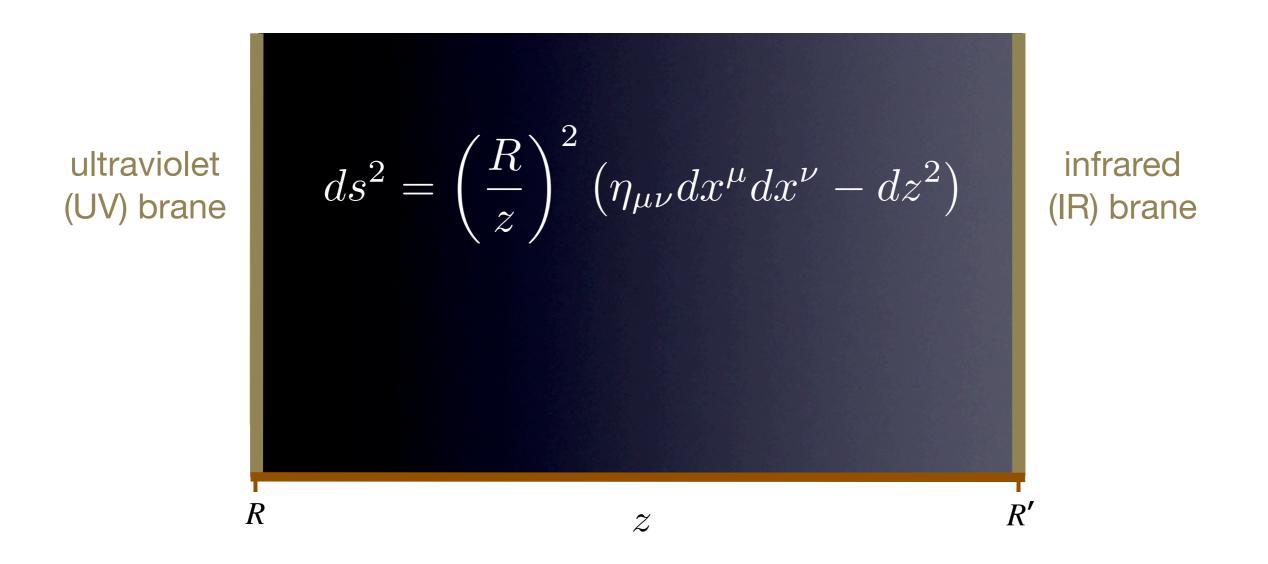
# Higgs and flavor physics as indirect BSM probes



Generic bounds on New Physics scale (for g<sub>X</sub>~1)

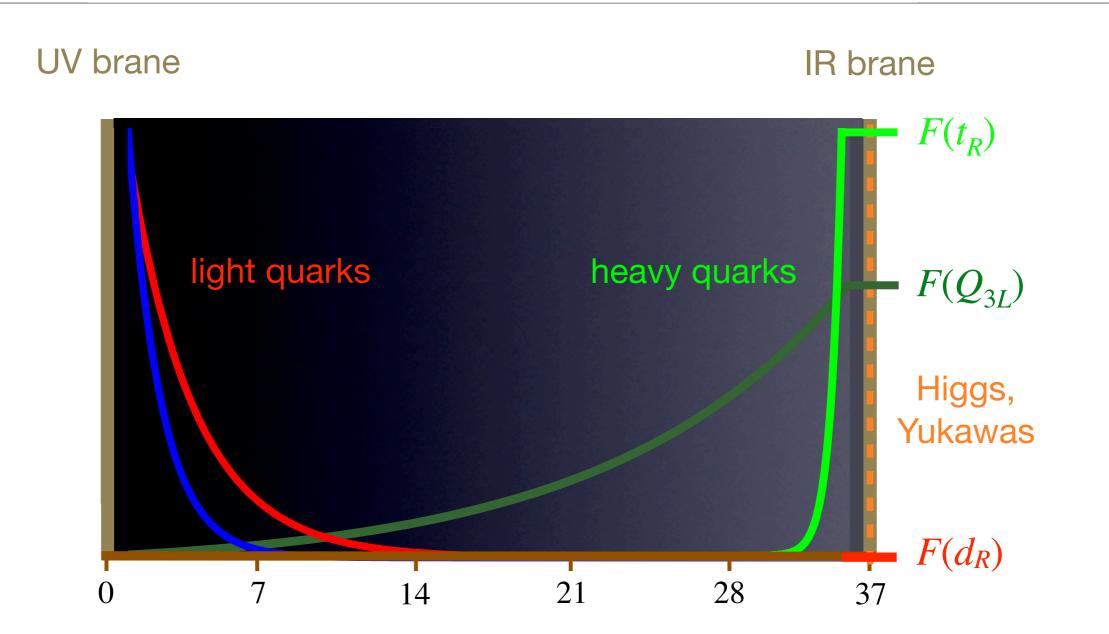


#### Flavor structure in RS models



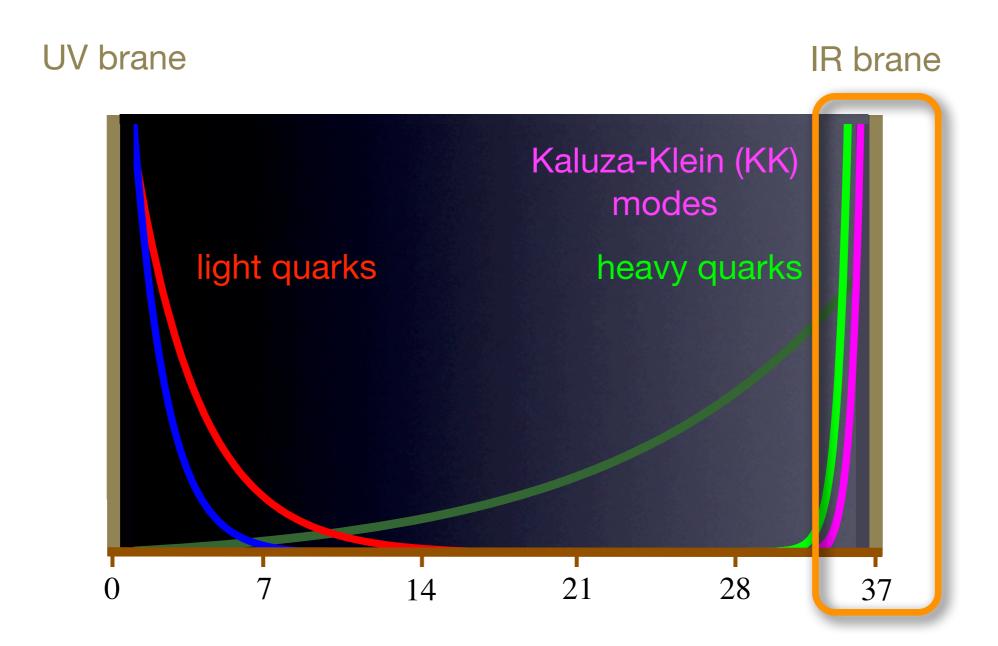
Randall-Sundrum (RS) models with a warped extra dimension address, at the same time, the **gauge hierarchy problem** and the **flavor problem** (hierarchies in the spectrum of quark masses and mixing angles)

#### Flavor structure in RS models



Localization of fermions in extra dimension depends exponentially on O(1) parameters related to the 5D **bulk masses**. Overlap integrals  $F(Q_L)$ ,  $F(q_R)$  with Higgs profile are **exponentially small** for light quarks, while O(1) for top quark: effective Yukawa couplings exhibit realistic hierarchies

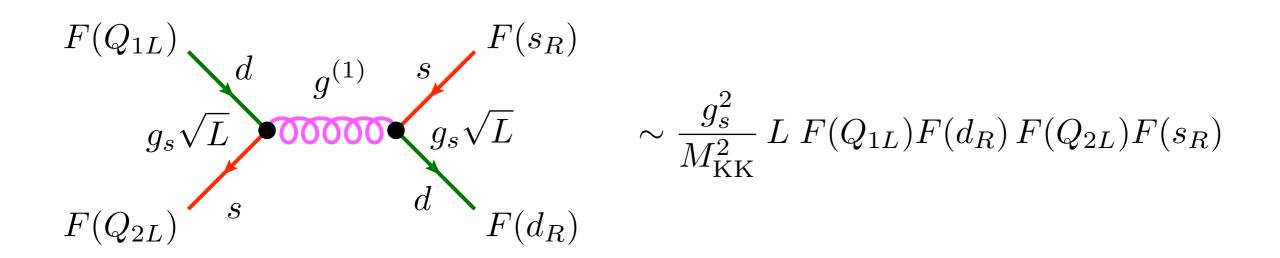
#### Flavor structure in RS models



Kaluza-Klein (KK) excitations of SM particles live close to the IR brane

Davoudiasl, Hewett, Rizzo (1999); Pomarol (1999)

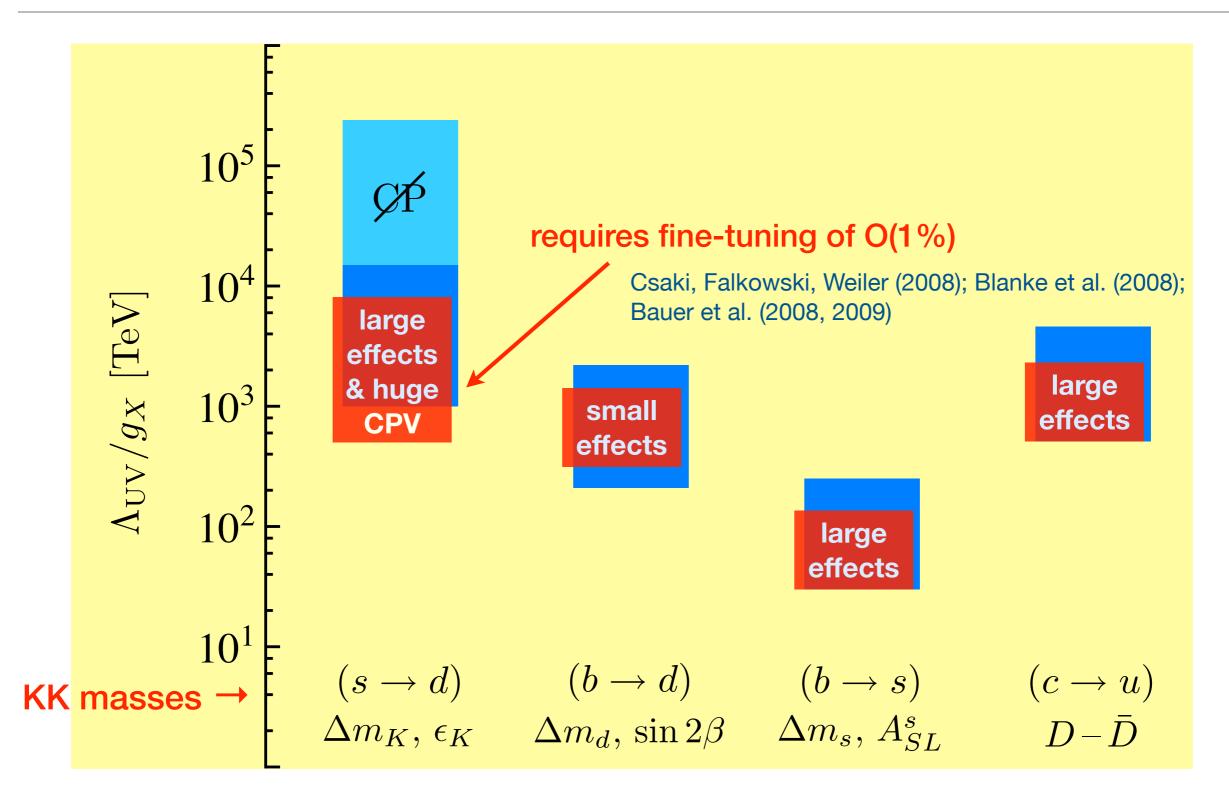
# RS-GIM protection of FCNCs



- Tree-level quark FCNCs induced by virtual exchange of Kaluza-Klein (KK)
   gauge bosons (including gluons!)
   Huber (2003); Burdman (2003); Agashe et al. (2004); Casagrande et al. (2008)
- Resulting FCNC couplings depend on same exponentially small overlap integrals  $F(Q_L)$ ,  $F(q_R)$  that generate fermion masses
- FCNCs involving light quarks are strongly suppressed: RS-GIM mechanism
   Agashe et al. (2004)

This mechanism suffices to suppress all but one of the dangerous FCNC couplings!

# RS-GIM protection of FCNCs



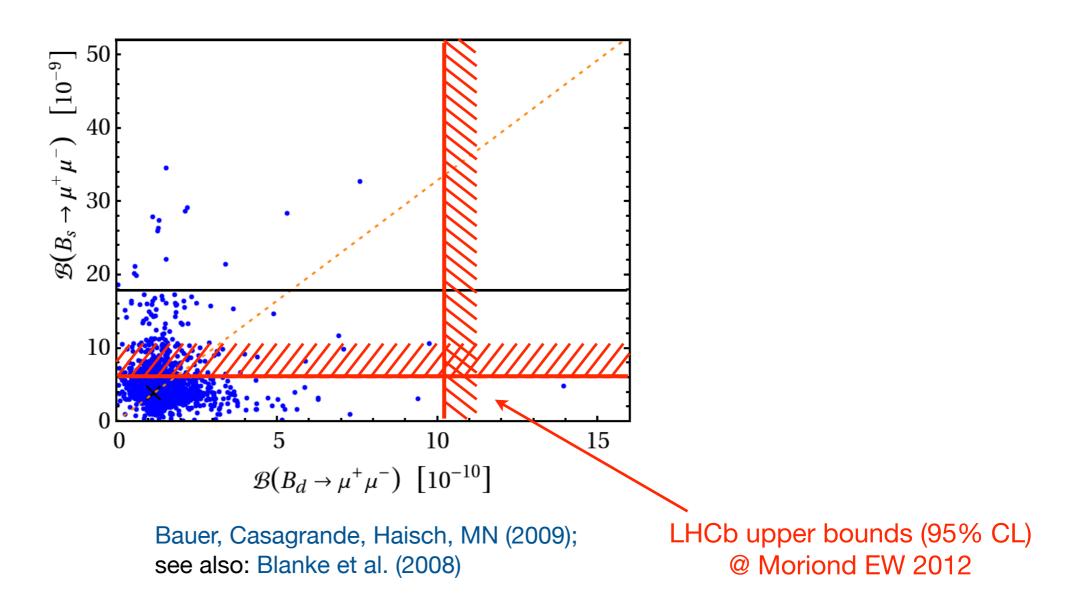
RS-GIM protection with KK masses of order few TeV

#### RS-GIM protection of FCNCs M. Bauer, R. Malm, MN arXiv: 1110.0471 (PRL) $SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_C$ broken by boundary conditions $10^5$ large effects requires fine-tuning of O(1%) or addl. model building $10^4$ Csaki, Falkowski, Weiler (2008); Blanke et al. (2008); $\Lambda_{\mathrm{UV}}/g_X \; [\mathrm{TeV}]$ Bauer et al. (2008, 2009) large $10^3$ small effects effects $10^2$ large effects $10^1$ KK masses $\rightarrow$ $(s \to d)$ $(b \to d)$ $(b \to s)$ $(c \to u)$ $\Delta m_K, \epsilon_K$ $\Delta m_d, \sin 2\beta$ $\Delta m_s, A_{SL}^s$ $D - \bar{D}$

RS-GIM protection with KK masses of order few TeV

# Example: Rare leptonic B<sub>s/d</sub>→µ<sup>+</sup>µ<sup>-</sup> decays

Rare decays  $B_{d,s} \rightarrow \mu^+ \mu^-$  could be significantly affected, but RS-GIM protection is sufficient to prevent too large deviations from SM are not generic:



Recent LHC(b) results on B<sub>s</sub>→µ<sup>+</sup>µ<sup>-</sup> begin cutting into the interesting parameter space

# Higgs Properties as an Indirect Probe for New Physics

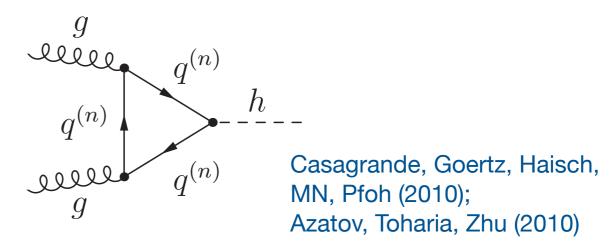


# Higgs physics as an indirect BSM probe

Higgs discovery will mark the birth of the hierarchy problem:

- one of the main motivations for physics beyond the SM
- detailed study of Higgs properties (mass, width, cross section, branching fractions) will help to probe whether the Higgs sector is as simple as predicted by the SM
- Higgs couplings to photons and gluons are loop-suppressed in the SM and hence are particularly sensitive to the presence of new particles

In RS models, large number of bulk fermionic fields in 5D theory gives rise to large loop effects, which change the effective hγγ and hgg couplings



- KK towers of light quarks contribute as much as those of heavy quarks
- effect even more pronounced in models with custodial protection

Much like flavor physics, precision Higgs physics probes quantum effects of new particles!

# Higgs physics as an indirect BSM probe

- Have recently obtained closed, analytic expressions for the Higgs production cross section as well as the main Higgs decay rates in minimal RS model, obtained by summing the infinite towers of KK states
- Work with a IR brane-localized Higgs sector, but regularize the Higgs profile carefully in an intermediate step:
   Azatov, Toharia, Zhu (2009)

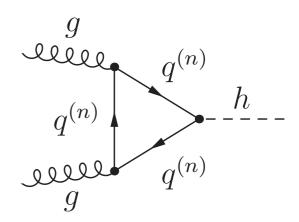


• Despite of the fact that the **sum over KK modes converges** (and is well approximated by the sum over the first few modes), it is important to implement the (physical) **UV cutoff** inherent in RS models

# Higgs physics as an indirect BSM probe

RS model is an effective theory defined with a **physical**, **5D** position-dependent cutoff, the warped Planck scale:  $\Lambda_{\rm UV} \sim M_{\rm Pl} \, e^{-\sigma(\phi)}$ 

 for loop graphs including a Higgs boson, the cutoff lies in the several TeV range



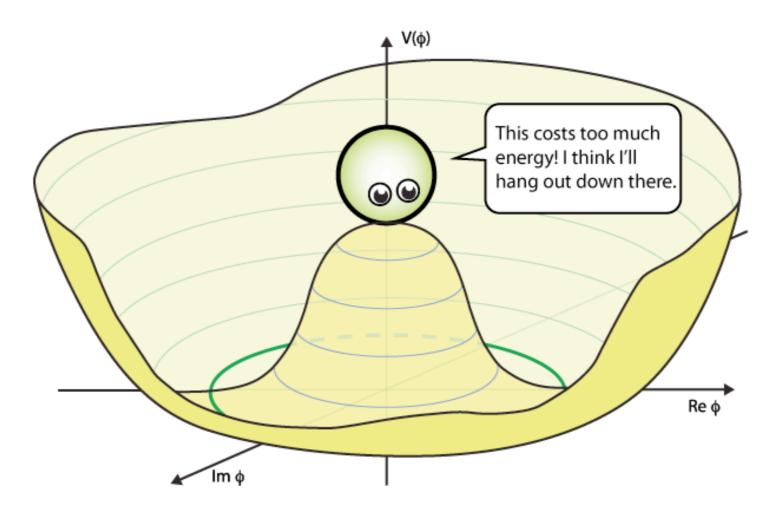
#### **UV regularization** needed for:

- gauge invariance (even in 4D)
- regularization of KK sum (superficially logarithmic divergence)
- eliminating UV contributions from super-massive KK states ~Mκκ/η

Have explored dimensional regularization & hard momentum cutoff, both in 4D picture (KK sum) and 5D loop calculation

• find same results in all cases! Carena, Casagrande, Goertz, Haisch, MN: arXiv:1204.0008 & work in preparation

# Phenomenology

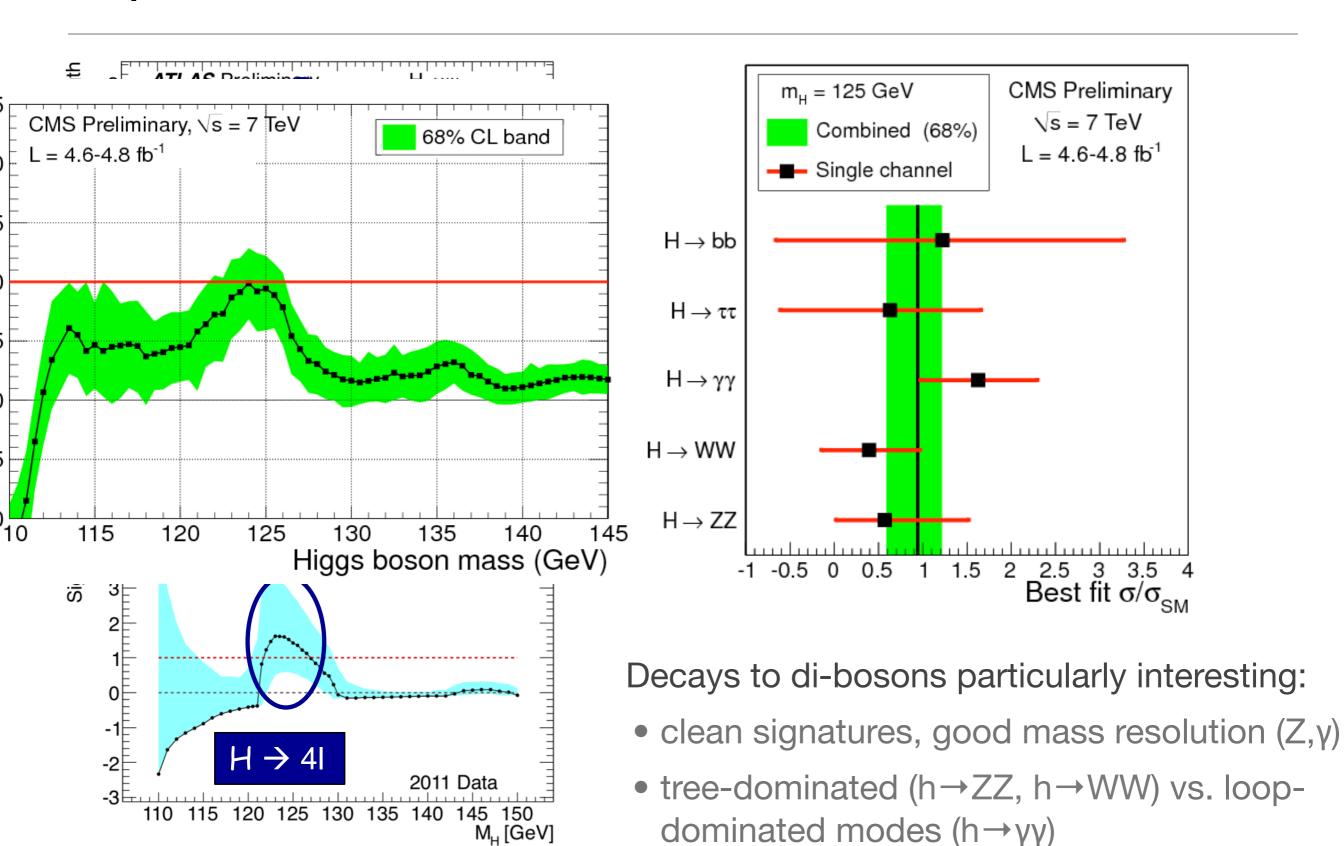


Goertz, Haisch, MN: arXiv:1112.5099

Carena, Casagrande, Goertz, Haisch, MN: arXiv:1204.0008



# Implications of recent LHC data



#### Higgs production cross section:

$$R_h = \frac{\sigma(gg \to h)_{RS}}{\sigma(gg \to h)_{SM}} = \kappa_g^2$$

Modifications due to top Yukawa and Kaluza-Klein tower of quark states:

$$\kappa_g \approx \operatorname{Re}(\kappa_t) \bigcirc \sum_{q=u,d} \operatorname{Tr} f(\boldsymbol{X}_q)$$

#### modified top contribution: reduced!

$$\kappa_t \approx 1 - \frac{v_{\mathrm{RS}}^2}{3M_{\mathrm{KK}}^2} \frac{\left(\boldsymbol{Y}_u \boldsymbol{Y}_u^{\dagger} \boldsymbol{Y}_u\right)_{33}}{\left(\boldsymbol{Y}_u\right)_{33}}$$

dominant effect described by 5D Yukawa matrices!

Higgs production cross section:

$$R_h = \frac{\sigma(gg \to h)_{RS}}{\sigma(gg \to h)_{SM}} = \kappa_g^2$$

Modifications due to top Yukawa and Kaluza-Klein tower of quark states:

 $\kappa_g \approx \operatorname{Re}(\kappa_t) \bigcirc \sum \operatorname{Tr} f(\boldsymbol{X}_q)$ 

q=u,d

modified top contribution: reduced!

$$\kappa_t \approx 1 - \frac{v_{\mathrm{RS}}^2}{3M_{\mathrm{KK}}^2} \, \frac{\left(\boldsymbol{Y}_u \boldsymbol{Y}_u^{\dagger} \boldsymbol{Y}_u\right)_{33}}{\left(\boldsymbol{Y}_u\right)_{33}}$$

dominant effect described by 5D Yukawa matrices! effect of KK tower: destructive!

$$f(x) = \frac{x \tanh x}{\cosh 2x}$$

$$\boldsymbol{X}_q = \frac{v_{\mathrm{RS}}}{\sqrt{2}M_{\mathrm{KK}}} \sqrt{\boldsymbol{Y}_q \, \boldsymbol{Y}_q^{\dagger}}$$

anarchical 5D Yukawa matrices

→ equal contributions for KK modes of all SM fermions!

#### Higgs production cross section:

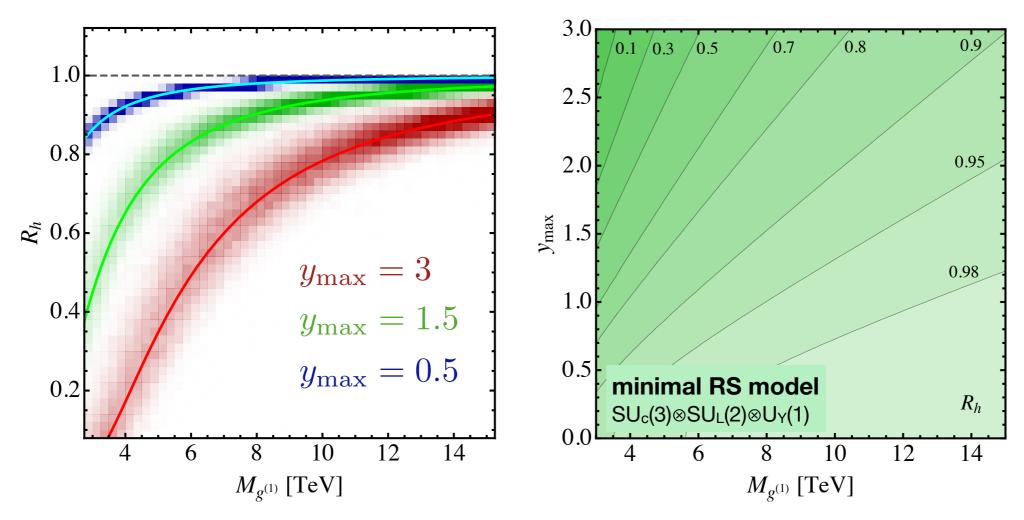
$$R_h = \frac{\sigma(gg \to h)_{RS}}{\sigma(gg \to h)_{SM}} = \kappa_g^2$$

Modifications due to top Yukawa and Kaluza-Klein tower of quark states:

$$\kappa_g \approx \operatorname{Re}(\kappa_t) \bigcirc \sum_{q=u,d} \operatorname{Tr} f(\boldsymbol{X}_q)$$

- Both effects described in terms of fundamental 5D parameters
- For random Yukawa matrices, the matrix-valued functions are to a good approximation determined by the maximum value allowed for the random variables  $0 < |(Y_q)_{ij}| < y_{max}$  (most previous analyses assumed  $y_{max}=3$ )
- Results can therefore be displayed as functions of  $y_{\rm max}$  and the mass of the lightest Kaluza-Klein gluon ( $m_{g^{(1)}}\approx 2.45M_{\rm KK}$ )

#### Find a **significant suppression** of the cross section:



- If a strong suppression is observed, then this can be interpreted as a **hint for** existence of WEDs and translated into parameter space of such models
- If σ(gg→h) is close to SM prediction, this would imply tight bounds on model parameters, perhaps moving KK masses out of LHC reach for direct production

# Higgs phenomenology in RS: decays

#### Higgs decays:

$$\frac{\Gamma(h \to VV)_{\rm RS}}{\Gamma(h \to VV)_{\rm SM}} = \kappa_V^2, \qquad \frac{\Gamma(h \to \gamma\gamma)_{\rm RS}}{\Gamma(h \to \gamma\gamma)_{\rm SM}} = \kappa_\gamma^2$$

Modification for massive vector bosons (V = Z,W) is a small effect:

$$\kappa_V \approx 1 - \frac{m_V^2}{M_{\mathrm{KK}}^2} \left( \stackrel{\longleftarrow}{L} - 1 \right) \qquad L = \ln \frac{M_{\mathrm{Pl}}}{10 \, \mathrm{TeV}} \approx 35$$

• Modification for photons can be significant (A<sub>W</sub> = 6.27 is a loop function):

$$\kappa_{\gamma} \approx \frac{1}{A_W - \frac{4N_c}{9}} \left[ \kappa_W A_W + \frac{21}{8} \, \frac{m_W^2}{M_{\rm KK}^2} \, (L-1) \right] \qquad \text{effect of KK modes W}^{(n)}$$

effect of SM quarks and KK fermions (quarks and leptons)

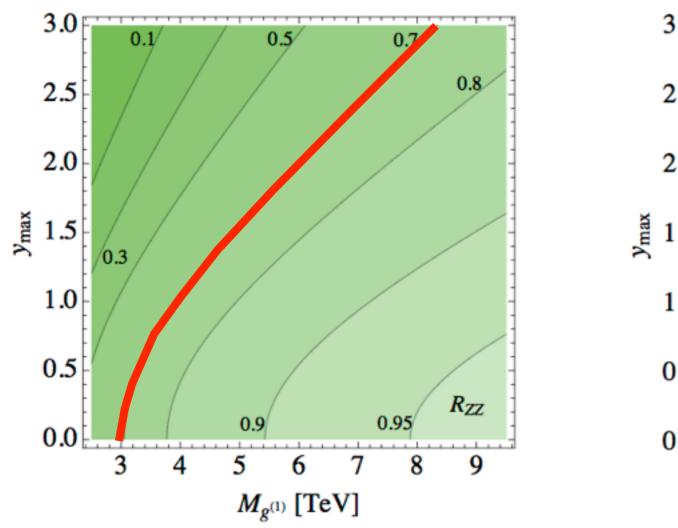
Production cross section times branching fraction:

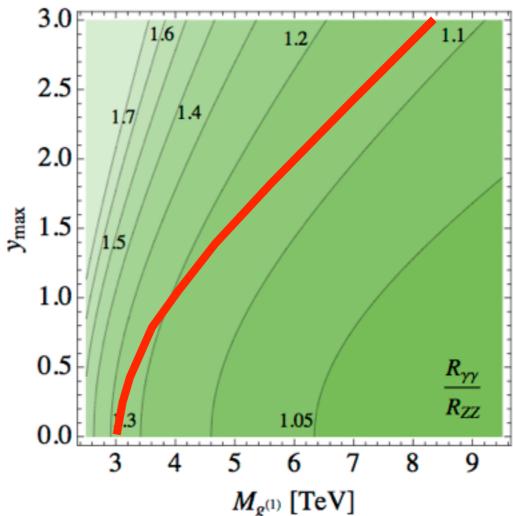
$$R_f = \frac{\left[\sigma(pp \to h) \operatorname{Br}(h \to f)\right]_{RS}}{\left[\sigma(pp \to h) \operatorname{Br}(h \to f)\right]_{SM}}$$

- Phenomenologically, the most interesting ratios are  $R_{\gamma\gamma},\,R_{ZZ}$  and their ratio (experimentally clean signatures, good mass resolution)
- ZZ mode probes loop effects via Higgs production, while di-photon mode is sensitive to loop effects in both production and decay
- Preliminary LHC data (if interpreted invoking a Higgs hypothesis) may indicate that  $R_{ZZ}\lesssim 1$ , while  $R_{\gamma\gamma}/R_{ZZ}>1$

Caveat: Discussion below is illustrative -- serious analysis must await Higgs discovery and reliable rate measurements!

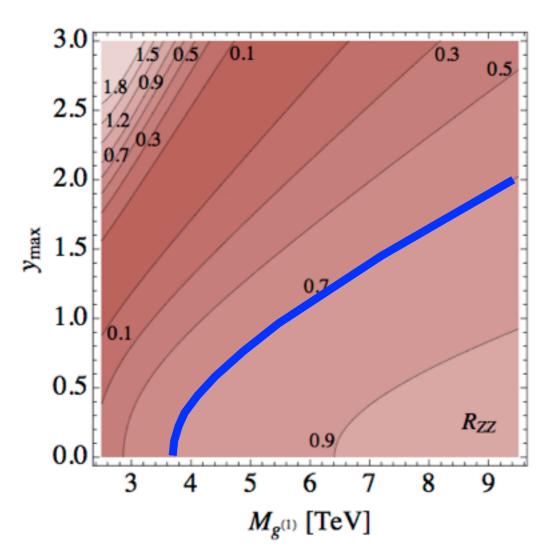
Minimal RS model with bulk fermions and gauge fields:

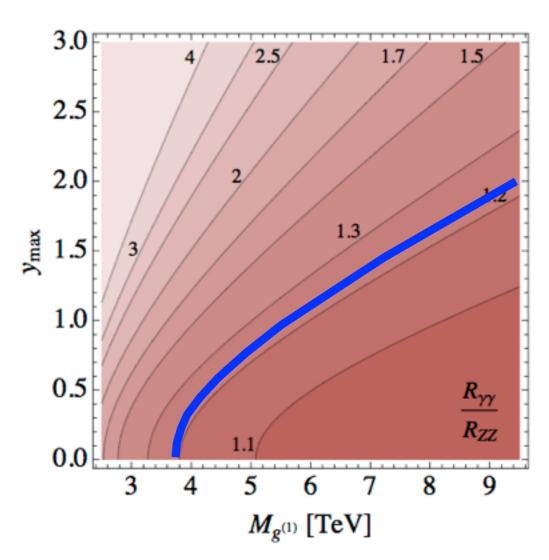




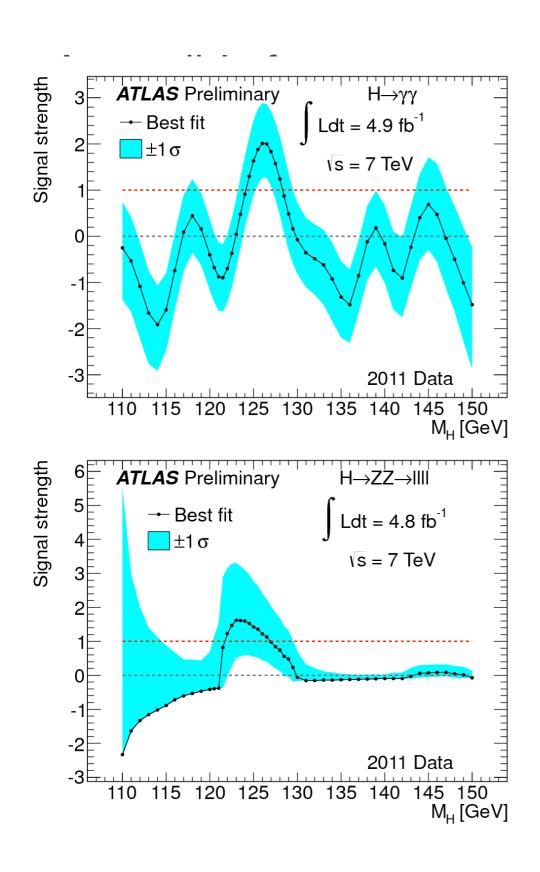
• Assuming  $y_{max}$ =3, a measurement  $R_{ZZ} \approx 0.7$  along with a slight enhancement of the di-photon over the ZZ channel would then imply KK masses  $\approx 8$  TeV, far outside reach for direct production at the LHC (a lower bound  $R_{ZZ} > 0.7$  would imply very strong bounds)

Extended RS model with custodial symmetry:



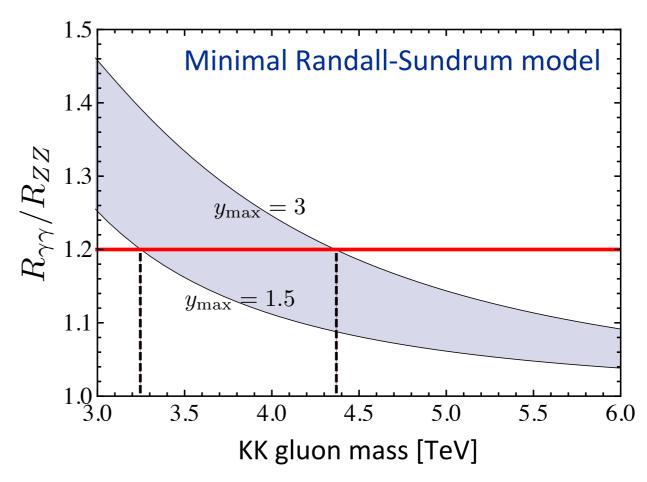


• Even with reduced  $y_{max}=2$ , a measurement  $R_{ZZ}\approx 0.7$  along with a slight enhancement of the di-photon over the ZZ channel would then imply KK masses  $\approx 10$  TeV, far outside reach for direct production at the LHC (a lower bound  $R_{ZZ}>0.7$  would imply very strong bounds)



#### olay:

Consider **double ratios**, which are insensitive to New Physics effects in Higgs production, e.g.:



F. Goertz, U. Haisch, M. Neubert (2011)

### Conclusions

- Warped extra-dimension models provide an appealing framework for addressing the hierarchy problem and the flavor puzzle within the same geometrical approach
- Much like rare FCNC processes, Higgs production in gluon-gluon fusion and Higgs decays into the di-photon final state are loopsuppressed processes, which are sensitive to new heavy particles and probe EWSB sector at the quantum level
- Find that the contribution of the Kaluza-Klein towers of SM quarks is independent of the quark mass and given in terms of fundamental
   5D Yukawa matrices
- Effects are enhanced by the large multiplicity of 5D fermion states and probe regions of parameter space not accessible to direct searches
- Interesting to study other WED models: bulk Higgs, gauge-Higgs unification, generalized Yukawa couplings ... → model dependence?